



中山大學天琴中心

TIANQIN CENTER FOR GRAVITATIONAL PHYSICS, SYSU



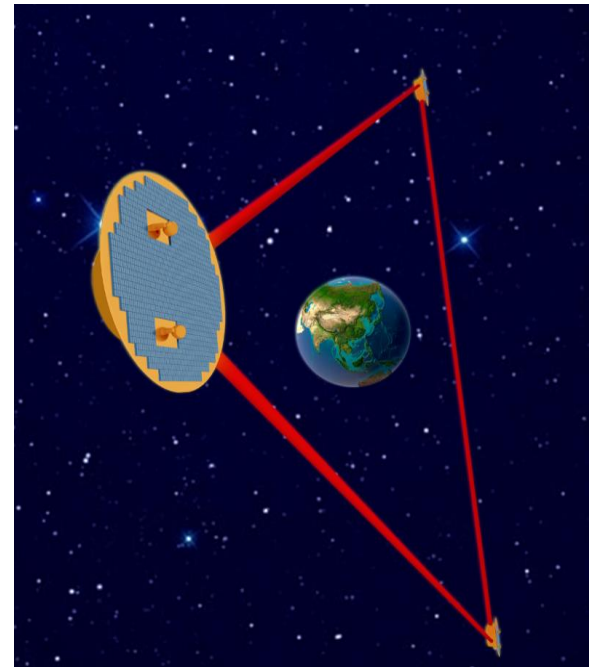
# Fundamental Physics and Cosmology with TianQin

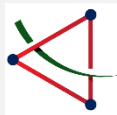
**Fa Peng Huang**

<https://fapenghuang.github.io/group/>

Sun Yat-sen University, TianQin center

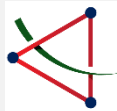
The Rencontres du Vietnam on  
Cosmology @ Quy Nhon, Vietnam 2025.08.12



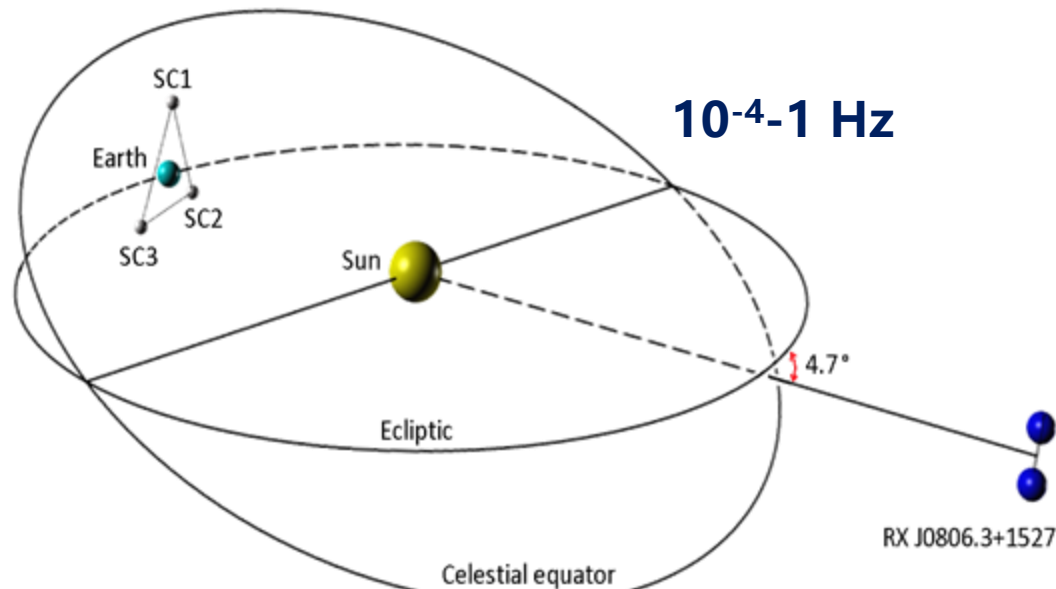


# Outline

1. **TianQin and gravitational wave (GW) in a nutshell**
2. **Higgs potential, Electroweak (EW) phase transition/baryogenesis at TianQin**
3. **Dark matter (DM) signals at TianQin (focus)**
  - ✓ **Ultralight DM**
  - ✓ **Heavy DM**
3. **Summary and outlook**



# What is TianQin ?

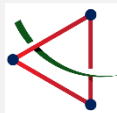


- Expected in 2035
- Geocentric orbit, normal triangle constellation, radius  $\sim 10^5$  km
- Unique frequency band, easier for deployment, tracking, control, and communication

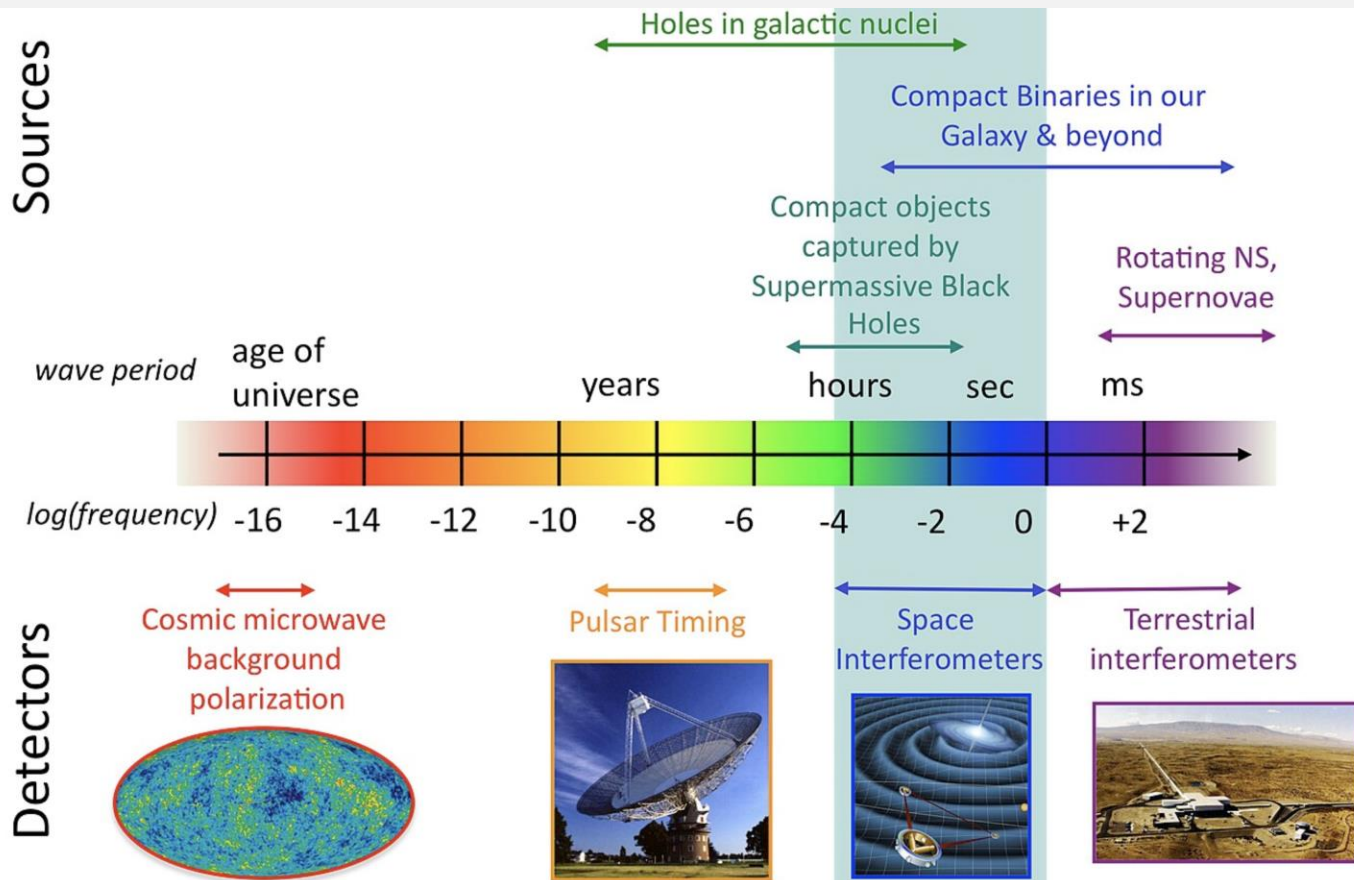


“天琴” (TianQin) “Harpe in space”

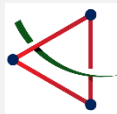
*J. Luo et al. TianQin: a space-borne gravitational wave detector,  
Class. Quant. Grav. 33 (2016) no.3, 035010.*



# Why needs TianQin?

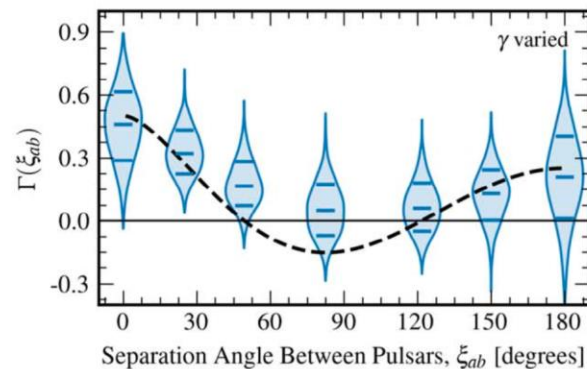
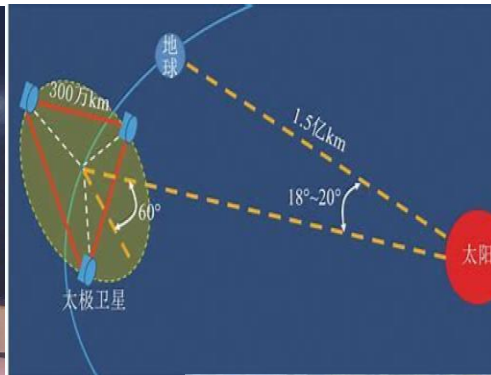


from  
Wiki



# GW experiments

## LISA/TianQin/Taiji ~2035



“TianQin”  
“Harpe in space”

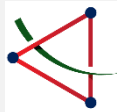
2023 June 29<sup>th</sup>: NANOGrav,  
EPTA, InPTA, Parkes PTA, CPTA



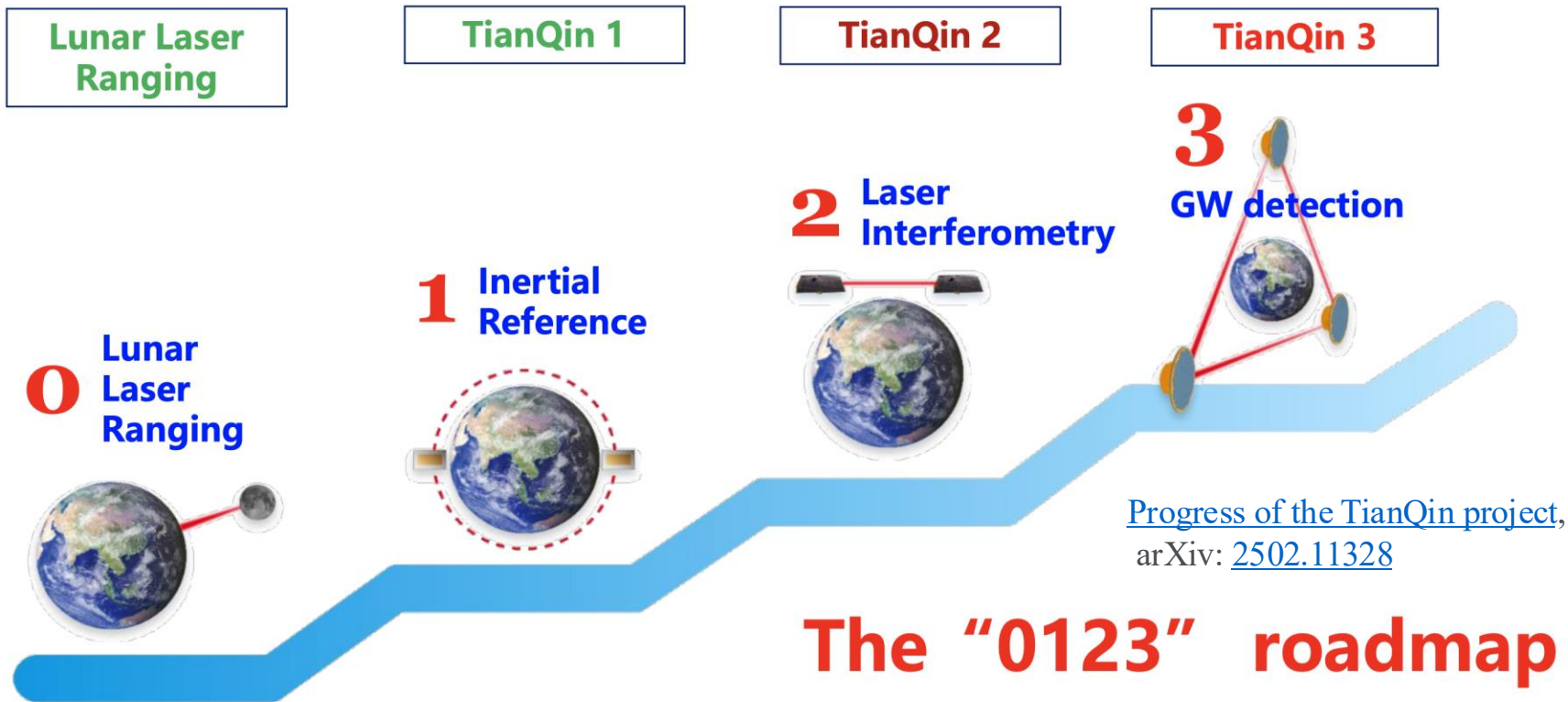
FAST

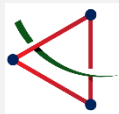


SKA



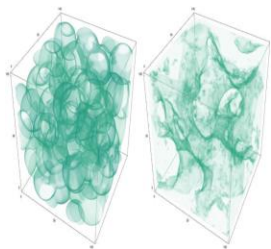
# Working toward TianQin-2035





# GW science at TianQin

## mHZ GW science @TianQin



early universe (EW  
phase transition...)



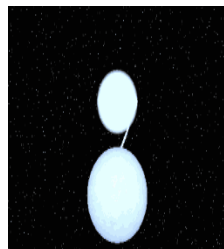
MBHB



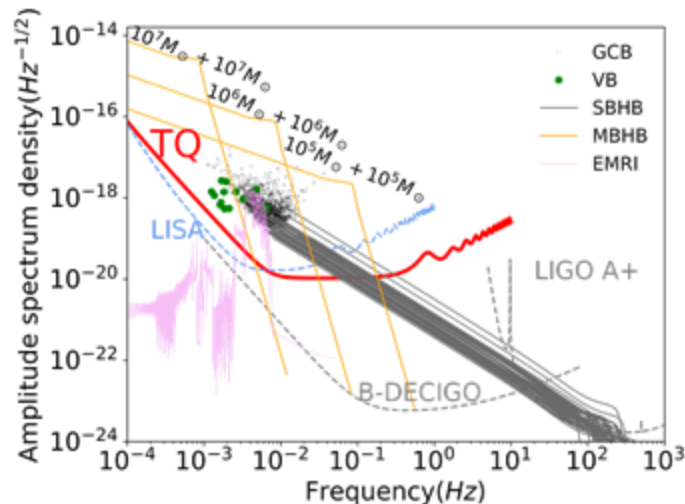
EMRI



SBHB



GCB



*Acta Scientiarum Naturalium Universitatis Sunyatseni,*  
2021, 60(1-2): 1-19

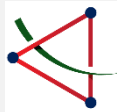
- **Astrophysics:** Origin and evolution of SMBH
- **Fundamental physics:** gravity theory/model
- **cosmology:** dark matter, EW phase transition...

- **Detectable sources**

(event rates)

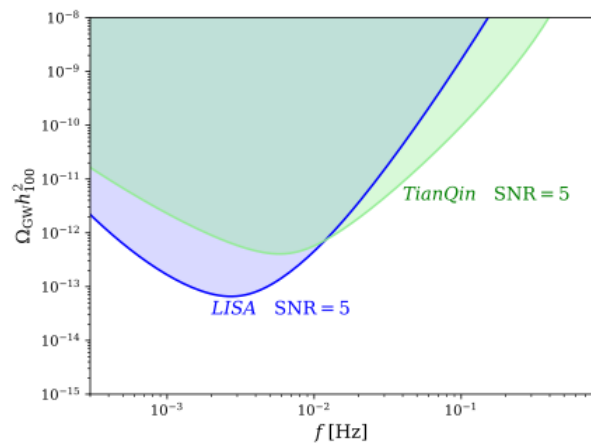
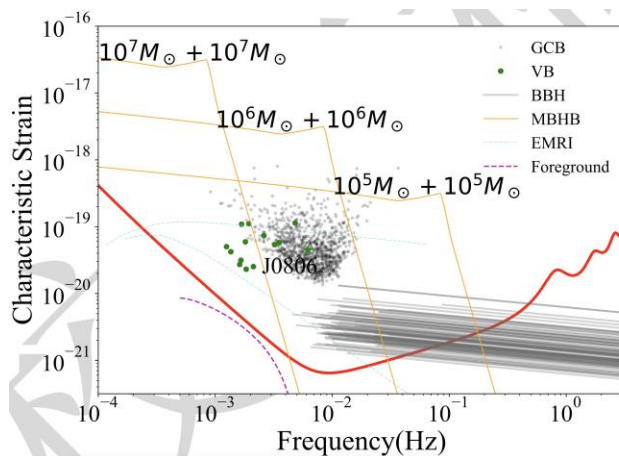
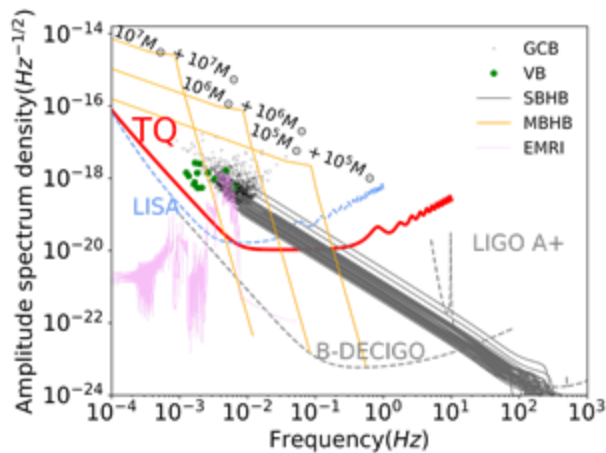
Detectability

- **SNR**
- **Precision of parameter reconstruction**



# TianQin sensitivity/Detectability

$$\Omega_{\text{gw}}(f) = \frac{2\pi^2}{3H_0^2} f^2 h_c^2(f) = \frac{2\pi^2 S_h(f)}{3H_0^2} f^3$$





## Correlation and complementary test with GW

Higgs physics (potential), Z-pole physics, TeV new physics

## Baryon asymmetry of our Universe (EW baryogenesis)

How to probe?  
GW@TianQin

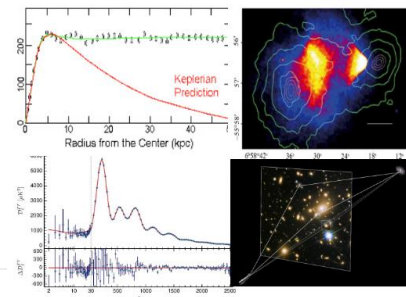
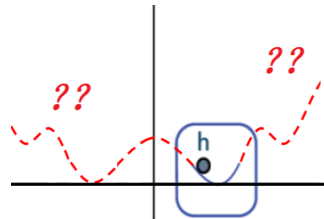
cosmic symmetry  
breaking relics

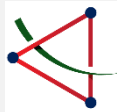
DM formation mechanism;  
New approach to explore DM

[Fundamental Physics and Cosmology with TianQin](#), arXiv: [2502.20138](#)

2025/08/12

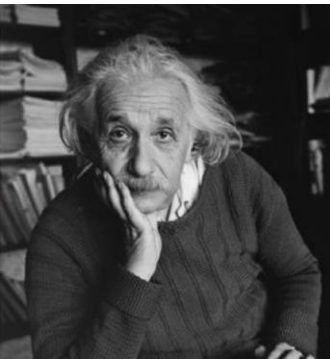
Fa Peng Huang, Fundamental Physics and Cosmology with TianQin





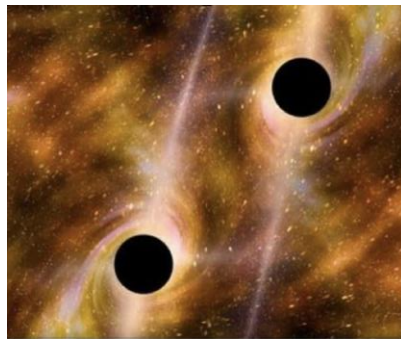
# What is GW ?

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$



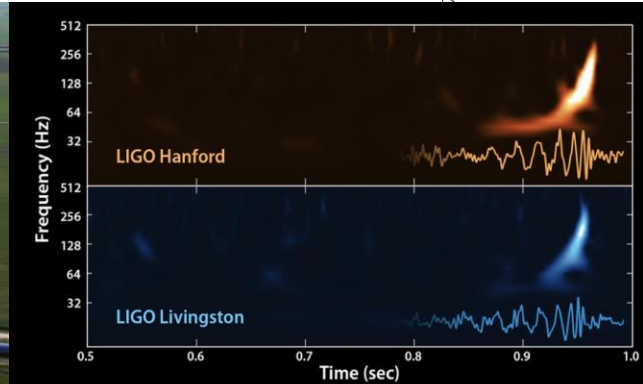
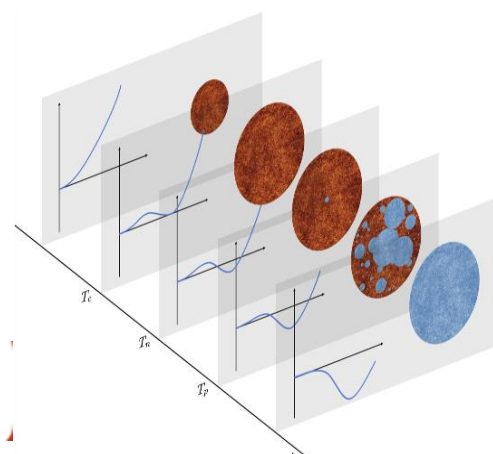
Isolated  
sources:  
quadrupole  
radiation

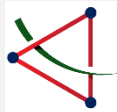
$$h_{ij} \simeq \frac{2G}{c^4 r} \ddot{Q}_{ij}^{TT}(t - r/c)$$



Stochastic  
sources:  
anisotropic  
stress  
tensor

$$\Pi_{ij}(\mathbf{x}, t)$$





# General GW in the early universe

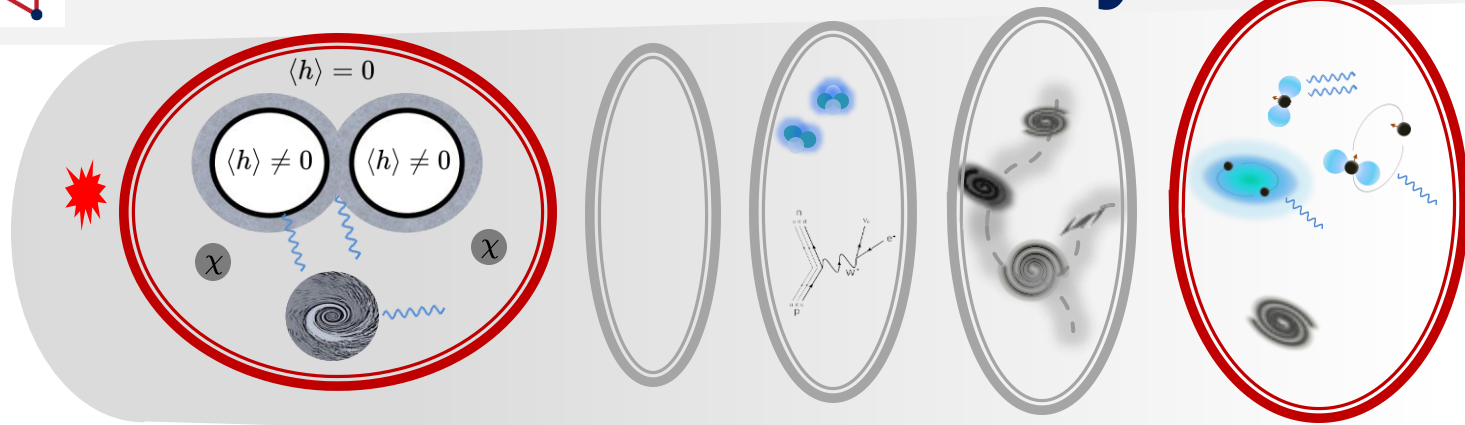
$$\ddot{h}_{ij}(\mathbf{x}, t) + 3H \dot{h}_{ij}(\mathbf{x}, t) - \frac{\nabla^2}{a^2} h_{ij}(\mathbf{x}, t) = 16\pi G \Pi_{ij}(\mathbf{x}, t)$$

- ✓ phase transition: TeV physics (focus)
- ✓ cosmic defects: cosmic string, domain wall...

## Possible sources of **tensor anisotropic stress** in the early universe

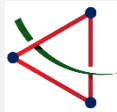
- Scalar field gradients  $\Pi_{ij} \sim [\partial_i \phi \partial_j \phi]^{TT}$  Collisions of bubble walls, cosmic string
- Bulk fluid motion  $\Pi_{ij} \sim [\gamma^2 (\rho + p) v_i v_j]^{TT}$  Sound waves and turbulence
- Gauge fields  $\Pi_{ij} \sim [-E_i E_j - B_i B_j]^{TT}$  Primordial magnetic fields (MHD turbulence)
- Second order scalar perturbations,  $\Pi_{ij}$  from a combination of  $\partial_i \Psi, \partial_i \Phi$  induced GW
- ... [arXiv:1801.04268](https://arxiv.org/abs/1801.04268)

# GW at TianQin from early to late universe



- EW Phase transition GW from EW baryogenesis and DM production induced by Higgs effective potential
- GW from cosmic string
- GW from domain wall
- induced GW
- ....
- various type of astrophysical binaries
- superradiance of axion (DM)
- DM imprints in the waveform of binaries
- ....

..... [Fundamental Physics and Cosmology with TianQin](#), arXiv: [2502.20138](#)

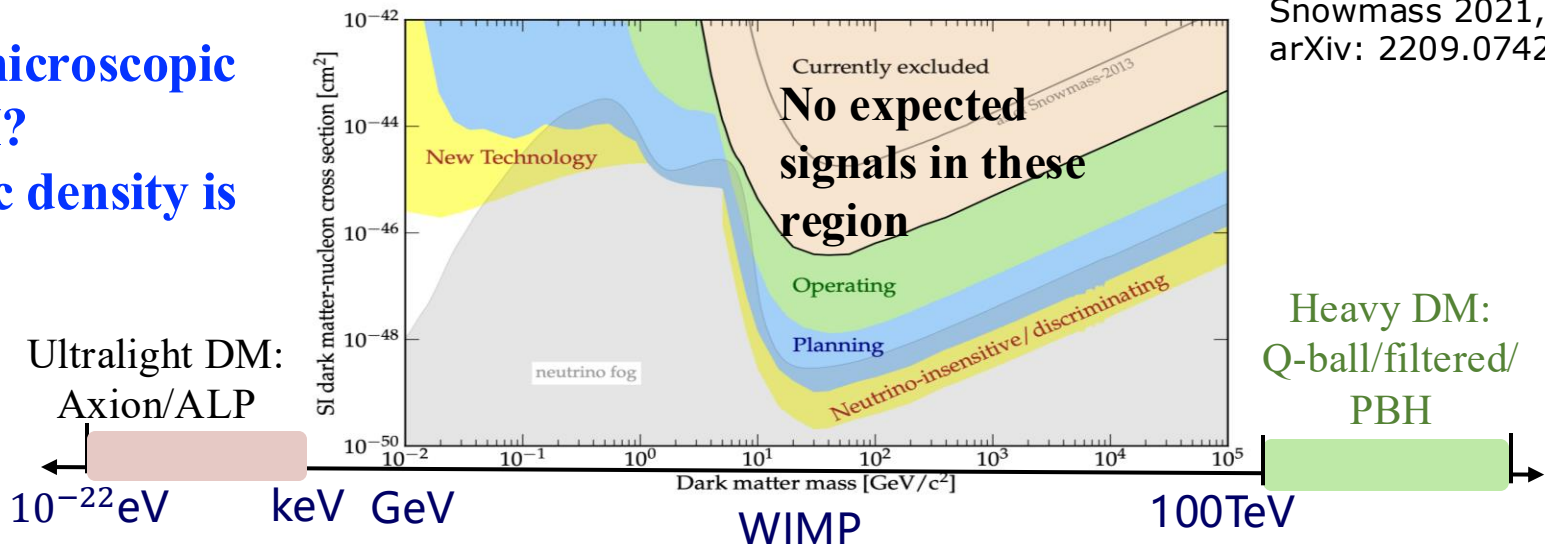


# DM search at TianQin

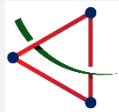
## DM theory and experiments status

arXiv: 1904.07915  
Snowmass 2021,  
arXiv: 2209.07426

What is the microscopic  
nature of DM?  
How DM relic density is  
produced?



- new DM mechanism beyond thermal freeze out: **cosmic phase transition**, Hawking radiation, superradiance...
- new detection method: LISA, **TianQin**, aLIGO, SKA, NanoGrav, Cosmic Explorer, **Einstein telescope**

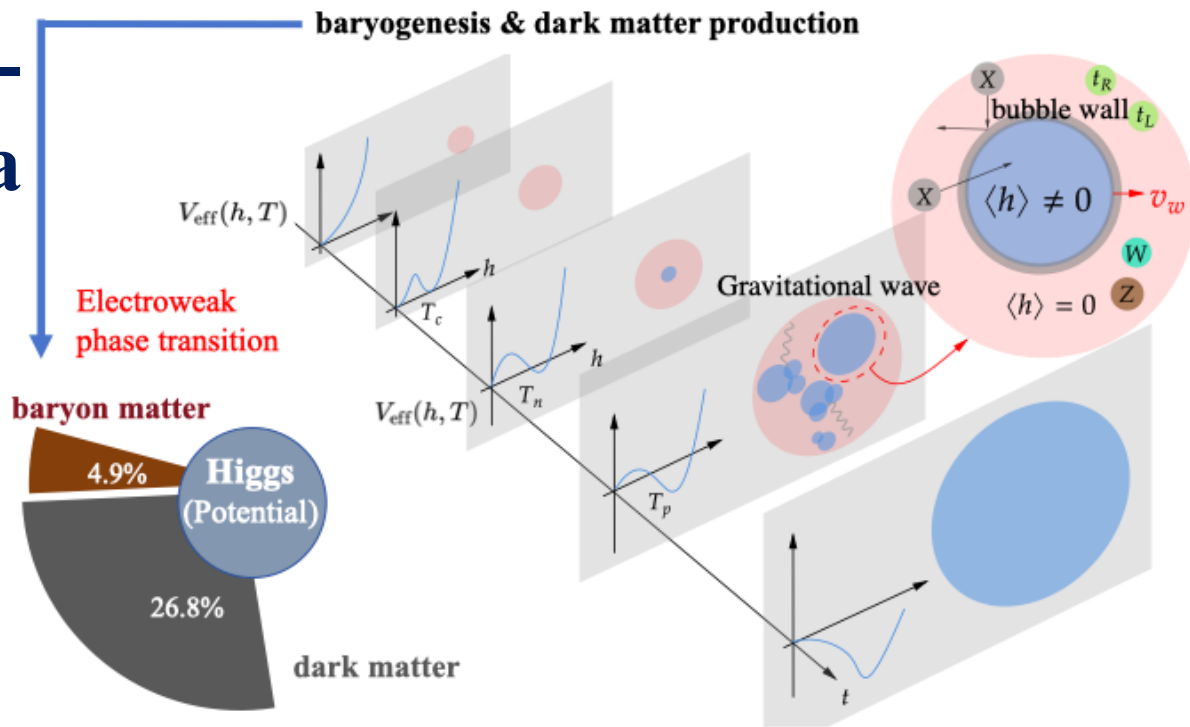


# EW phase transition at TianQin

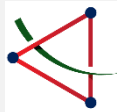
## Cosmology in post-Higgs and GW Era

The observation of **Higgs@LHC** and **GW@LIGO** initiates a new era of exploring DM and baryogenesis by GW.

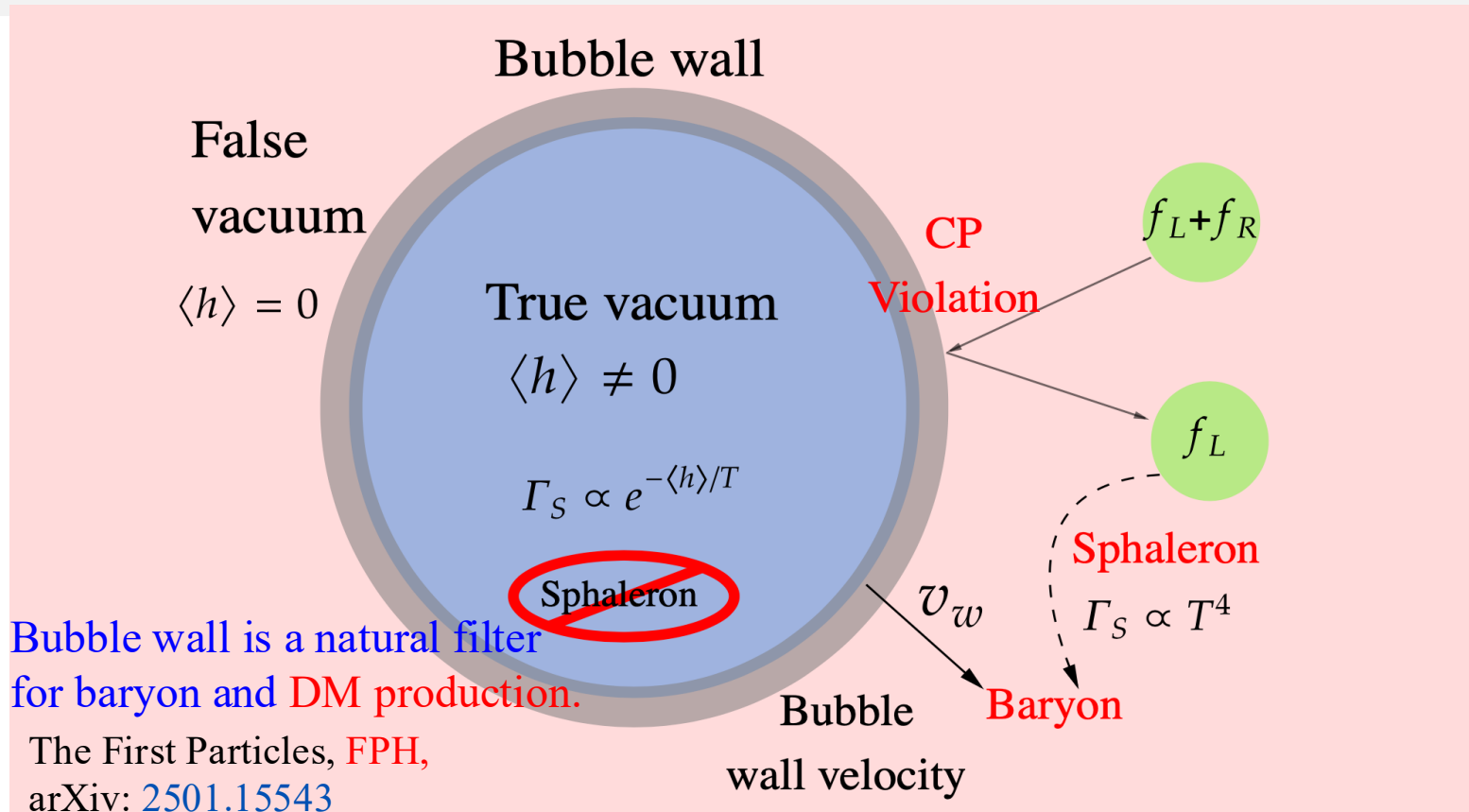
First-order phase transition by Higgs could provide a new approach for DM production.

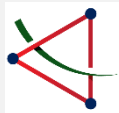


The First Particles, **FPH**, arXiv: [2501.15543](https://arxiv.org/abs/2501.15543)

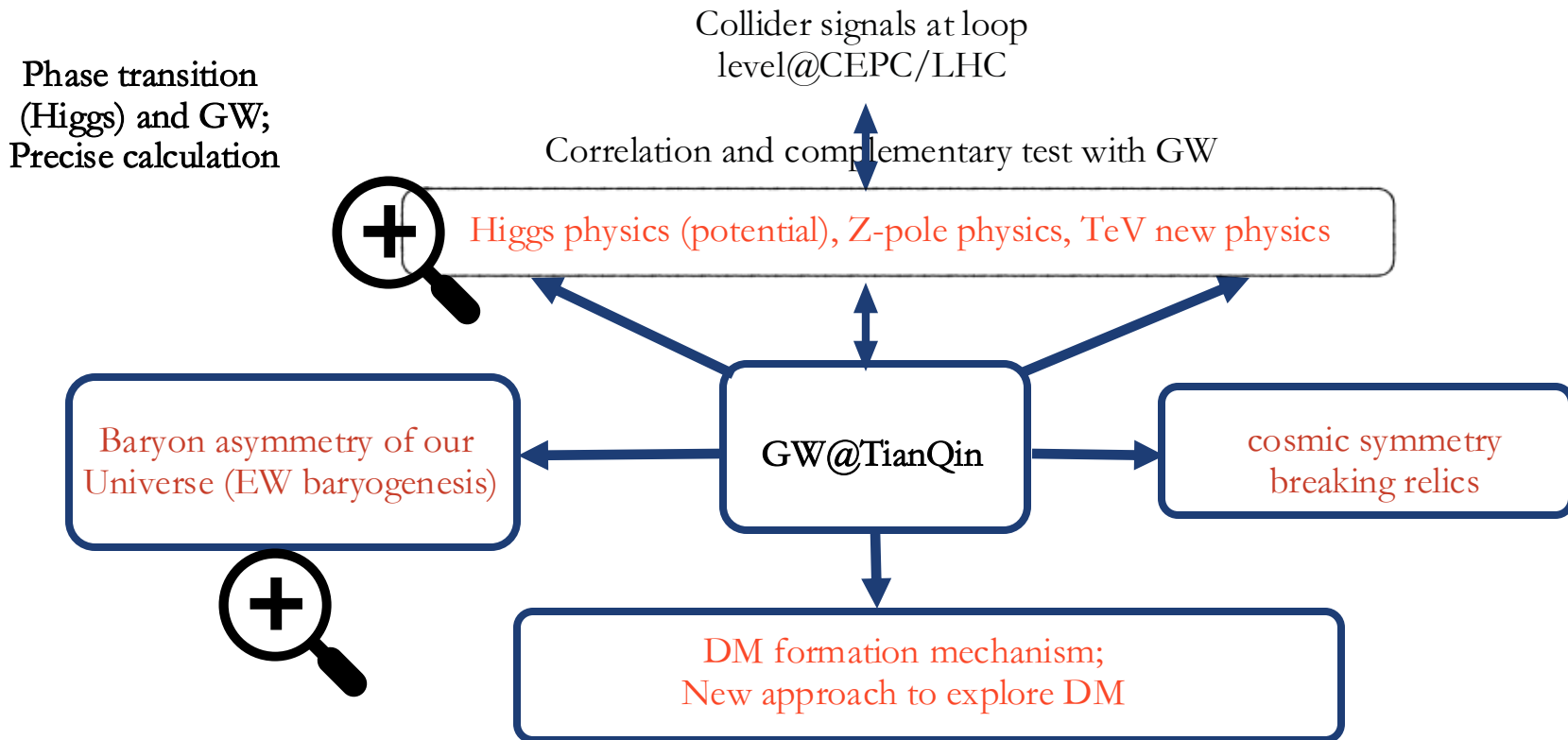


# DM/baryogenesis from phase transition

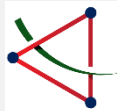




# Particle cosmology at TianQin



[Fundamental Physics and Cosmology with TianQin](#), arXiv: [2502.20138](#)



# EW baryogenesis

The baryon asymmetry of the universe is a long standing problem.

After discovery of Higgs @LHC & GW @aLIGO,  
EW baryogenesis becomes a timely and testable scenario.

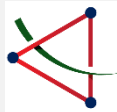


SM technically has all the 3 elements for baryogenesis  
(Sakharov conditions)

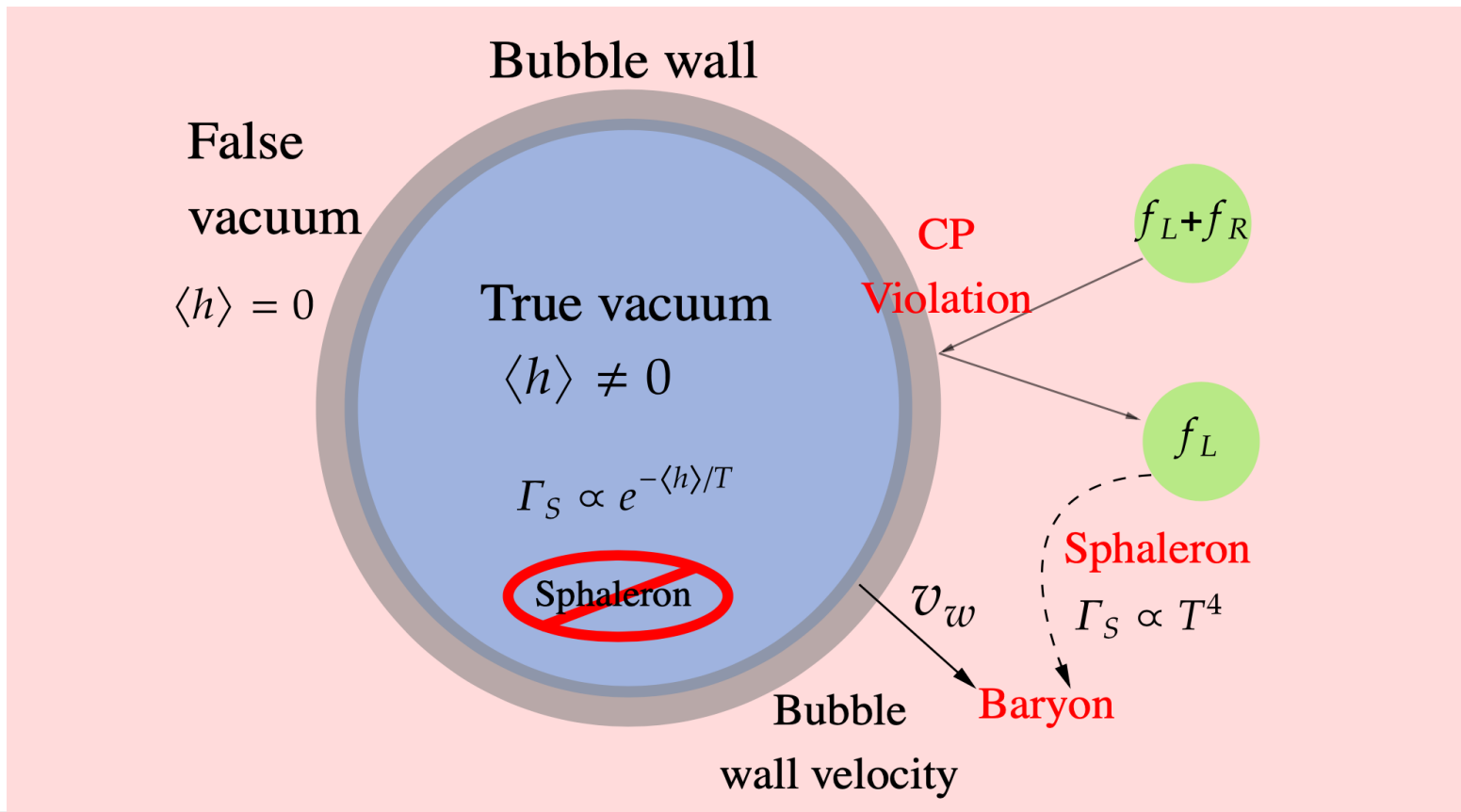
- B violation from anomaly in B+L current
- C and CP-violation: CKM matrix, but too weak, need new CP-violating sources
- **Departure from thermal equilibrium: first-order phase transition with expanding Higgs bubble wall**

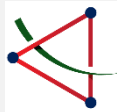
$$\eta_B = n_B/n_\gamma = 5.8 - 6.6 \times 10^{-10}$$

(CMB, BBN)



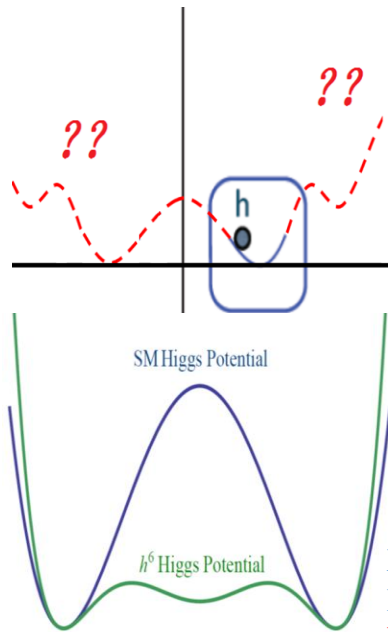
# EW baryogenesis





# Higgs potential

What is the shape of Higgs potential?



Current data tells us nothing but the quadratic oscillation around the VEV 246 GeV with 125 GeV mass.

$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4$$

or 
$$V(h) = \frac{1}{2}\mu^2 h^2 - \frac{\lambda}{4}h^4 + \frac{1}{\Lambda^2}h^6$$

**Produce a first-order EW phase transition,  
deviation of Higgs trilinear coupling and GW**

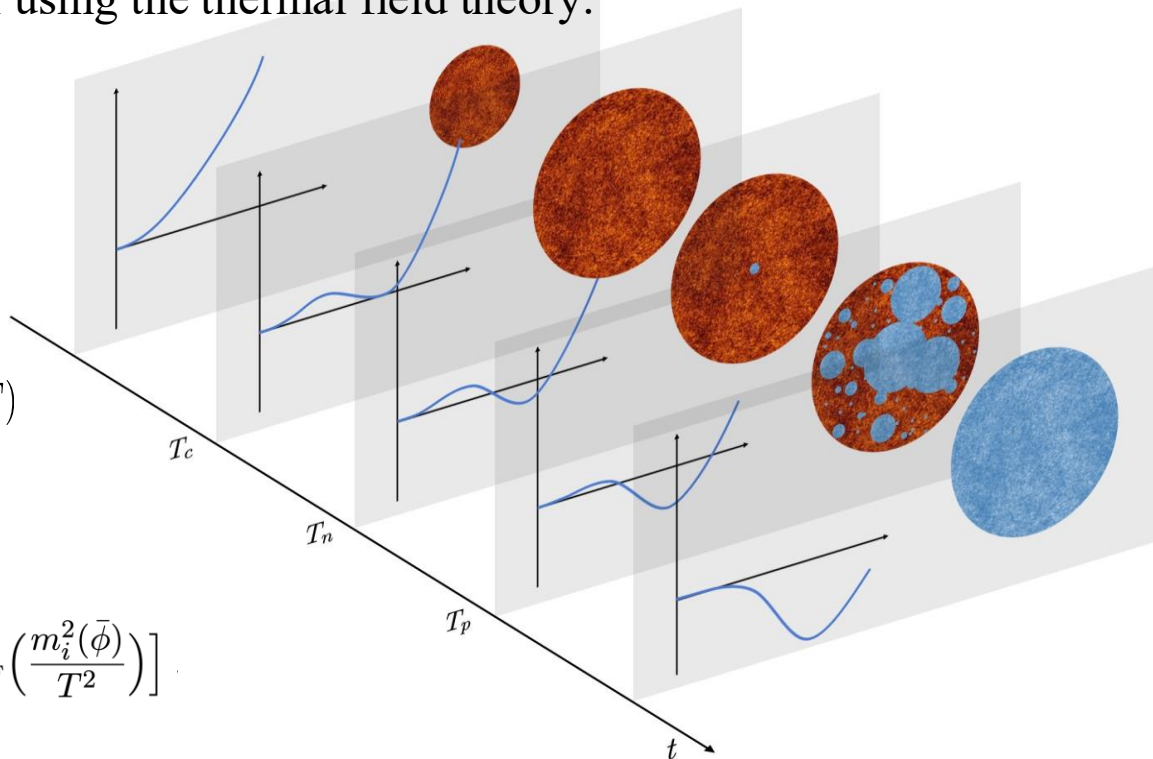
Xinmin Zhang Phys.Rev. D47 (1993) 3065-3067; C. Grojean, G. Servant, J. Well PRD71(2005)036001  
D.J.H. Chung, A. J. Long, Lian-tao Wang Phys.Rev. D87(2013) 023509  
**FPH**, et.al, Phys.Rev.D94(2016)no.4,041702 ; **FPH**, et.al, Phys.Rev.D93 (2016) no.10,103515  
arXiv:1511.06495, Nima Arkani-Hamed et. al.; PreCDR of CEPC; arXiv: [1811.10545](https://arxiv.org/abs/1811.10545), CDR of CEPC



# Phase transition in a nutshell



Calculate the finite-temperature effective potential using the thermal field theory:

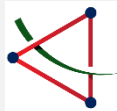


$$\Gamma = \Gamma_0 e^{-S(T)}$$

$$S(T) = \int d^4x \left[ \frac{1}{2} \left( \frac{\partial \phi}{\partial x} \right)^2 + V_{\text{eff}}(\phi, T) \right]$$

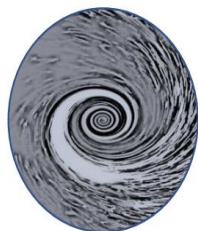
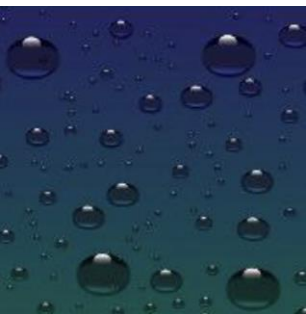
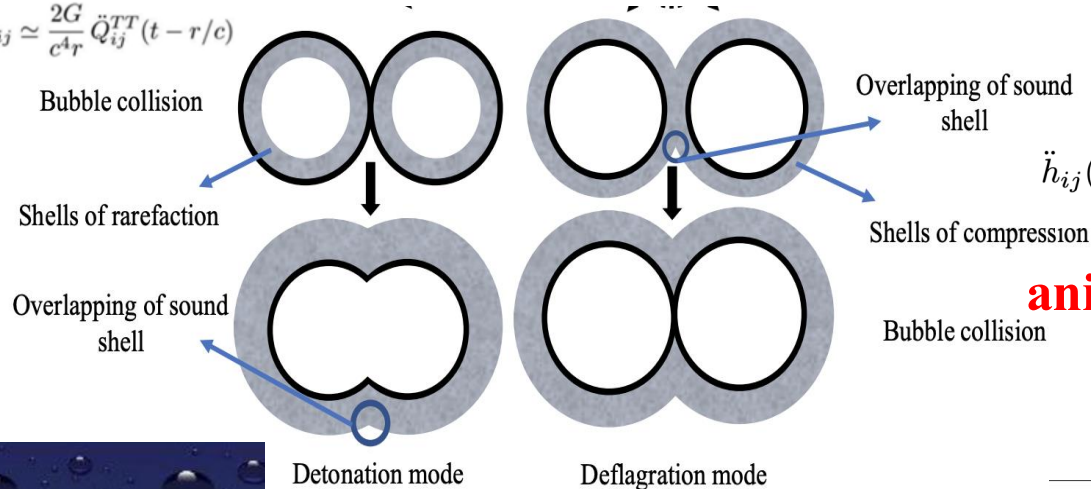
$$V_{\text{eff}}^{(1)}(\bar{\phi}) = \sum_i n_i \left[ \int \frac{d^D p}{(2\pi)^D} \ln(p^2 + m_i^2(\bar{\phi})) + J_{\text{B,F}}\left(\frac{m_i^2(\bar{\phi})}{T^2}\right) \right]$$

Xiao Wang, **FPH**, Xinmin Zhang, JCAP05(2020)045



# Phase transition GW in a nutshell

$$h_{ij} \simeq \frac{2G}{c^4 r} \ddot{Q}_{ij}^{TT}(t - r/c)$$



Turbulence

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$\ddot{h}_{ij}(\mathbf{x}, t) + 3H \dot{h}_{ij}(\mathbf{x}, t) - \frac{\nabla^2}{a^2} h_{ij}(\mathbf{x}, t) = 16\pi G \Pi_{ij}(\mathbf{x}, t)$$

**anisotropic stress tensor:  
source of GW**

General form  $\Pi_{ij}$

$$[\partial_i \phi \partial_j \phi]^{TT}$$

$$[\gamma^2(\rho + p)v_i v_j]^{TT}$$

$$[-E_i E_j - B_i B_j]^{TT}$$

$$\partial_i \Psi, \partial_i \Phi$$

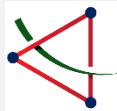
**E. Witten, Phys. Rev. D 30, 272 (1984)**

**C. J. Hogan, Phys. Lett. B 133, 172 (1983);**

**M. Kamionkowski, A. Kosowsky and M. S. Turner, Phys. Rev. D 49, 2837 (1994))**

**EW phase transition  
GW becomes more  
interesting and  
realistic after the  
discovery of  
Higgs by LHC and  
GW by LIGO.**

**Xiao Wang, FPH, Xinmin Zhang, JCAP05(2020)045**



# Phase transition dynamics

In theory, the most important and difficult phase transition parameter for GW, dynamical DM, baryogenesis is bubble wall velocity  $v_w$

In experiments, GW experiment is most sensitive to bubble wall velocity  $v_w$

arXiv: 2404.18703

Aidi Yang, **FPH**, JCAP 2025

*Finite-temperature effective potential*

$$V_{eff}(\phi, T)$$

$\alpha$

$T_p$

$R_* H_*$

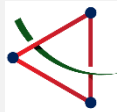
- (1). Daisy resummation problem
- (2). Gauge dependence problem
- (3). No perturbative calculations: lattice calculations and dim-reduction method

*Bubble wall velocity*  
*this talk*  $v_w$

*Energy budget*  
 $\kappa$

S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang, arXiv:2007.10343,  
Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith, arXiv:2009.14295v2  
Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903  
Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005...

F. Giese, T. Konstandin, K. Schmitz and J. van de Vis, arXiv:2010.09744  
Xiao Wang, **FPH** and Xinmin Zhang, Phys.Rev.D 103 (2021) 10, 103520  
Xiao Wang, Chi Tian, **FPH**, JCAP 07 (2023) 006

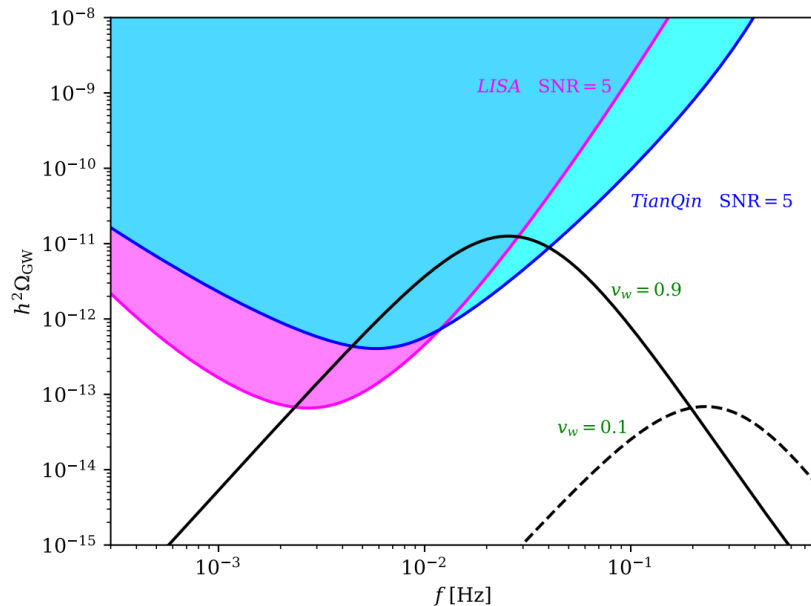
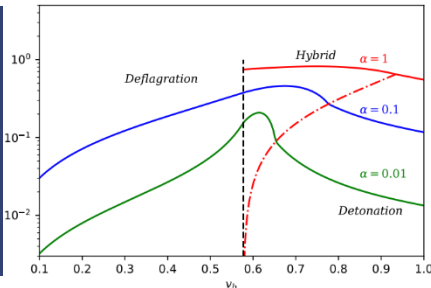
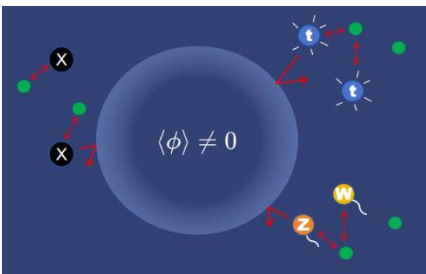


# Bubble wall is essential (like a filter)

The most essential parameter for  
phase transition GW, phase  
transition DM, baryogenesis  $v_w$

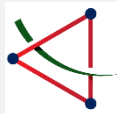
GW detection favor larger  $v_w$   
EW baryogenesis favor smaller  $v_w$   
Dynamical DM is sensitive to  $v_w$

S. Hoche, J. Kozaczuk, A. J. Long, J. Turner and Y. Wang, arXiv:2007.10343,  
Avi Friedlander, Ian Banta, James M. Cline, David Tucker-Smith,  
arXiv:2009.14295v2  
Xiao Wang, **FPH**, Xinmin Zhang, arXiv:2011.12903  
Siyu Jiang, **FPH**, xiao wang, Phys.Rev.D 107 (2023) 9, 095005

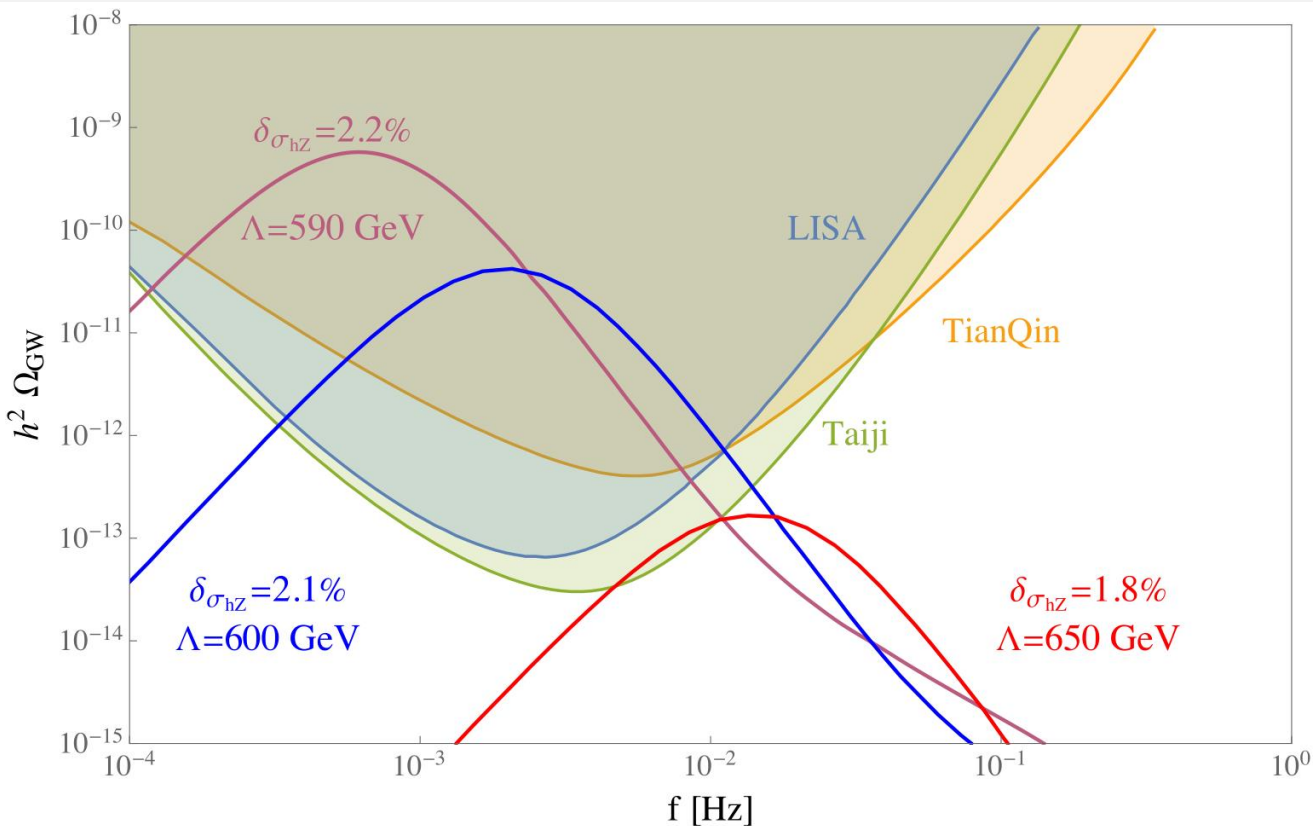


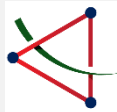
$$\rho_{DM}^4 v_w^{3/4} = 73.5 (2\eta_B s_0)^3 \lambda_S \sigma^4 \Gamma^{3/4}$$

**FPH**, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028;



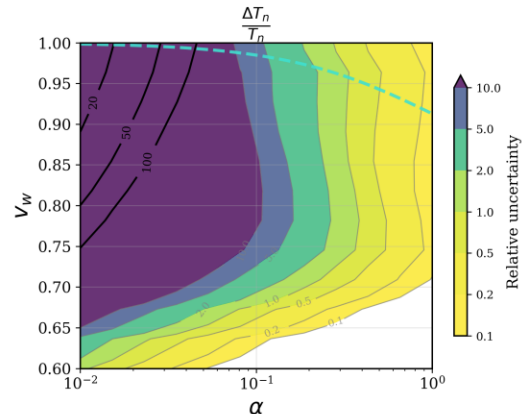
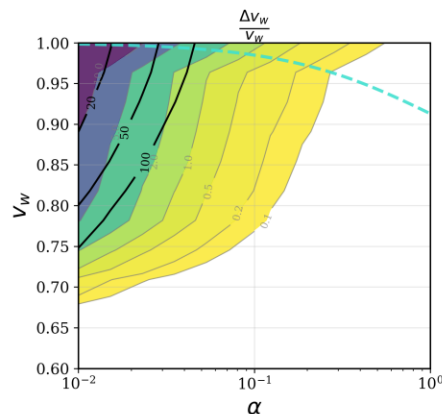
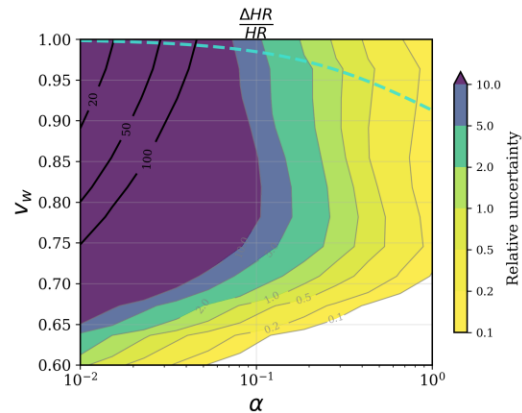
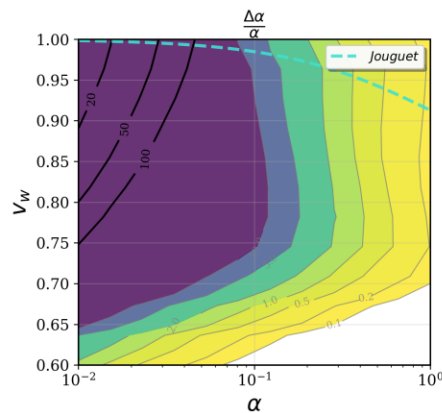
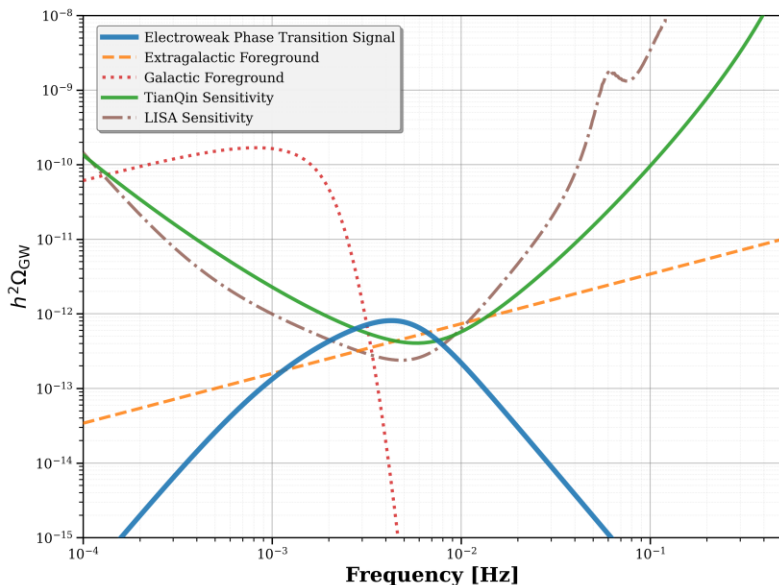
# Higgs potential at TianQin

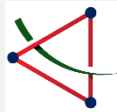




# Reconstruction at TianQin

**Electroweak Phase Transition Detection**





# Particle cosmology at TianQin

Phase transition  
(Higgs) and GW;  
Precise calculation

Collider signals at loop  
level@CEPC/LHC

Correlation and complementary test with GW

Higgs physics (potential), Z-pole physics, TeV new physics

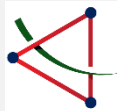
Baryon asymmetry of our  
Universe (EW baryogenesis)

GW@TianQin

cosmic symmetry  
breaking relics

DM formation mechanism;  
New approach to explore DM

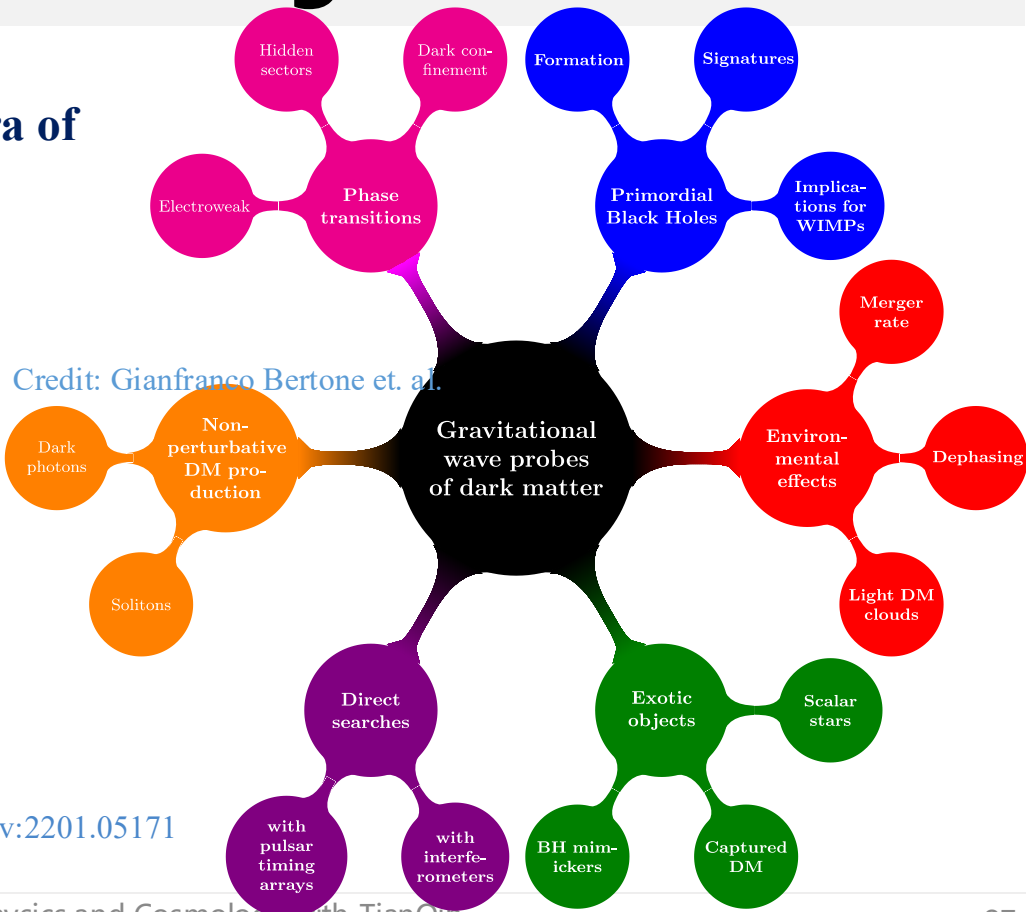
[Fundamental Physics and Cosmology with TianQin](#), arXiv: [2502.20138](#)

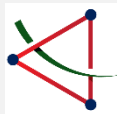


# New DM mechanism/signal at TianQin

- The observation of GW@LIGO and possible hints at PTAs initiate a new era of exploring DM by GW.
- DM can trigger a first-order phase transition in the early universe and detectable GW signals.

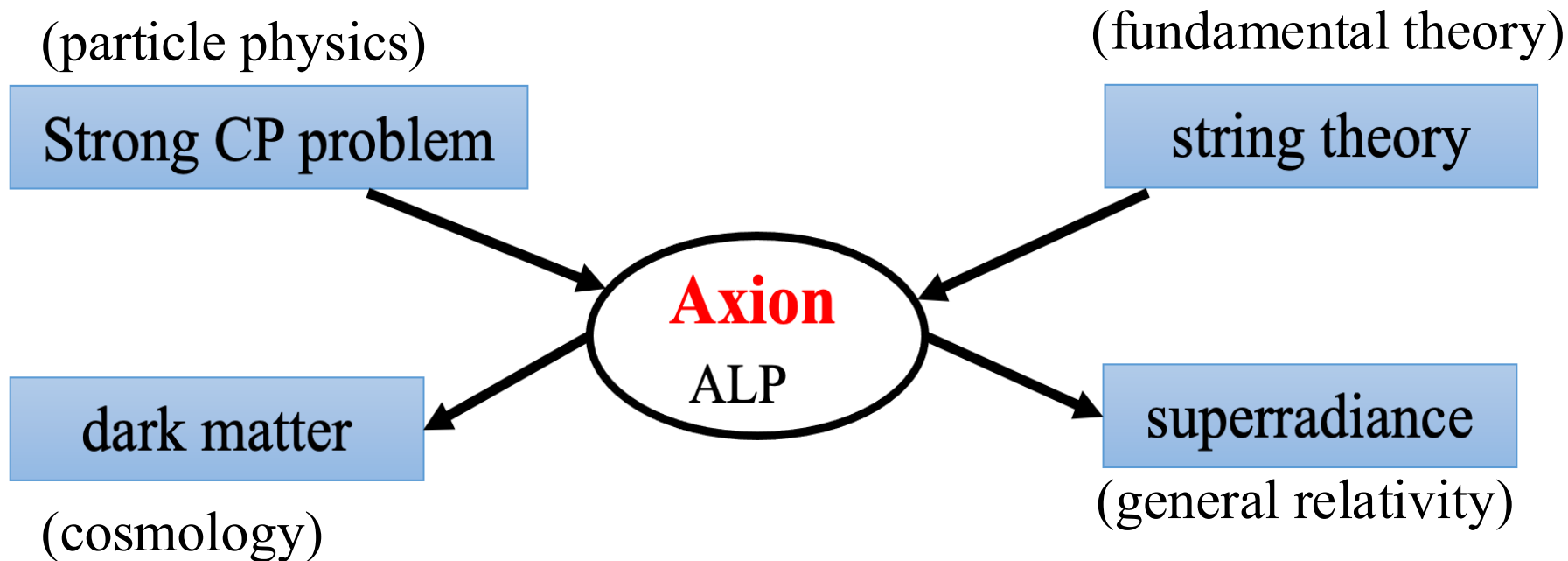
J.Jaekel, V. V. Khoze, M. Spannowsky,  
Phys.Rev. D94 (2016) no.10, 103519  
Zhaofeng Kang, et.al. arXiv:2101.03795, arXiv:2003.02465  
Yan Wang, Chong Sheng Li, and **FPH**, arXiv:2012.03920  
**FPH**, Eibun Senaha Phys.Rev. D100 (2019) no.3, 03501  
**FPH** PoS ICHEP2018 (2019) 397  
**FPH**, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028  
**FPH**, Jiang-Hao Yu, Phys.Rev. D98 (2018) no.9, 095022  
**FPH**, Xinmin Zhang, Phys.Lett. B788 (2019) 288-  
Haipeng An, et.al, arXiv: 2208.14857, arXiv:2009.12381, arXiv:2201.05171  
....



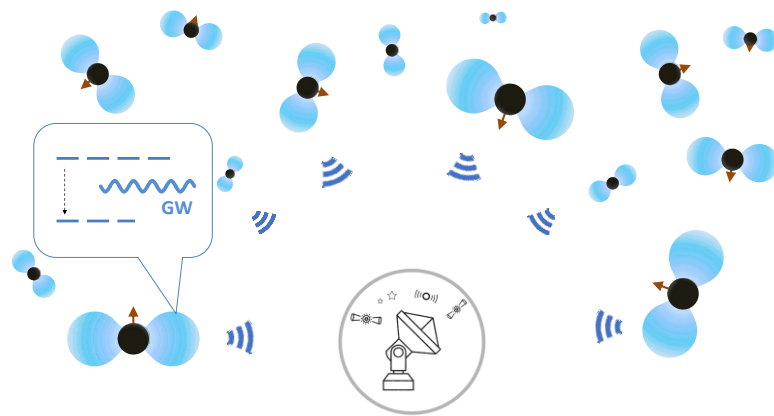
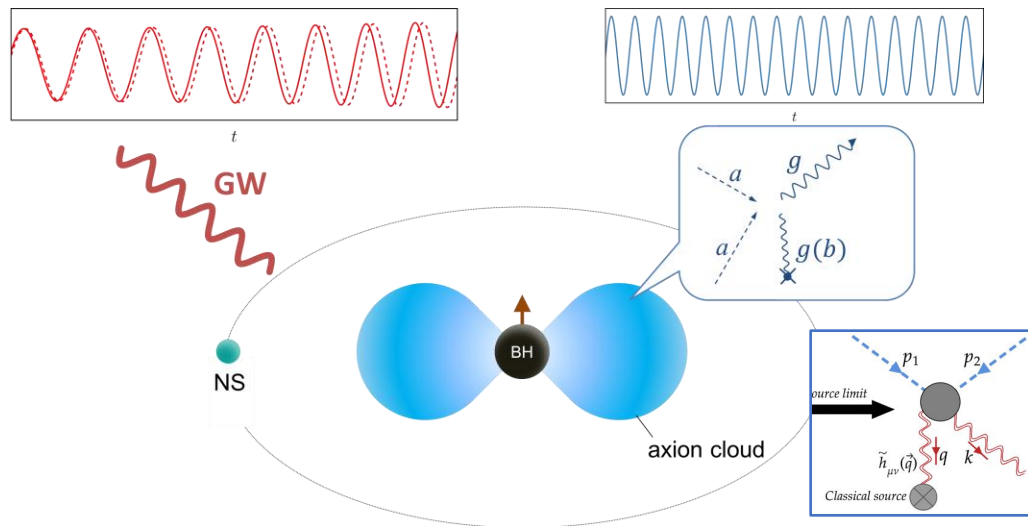


# Axion particle cosmology

Ultralight axion is a promising DM candidate.



# GW of axion from superradiance

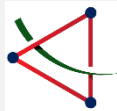


Jing Yang, Ning Xie, **FPH** arXiv:2306.17113,  
arXiv:2404.18703 Aidi Yang, **FPH**

Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024);

Jing Yang, **FPH**

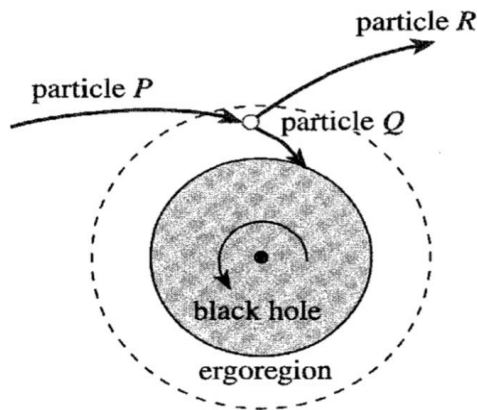
Phys.Rev.D 108 (2023) 10, 103002



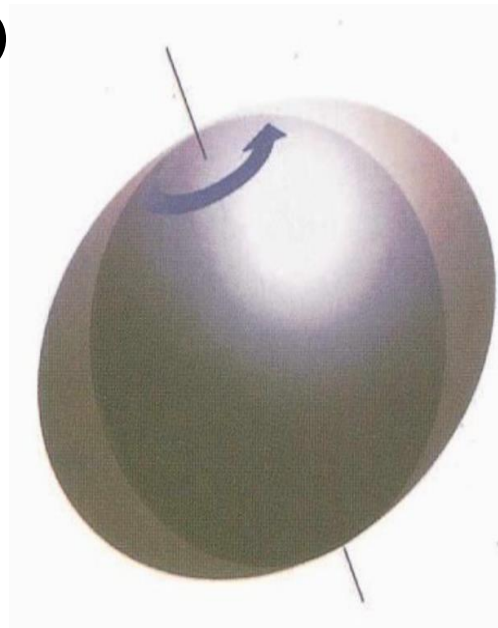
# What is superradiance ?

When Klein (-Gordon) meets Kerr (Dirac Prize 2025)  
superradiance occurs

$$\Delta \frac{d}{dr} \left( \Delta \frac{dR}{dr} \right) + \left[ \omega^2 (r^2 + a^2)^2 - 4aMr m \omega + a^2 m^2 - \Delta (m_a^2 r^2 + a^2 \omega^2 + \lambda) \right] R = 0$$



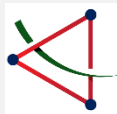
Penrose '69 '71  
Zel'dovich '72  
Starobinsky '73



S. Hawking

Exponential growth solution of Klein-Gordon  
equation due to the boundary condition of Kerr BH.

**Ultralight axion** can form **axion cloud** around rotating  
BH, forming **Gravitational atom (GA)**.



# GW of axion from superradiance

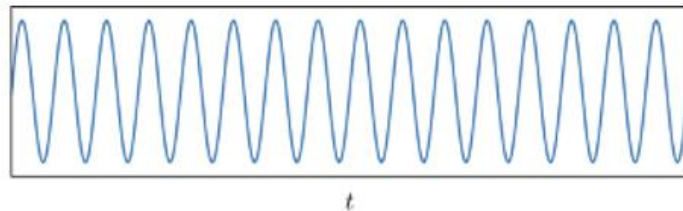
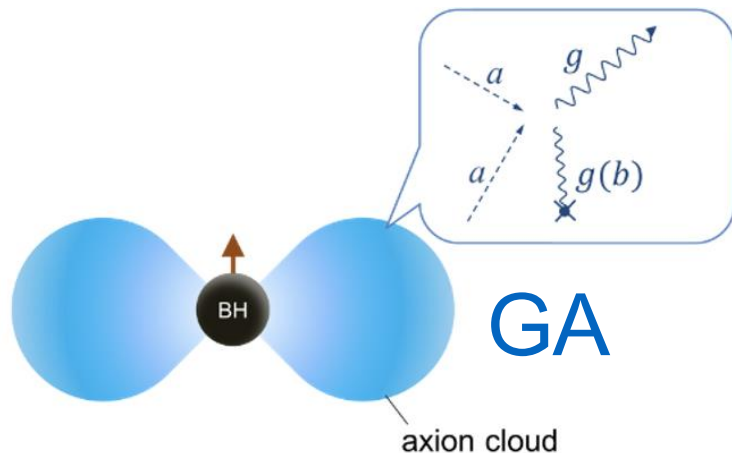
## Axions can annihilate to GWs

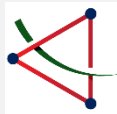
A. Arvanitaki and S. Dubovsky, Phys. Rev. D 83, 044026 (2011)

R. Brito, V. Cardoso and P. Pani, Class. Quant. Grav. 32, no.13, 134001 (2015)

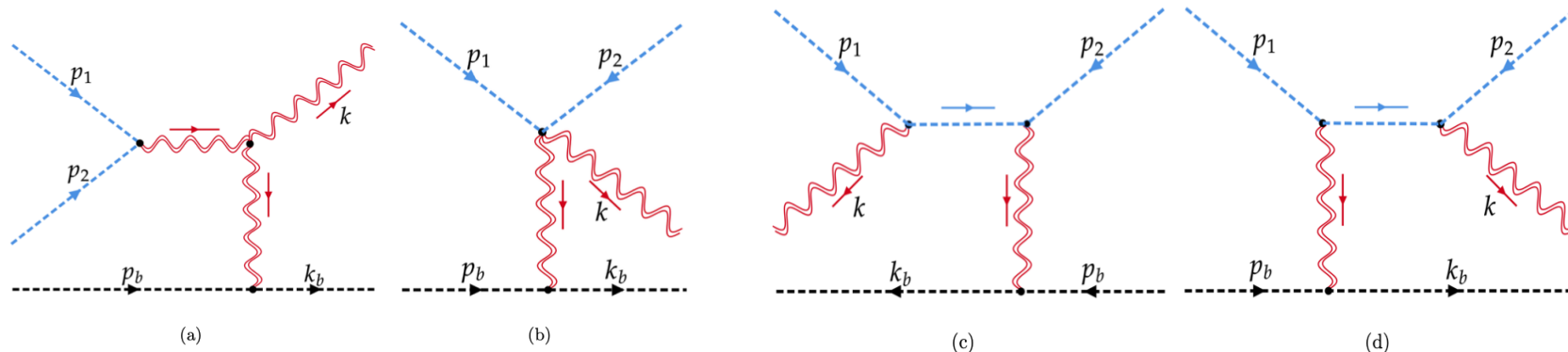
H. Yoshino and H. Kodama, PTEP 2014, 043E02 (2014)

**Jing Yang, FPH**, Phys.Rev.D 108 (2023) 10, 103002

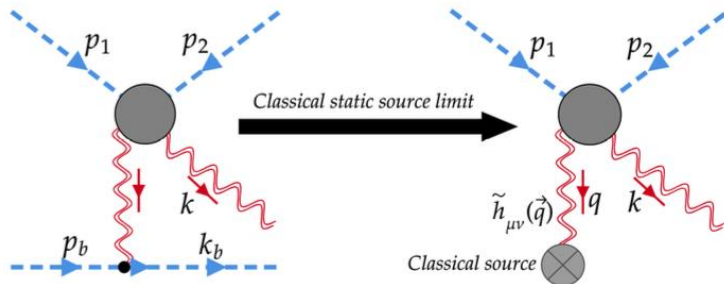




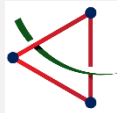
# Microscopic physics



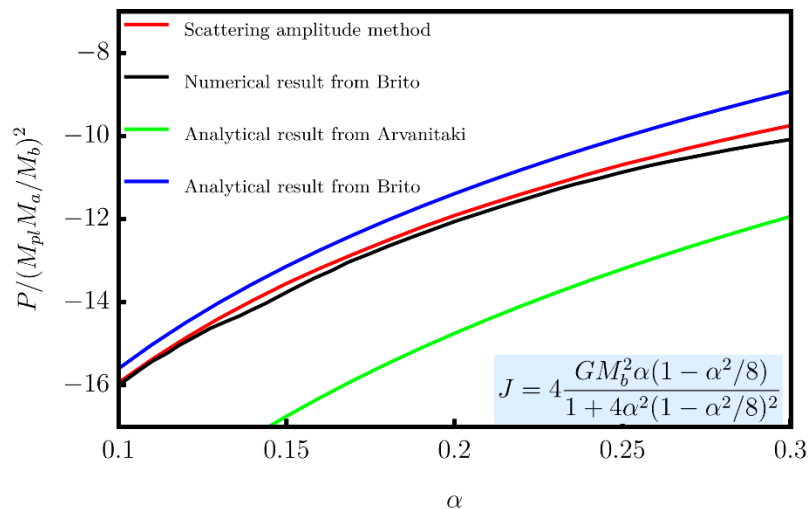
$$M(p_b, p_1, p_2 \rightarrow k, k_b)$$



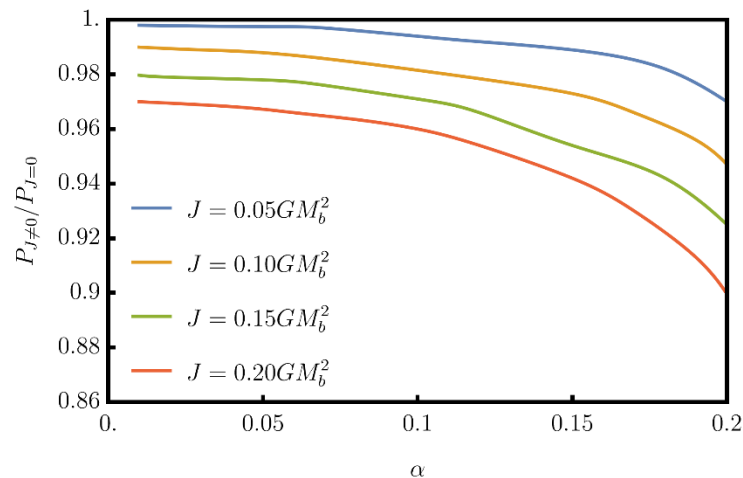
**Jing Yang, FPH,**  
Phys.Rev.D 108 (2023) 10, 103002



# GW radiation from axion annihilation

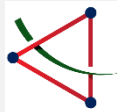


$$\alpha = GM_b m_a \quad M_b = 100 M_{\text{sun}} \quad M_a = M_{\text{sun}}$$



**Jing Yang, FPH,**  
 Phys.Rev.D 108 (2023) 10, 103002

- ✓ monochromatic GW signal  $\omega_{\text{ann}} \sim 2 m_a$
- ✓ gradually depletion of axion cloud (DC) and reduce GA mass



# GW radiation from axion annihilation

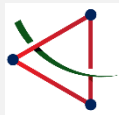
- ✓ Simple and straightforward.
- ✓ Easy to include Kerr metric effects.
- ✓ Microscopic physics is intuitive.
- ✓ It is clearly and simple to demonstrate the analytic approximation formulae.

$$P = \frac{(M_a/\text{GeV})^2 \alpha^{14}}{(M_b/\text{GeV})^6 (2 + \alpha^2)^{11} (4 + \alpha^2)^4} \left[ (M_b/\text{GeV})^4 (9.671 \times 10^{41} + 5.577 \times 10^{42} \alpha^2 + 1.474 \times 10^{43} \alpha^4 + 2.361 \times 10^{43} \alpha^6) + J(M_b/\text{GeV})^2 \alpha (-3.839 \times 10^{80} - 2.111 \times 10^{81} \alpha^2 - 5.329 \times 10^{81} \alpha^4 - 8.165 \times 10^{81} \alpha^8) + J^2 \alpha^2 (3.809 \times 10^{118} + 2.184 \times 10^{119} \alpha^2 + 5.799 \times 10^{119} \alpha^4 + 9.450 \times 10^{119} \alpha^6) \right] \text{GeV}^2.$$

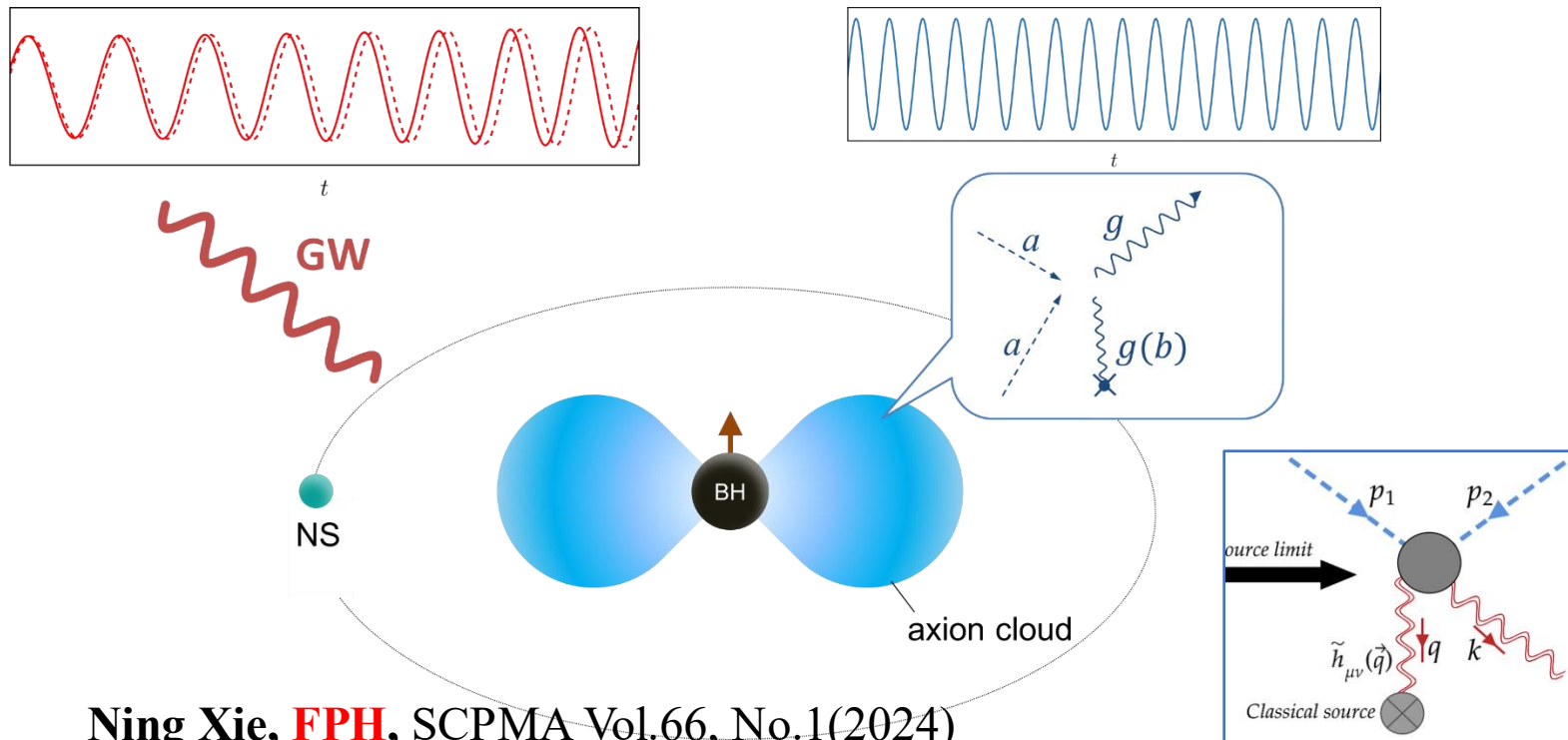
Important for the GW and axion search. More precise calculations and more broad applications are working in progress.

**Jing Yang, FPH,**

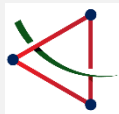
Phys.Rev.D 108 (2023) 10, 103002



# GW of axion from superradiance



Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024)



# GW of axion from superradiance

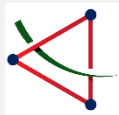
Without ultralight axions  $-\frac{dE_0}{dt} = \mathcal{P}_{\text{GW}} \quad \mathcal{P}_{\text{GW}} = \frac{32}{5} \mu^2 r^4 \omega^6$

With ultralight axions

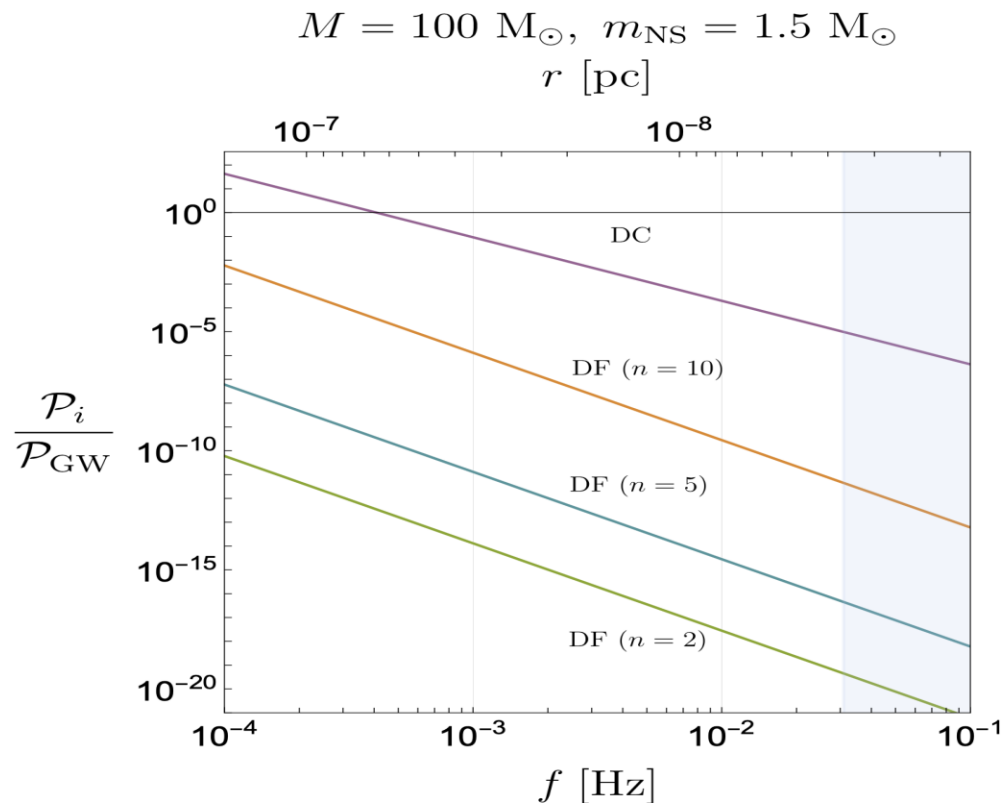
$$-\frac{dE}{dt} = (\mathcal{P}_{\text{GW}} + \boxed{\mathcal{P}_{\text{DC}}} + \mathcal{P}_{\text{DF}} + \mathcal{P}_{\text{DR}})$$

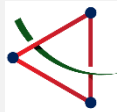
dynamical friction (DF), depletion of axion cloud (DC), dipole radiation(DR)

**Ning Xie, FPH**, SCPMA Vol.66, No.1(2024)



# GW of axion from superradiance



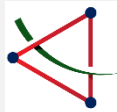


# GW of axion from superradiance

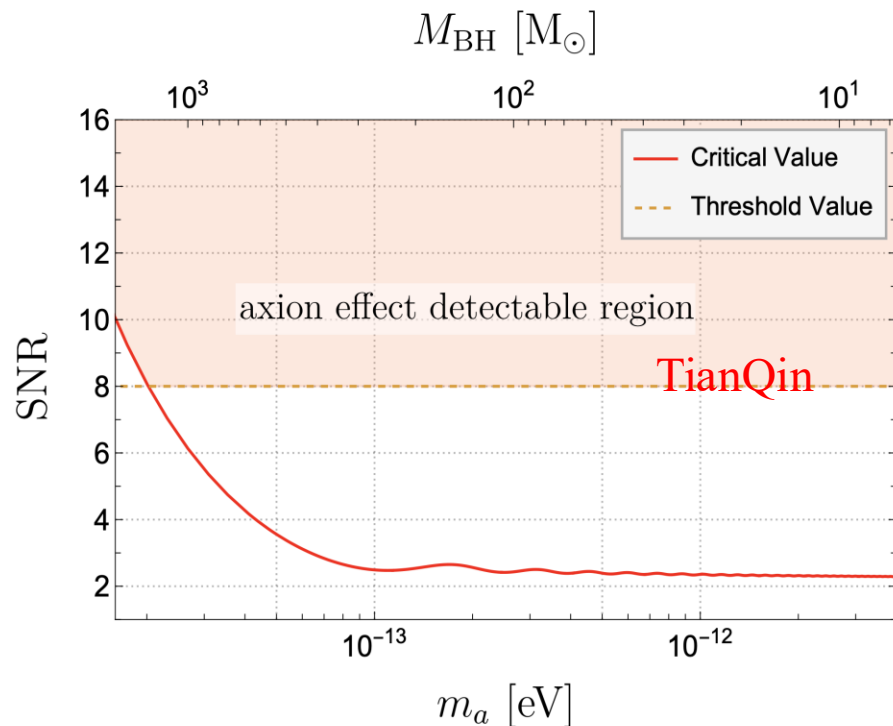
$$\frac{dr}{dt} = \left( -\frac{M m_{\text{NS}}}{2r^2} \right)^{-1} (\mathcal{P}_{\text{GW}} + \mathcal{P}_{\text{DC}} + \mathcal{P}_{\text{DF}} + \mathcal{P}_{\text{DR}})$$

$$\Delta\phi \sim 15\pi \left( \frac{m_a}{10^{-12} \text{ eV}} \right) \left( \frac{f_T}{10^{-2} \text{ Hz}} \right) \left( \frac{T}{5 \text{ yrs}} \right)^2$$

Ning Xie, **FPH**, SCPMA Vol.66, No.1(2024)



# Complementary search: TianQin+PTA

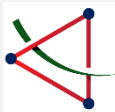


Axions modify the rate of binary period change

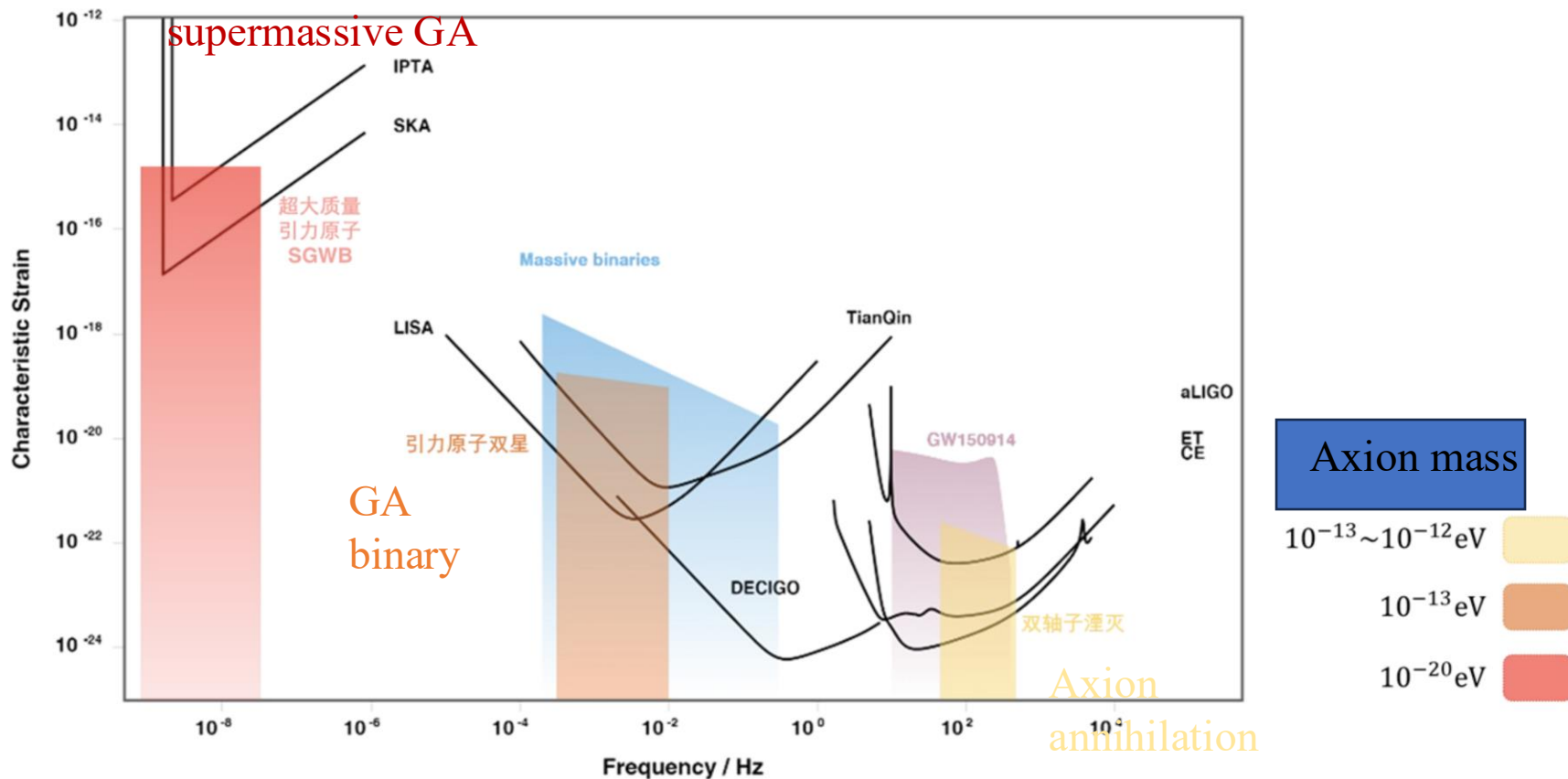
$$\Delta\dot{P} = \left| \dot{P} - \dot{P}_{\text{vac}} \right| \approx 10^{-12} \text{ s/s}$$

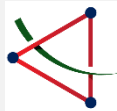
Future Pulsar timing measurement precision, such as SKA

$$10^{-15} \text{ s/s}$$



# GW detection of ultralight axion





# Heavy DM from cosmic phase transition

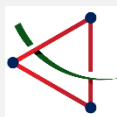
Renaissance of quark nugget DM idea by E. Witten.

Recently, dynamical DM formed by phase transition has become a new idea for heavy. Bubble wall in can be the “filter” to obtain the needed heavy DM when avoiding the unitarity constraints.

E. Krylov, A. Levin, V. Rubakov, Phys.Rev.D 87 (2013) 8, 083528  
FPH, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028  
arXiv:1912.04238, Dongjin Chway, Tae Hyun Jung, Chang Sub Shin  
Phys.Rev.Lett. 125 (2020) 15, 151102 , M. J. Baker, J. Kopp, and A. J. Long  
arXiv:2101.05721, Aleksandr Azatov, Miguel Vanvlasselaer, Wen Yin  
arXiv:2103.09827, Pouya Asadi , Eric D. Kramer, Eric Kuflik, Gregory W. Ridgway, Tracy R. Slatyer, J. Smirnov  
arXiv:2103.09822, Pouya Asadi , Eric D. Kramer, Eric Kuflik, Gregory W. Ridgway, Tracy R. Slatyer, J. Smirnov  
Siyu Jiang, FPH, Chong Sheng Li, arXiv:2305.02218  
Siyu Jiang, FPH, Pyungwon Ko, arXiv:2404.16509  
more than 100 papers in recent 5 years



| phase transition in the early universe | Coffee making process |
|----------------------------------------|-----------------------|
| Bubble wall                            | filter                |
| Case I: (gauged) Q-ball DM             | Large coffee beans    |
| Case II: filtered DM                   | Coffee                |
| Phase transition GW                    | Aroma                 |



# Case I: Q-ball DM

# What is Q-ball?

PHYSICS REPORTS (Review Section of Physics Letters) 221, Nos. 5 & 6 (1992) 251-350, North-Holland

PHYSICS REPORTS

Nuclear Physics B262 (1985) 263-283  
© North-Holland Publishing Company

Nontopological solitons\*

T.D. Lee

*Department of Physics, Columbia University, New York, NY 10027, USA*

and

Y. Pang

*Brookhaven National Laboratory, Upton, NY 11973, USA*

Received May 1992; editor: D.N. Schramm

**Q-BALLS\***

Sidney COLEMAN

*Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138, USA*

Q-ball is the most typical non-topological soliton, initially proposed by Prof. Tsung-Dao Lee and Sidney Coleman. In quantum field theory, a spherically symmetric extended body that forms a non-topological soliton structure with a conserved global quantum number  $Q$  is called a Q-ball.

$$\phi = (\phi_R + i\phi_I)/\sqrt{2} \quad Q = \int j^0 dx = \int (\phi_I \dot{\phi}_R - \phi_R \dot{\phi}_I) dx.$$

$$\delta(E - \omega Q) = 0$$



$$E = \int \left\{ \frac{1}{2} [\dot{\phi}_R^2 + \dot{\phi}_I^2 + (\nabla \phi_R)^2 + (\nabla \phi_I)^2] + U \left[ \frac{1}{2} (\phi_R^2 + \phi_I^2) \right] \right\} dx$$

$$\phi = f(r)e^{-i\omega t}$$

# Q-ball production mechanism

Q-ball production :

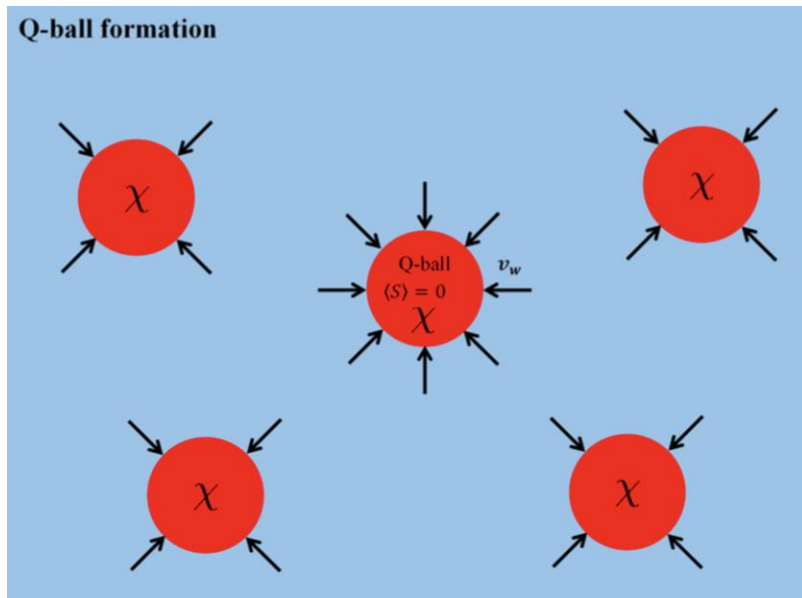
- (1) produce the charge asymmetry (i.e. locally produce lots of particles with the same charge to form Q-ball)
- (2) and packet the same sign charge in the small size after overcoming the Coulomb repulsive interaction.

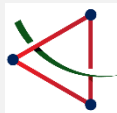
1. Supersymmetry? Affleck-Dine mechanism.

We do not observe the supersymmetry until now!

2. Q-ball formation based on first-order phase transition. 

This talk

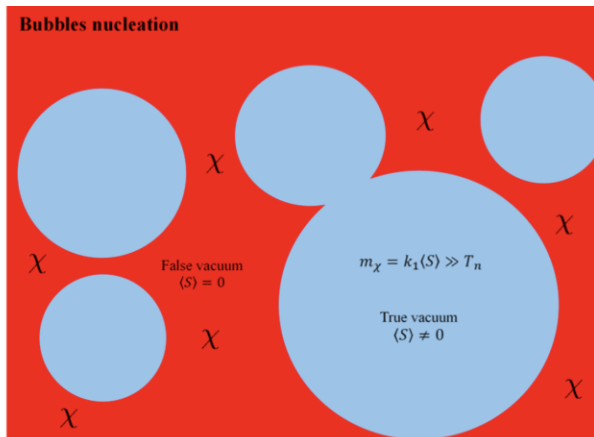




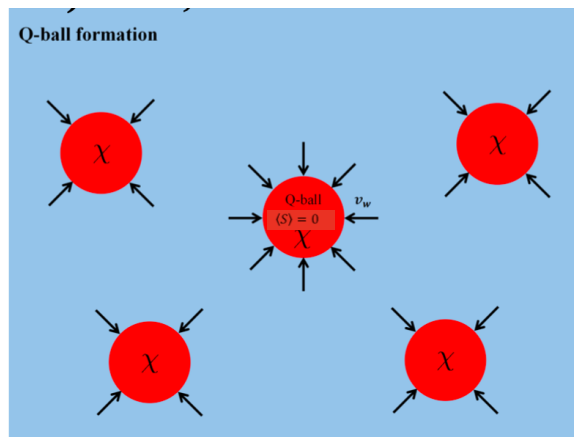
# Case I: Q-ball DM

**Global Q-ball DM:** The cosmic phase transition with Q-balls production can explain baryogenesis and DM simultaneously.

$$\rho_{DM}^4 v_w^{3/4} = 73.5 (2\eta_B s_0)^3 \lambda_S \sigma^4 \Gamma^{3/4}$$



(a) Bubble nucleation:  $\chi$  particles trapped in the false vacuum due to Boltzmann suppression



(b) Q-ball formation: After the formation of Q-balls, they should be squeezed by the true vacuum

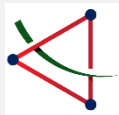


New DM production scenario by the bubbles.  
The global Q-ball model proposed by T.D. Lee

Friedberg-Lee-Sirlin model

[R. Friedberg, T.D. Lee and A. Sirlin. Rev. D 13 \(1976\) 2739](#)

**FPH**, Chong Sheng Li, Phys.Rev. D96 (2017) no.9, 095028;



# Case I: Gauged Q-ball DM

$$\langle h \rangle \neq 0$$

$$\langle \phi \rangle = 0$$

$$\langle h \rangle = 0$$

$$\langle \phi \rangle \neq 0$$

$$\langle A \rangle \neq 0$$

When the conserved U(1) symmetry is **local**,  
This introduces an extra **gauge field A**.

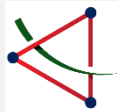
The **minimal model** achieving

$$\mathcal{L} = (D_\mu \phi)^\dagger (D^\mu \phi) + \frac{1}{2} \partial_\mu h \partial^\mu h - \frac{1}{4} \tilde{A}_{\mu\nu} \tilde{A}^{\mu\nu} - V(\phi, h)$$

$$V(\phi, h) = \frac{\lambda_{\phi h}}{2} h^2 |\phi|^2 + \frac{\lambda_h}{4} (h^2 - v_0^2)^2$$

Interestingly, this portal coupling also  
naturally induces a strong first-order phase  
transition.

Siyu Jiang, **FPH**, Pyungwon Ko, JHEP 07 (2024) 053



# Gauged Q-ball

$$\mathcal{L} = (D_\mu \phi)^\dagger (D^\mu \phi) + \frac{1}{2} \partial_\mu h \partial^\mu h - \frac{1}{4} \tilde{A}_{\mu\nu} \tilde{A}^{\mu\nu} - V(\phi, h)$$

$$V(\phi, h) = \frac{\lambda_{\phi h}}{2} h^2 |\phi|^2 + \frac{\lambda_h}{4} (h^2 - v_0^2)^2$$

$$\tilde{A}_t(r) = v_0 \frac{\tilde{g}}{\sqrt{2\lambda_h}} \mathcal{A}(\rho), \quad \phi(t, r) = \frac{v_0}{\sqrt{2}} \Phi(\rho) e^{-i\omega t}, \quad h(r) = v_0 \mathcal{H}(\rho)$$

Friedberg-Lee-Sirlin-Maxwell model

$$\frac{1}{\rho^2} \partial_\rho (\rho^2 \partial_\rho \mathcal{A}) + (\nu - \alpha^2 \mathcal{A}) \Phi^2 = 0,$$

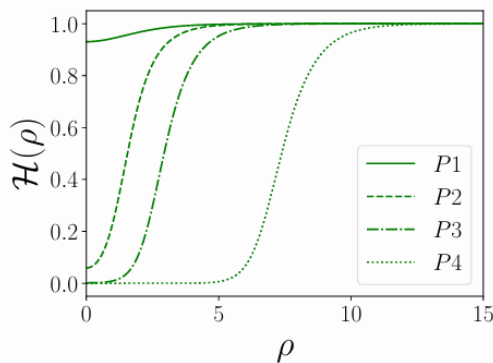
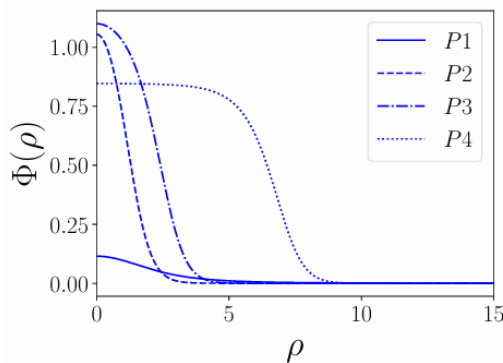
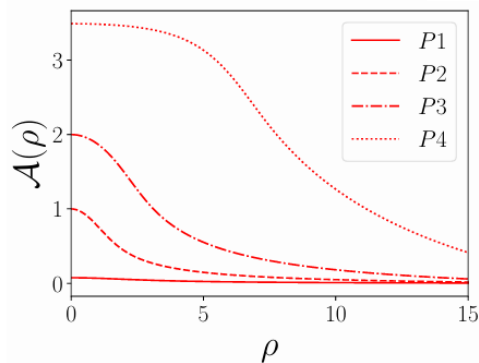
$$\alpha \equiv \frac{|\tilde{g}|}{\sqrt{2\lambda_h}}, \quad k \equiv \frac{\sqrt{\lambda_{\phi h}}}{2\sqrt{\lambda_h}} = \frac{m_\phi}{m_h}$$

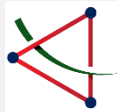
$$\frac{1}{\rho^2} \partial_\rho (\rho^2 \partial_\rho \Phi) + [(\nu - \alpha^2 \mathcal{A})^2 - k^2 \mathcal{H}^2] \Phi = 0,$$

$$\nu \equiv \frac{\omega}{\sqrt{2\lambda_h} v_0}$$

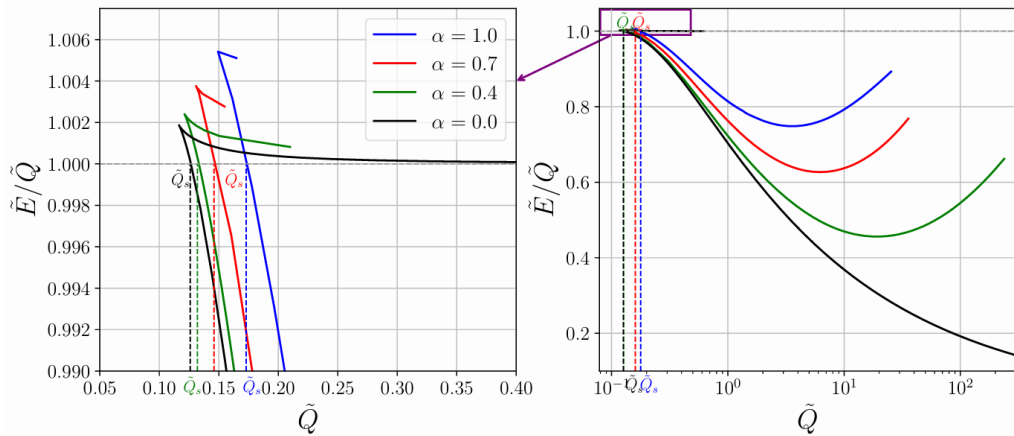
$$\frac{1}{\rho^2} \partial_\rho (\rho^2 \partial_\rho \mathcal{H}) - k^2 \mathcal{H} \Phi^2 - \frac{1}{2} \mathcal{H} (\mathcal{H}^2 - 1) = 0.$$

relaxation method





# Gauged Q-ball stability

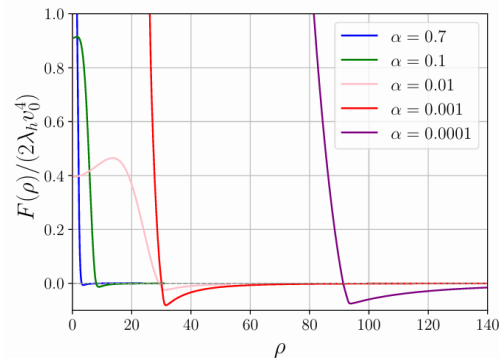
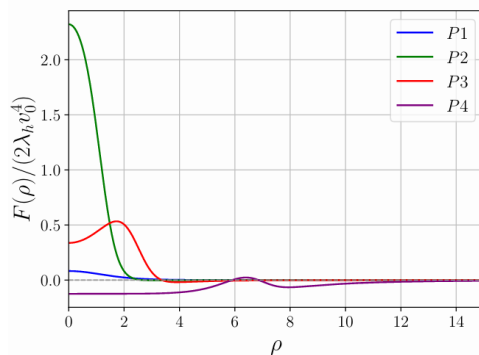


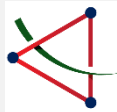
Quantum stability

$$E < m_\phi Q \quad \text{or} \quad \tilde{E}/\tilde{Q} < 1$$

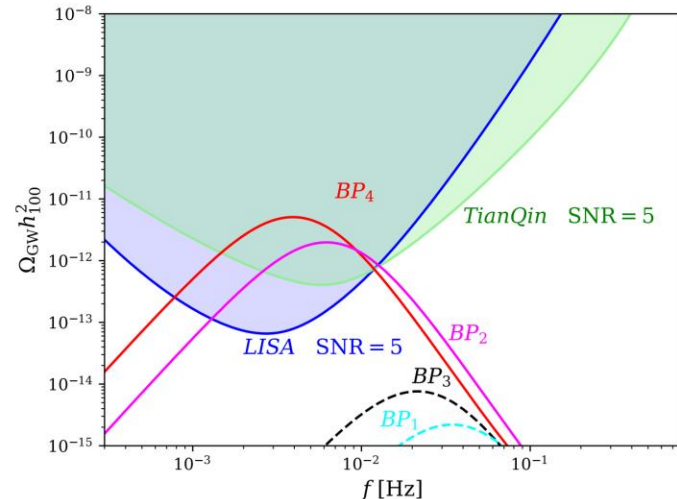
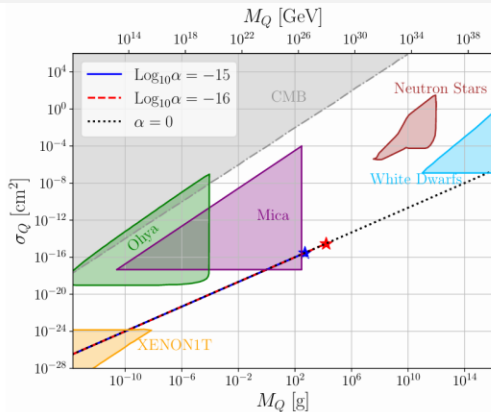
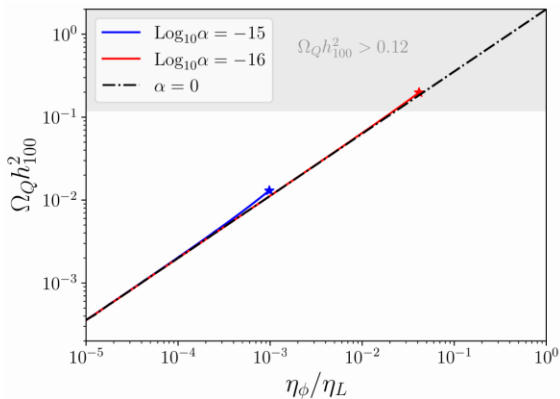
Stress stability

$$F(r) = \frac{2}{3}s(r) + p(r) > 0$$





# Gauged Q-ball DM at TianQin



$$\Omega_Q h_{100}^2$$

$$\simeq 2.81 \times \left( \frac{s_0 h_{100}^2}{\rho_c} \right) \left( \frac{\Gamma(T_*)}{v_w} \right)^{3/16} s_*^{-1/4} (F_\phi^{\text{trap}} \eta_\phi)^{3/4} \lambda_h^{1/4} v_0 \left( 1 + \frac{108^{1/4} \tilde{g}^2 F_\phi^{\text{trap}} \eta_\phi s_* v_w^{3/4}}{5.4 \pi^{7/4} \Gamma(T_*)^{3/4}} \right)$$

$F_\phi^{\text{trap}}$ : The fraction of particles trapped into the false vacuum. It is determined by the phase transition dynamics.

|        | $\lambda_{\phi h}$ | $T_p$ [GeV] | $\alpha_p$ | $\beta/H_p$ | $v_w$ | $F_\phi^{\text{trap}}$ | $\eta_\phi/\eta_L$ | $\delta\sigma_{Zh}$ | GW |
|--------|--------------------|-------------|------------|-------------|-------|------------------------|--------------------|---------------------|----|
| $BP_1$ | 6.8                | 69.8        | 0.12       | 540         | 0.1   | 0.932                  | 0.48               | -0.36%              | ●  |
| $BP_2$ | 6.8                | 70.4        | 0.12       | 578         | 0.6   | 0.805                  | 3.0                | -0.36%              | ●  |
| $BP_3$ | 7.0                | 63.0        | 0.15       | 372         | 0.1   | 0.965                  | 3.4                | -0.37%              | ●  |
| $BP_4$ | 7.0                | 63.9        | 0.15       | 403         | 0.6   | 0.858                  | 20.8               | -0.37%              | ●  |

Siyu Jiang, **FPH**,  
Pyungwon Ko, JHEP 07 (2024) 053

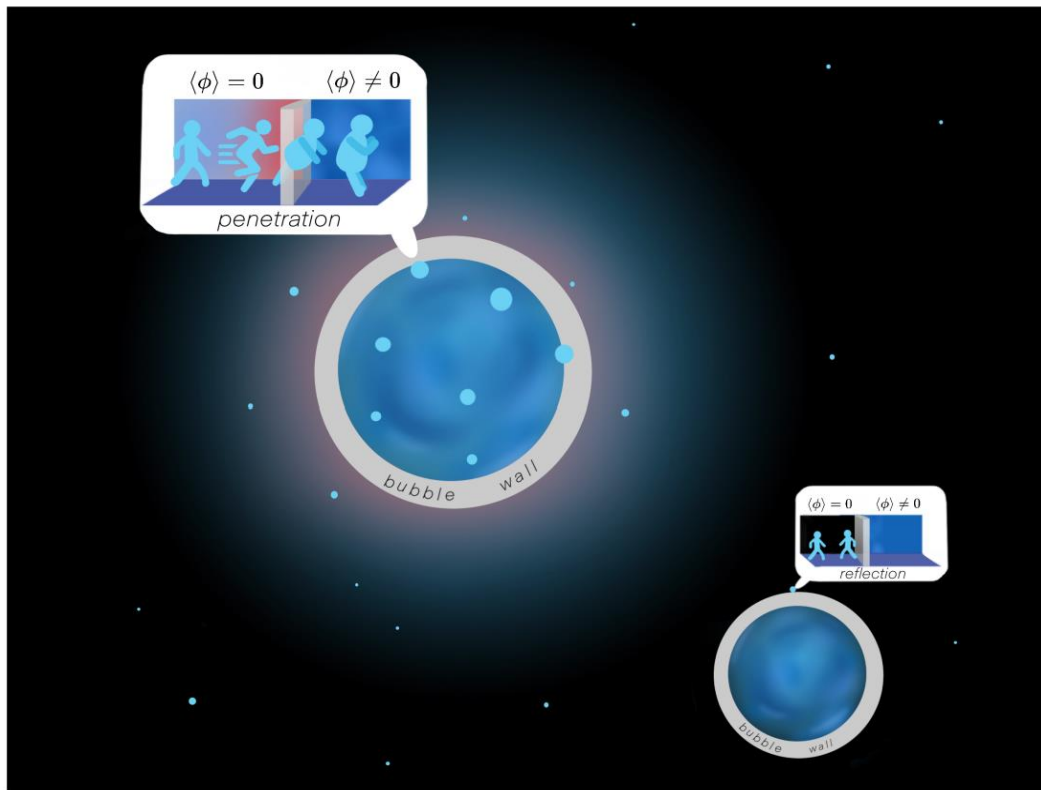


# Case II: filtered DM from phase transition

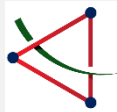


Bubble wall  
dynamics  
plays an essential  
role in the filtered  
DM mechanism.

**DM**



Siyu Jiang, **FPH**, Chong Sheng Li,  
Phys.Rev.D 108 (2023) 6, 063508

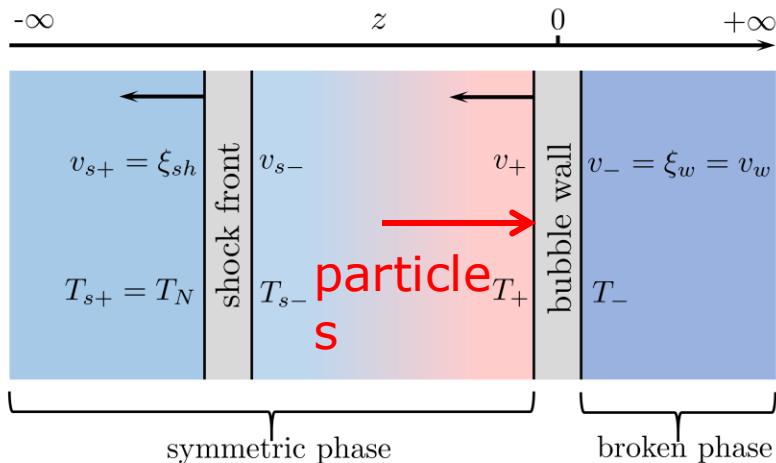


# Case II: filtered DM

Previous work:

$$\tilde{v}_{\text{pl}} = v_w, \quad T = T' = T_n$$

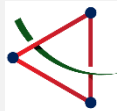
Phys.Rev.Lett. 125 (2020)  
15, 151102, M. J. Baker, J.  
Kopp, and A. J. Long



$$\tilde{v}_{\text{pl}} = \tilde{v}_+, \quad T = T_+, \quad T' = T_- \quad (\text{this work with hydrodynamic effects}).$$

$$J_w^{\text{in}} = \frac{g_\chi}{(2\pi)^2} \int_0^{-1} d\cos\theta \cos\theta \int_{-\frac{m_\chi^{\text{in}}}{\cos\theta}}^\infty dp \frac{p^2}{e^{\tilde{\gamma}_+(1+\tilde{v}_+\cos\theta)p/T_+}} = \frac{g_\chi T_+^3 (1 + \tilde{\gamma}_+ m_\chi^{\text{in}} (1 - \tilde{v}_+)/T_+)}{4\pi^2 \tilde{\gamma}_+^3 (1 - \tilde{v}_+)^2} e^{-\tilde{\gamma}_+ m_\chi^{\text{in}} (1 - \tilde{v}_+)/T_+}.$$

$$n_\chi^{\text{in}} = \frac{J_w^{\text{in}}}{\gamma_w v_w} \quad \Omega_{\text{DM}}^{(\text{hy})} h^2 = \frac{m_\chi^{\text{in}} (n_\chi^{\text{in}} + n_{\bar{\chi}}^{\text{in}})}{\rho_c/h^2} \frac{g_{\star 0} T_0^3}{g_\star (T_-) T_-^3} \simeq 6.29 \times 10^8 \frac{m_\chi^{\text{in}}}{\text{GeV}} \frac{(n_\chi^{\text{in}} + n_{\bar{\chi}}^{\text{in}})}{g_\star (T_-) T_-^3}$$



# Case II: filtered DM

$$T_{\phi}^{\mu\nu} = \partial^{\mu}\phi\partial^{\nu}\phi - g^{\mu\nu}\left[\frac{1}{2}(\partial\phi)^2 - V_{T=0}(\phi)\right] \quad \text{Energy-momentum tensor of scalar field}$$

$$T_{\text{pl}}^{\mu\nu} = \sum_i \int \frac{d^3k}{(2\pi)^3 E_i} k^{\mu} k^{\nu} f_i^{\text{eq}}(k) \quad \text{Energy-momentum tensor of fluid}$$

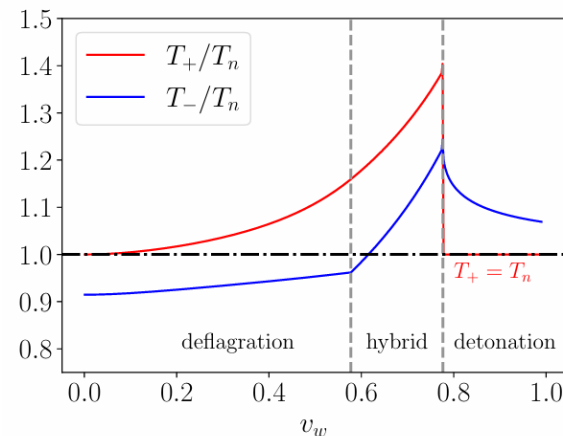
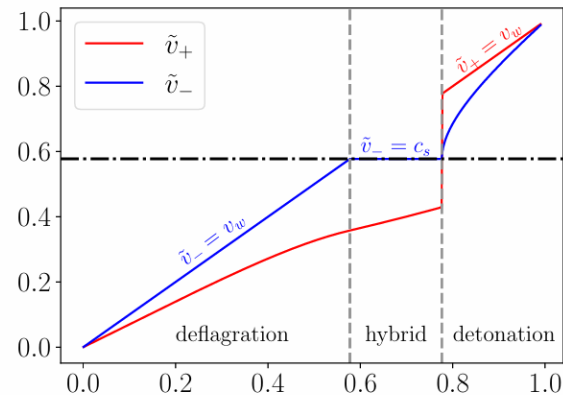
$$T_{\text{fl}}^{\mu\nu} = T_{\phi}^{\mu\nu} + T_{\text{pl}}^{\mu\nu} = \omega u^{\mu} u^{\nu} - p g^{\mu\nu} \quad \text{Energy-momentum conservation}$$

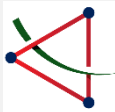
$$\omega_+ \tilde{v}_+^2 \tilde{\gamma}_+^2 + p_+ = \omega_- \tilde{v}_-^2 \tilde{\gamma}_-^2 + p_-, \quad \omega_+ \tilde{v}_+ \tilde{\gamma}_+^2 = \omega_- \tilde{v}_- \tilde{\gamma}_-^2$$

$$\alpha_+ \equiv \epsilon / (a_+ T_+^4)$$

$$r_{\omega} = \omega_+ / \omega_- = (a_+ T_+^4) / (a_- T_-^4)$$

$$\nabla_{\mu} T^{\mu\nu} = 0 \quad \longleftrightarrow \quad \begin{aligned} j_{\xi}^v &= \gamma^2 (1 - v\xi) \left[ \frac{\mu^2}{c_s^2} - 1 \right] \partial_{\xi} v \\ \frac{\partial_{\xi} \omega}{\omega} &= \left( 1 + \frac{1}{c_s^2} \right) \gamma^2 \mu \partial_{\xi} v . \end{aligned}$$





# Case II: filtered DM

Boltzmann equation

$$\mathbf{L}[f_\chi] = \mathbf{C}[f_\chi]$$

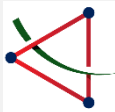
$$f_\chi = \mathcal{A}(z, p_z) f_{\chi,+}^{\text{eq}} = \mathcal{A}(z, p_z) \exp\left(-\frac{\tilde{\gamma}_+(E - \tilde{v}_+ p_z)}{T_+}\right)$$

$$\mathbf{L}[f_\chi] = \frac{p_z}{E} \frac{\partial f_\chi}{\partial z} - \frac{m_\chi}{E} \frac{\partial m_\chi}{\partial z} \frac{\partial f_\chi}{\partial p_z} \quad m_\chi(z) \equiv \frac{m_\chi^{\text{in}}(\phi_-)}{2} \left(1 + \tanh \frac{2z}{L_w}\right)$$

$$g_\chi \int \frac{dp_x dp_y}{(2\pi)^2} \mathbf{L}[f_\chi] \approx \left[ \left( \frac{p_z}{m_\chi} \frac{\partial}{\partial z} - \left( \frac{\partial m_\chi}{\partial z} \right) \frac{\partial}{\partial p_z} - \left( \frac{\partial m_\chi}{\partial z} \right) \frac{\tilde{\gamma}_+ \tilde{v}_+}{T_+} \right) \mathcal{A}(z, p_z) \right] \frac{g_\chi m_\chi T_+}{2\pi \tilde{\gamma}_+} e^{\tilde{\gamma}_+ (\tilde{v}_+ p_z - \sqrt{m_\chi^2 + p_z^2})/T_+}$$

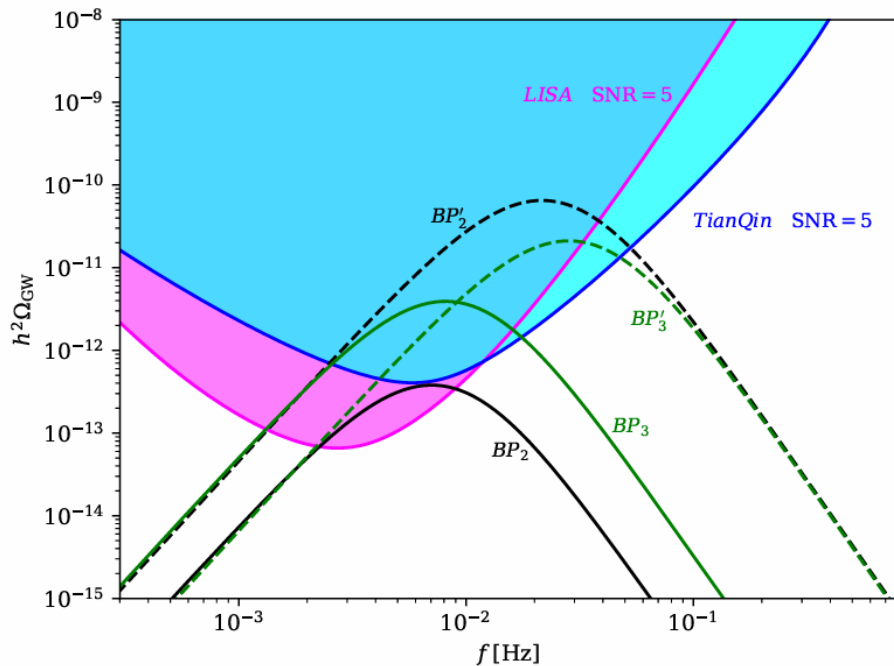
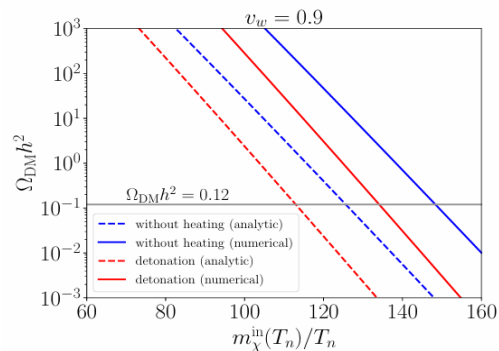
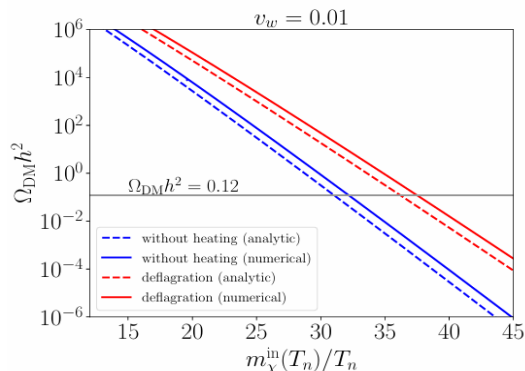
including  $\chi\bar{\chi} \leftrightarrow \phi\phi, \chi\phi \leftrightarrow \chi\phi, \chi\chi \leftrightarrow \chi\chi, \chi\bar{\chi} \leftrightarrow \chi\bar{\chi}, \dots$

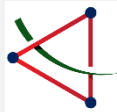
$$\begin{aligned} g_\chi \int \frac{dp_x dp_y}{(2\pi)^2} \mathbf{C}[f_\chi] &= -g_\chi g_{\bar{\chi}} \int \frac{dp_x dp_y}{(2\pi)^2 2E_p^{\mathcal{P}}} d\Pi_{q^{\mathcal{P}}} 4F \sigma_{\chi\bar{\chi} \rightarrow \phi\phi} \left[ f_{\chi_p} f_{\bar{\chi}_q,+}^{\text{eq}} - f_{\chi_p}^{\text{eq}} f_{\bar{\chi}_q}^{\text{eq}} \right] \\ &= -g_\chi g_{\bar{\chi}} \int \frac{dp_x dp_y}{(2\pi)^2 2E_p^{\mathcal{P}}} d\Pi_{q^{\mathcal{P}}} 4F \sigma_{\chi\bar{\chi} \rightarrow \phi\phi} \left[ \mathcal{A} f_{\chi_p,+}^{\text{eq}} f_{\bar{\chi}_q,+}^{\text{eq}} - f_{\chi_p}^{\text{eq}} f_{\bar{\chi}_q}^{\text{eq}} \right] \\ &\equiv \Gamma_{\text{P}}(z, p_z) \mathcal{A}(z, p_z) - \Gamma_{\text{I}}(z, p_z) , \end{aligned}$$



# Case II: filtered DM at TianQin

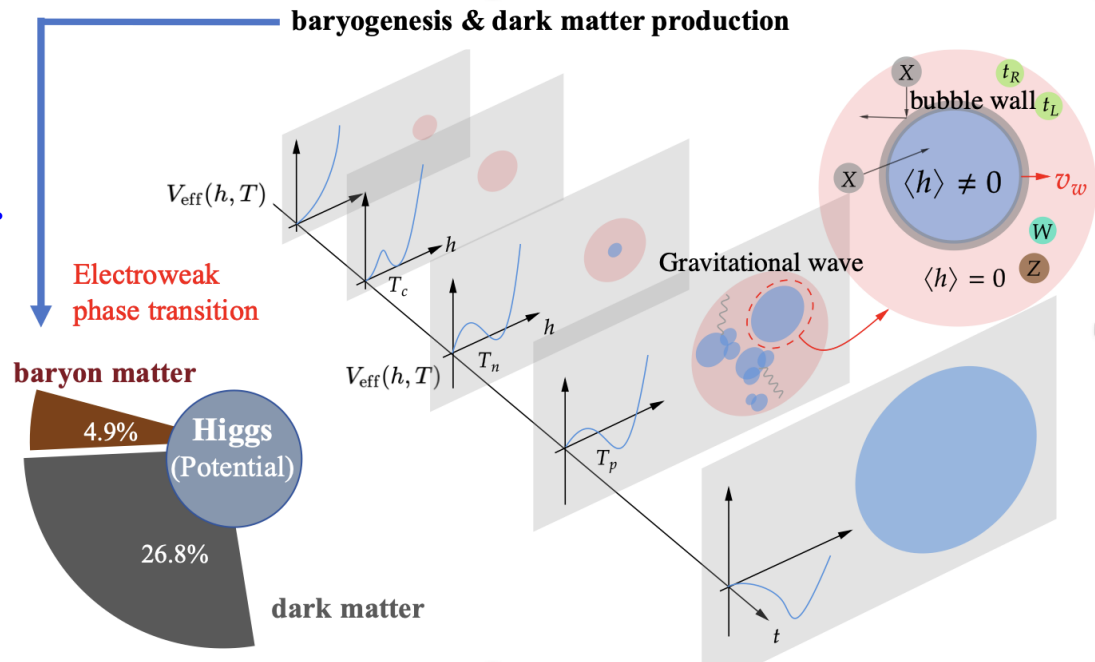
$$n_{\chi}^{\text{in}} = \frac{T_+}{\gamma_w \tilde{\gamma}_+} \int_0^\infty \frac{dp_z}{(2\pi)^2} \mathcal{A}(z \gg L_w, p_z) \exp \left[ \tilde{\gamma}_+ \left( \tilde{v}_+ p_z - \sqrt{p_z^2 + (m_{\chi}^{\text{in}})^2} \right) / T_+ \right] \left( \sqrt{p_z^2 + (m_{\chi}^{\text{in}})^2} + \frac{T_+}{\tilde{\gamma}_+} \right)$$

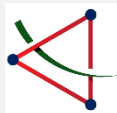




# Summary and outlook

- Cosmic phase transition can naturally explain DM and baryogenesis.
- The associated GW at TianQin provides new approaches to explore DM and cosmic symmetry breaking.



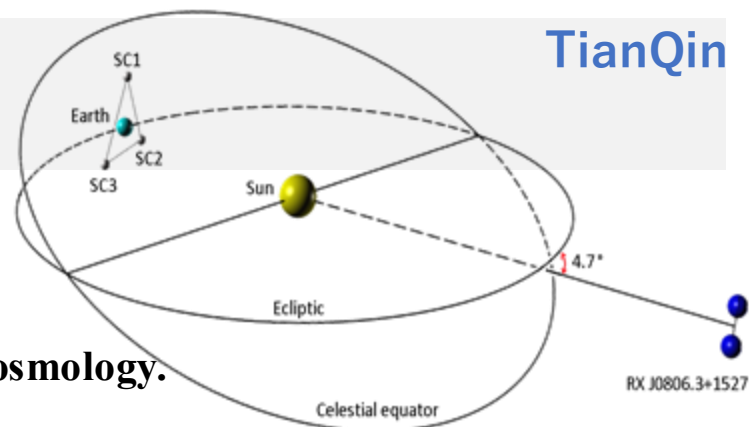


# Summary and outlook

Fundamental Physics and Cosmology with TianQin, arXiv: [2502.20138](#)

Besides direct detection of astrophysics GW sources  
and test of gravity theory,

**TianQin can explore many fundamental problems in particle cosmology.**



TianQin

Correlation and complementay test with CEPC  
Higgs physics (potential), Z-pole physics, TeV new physics

[Progress of the TianQin project](#), arXiv: [2502.11328](#)

Baryon asymmetry of our  
Universe (EW baryogenesis)

GW  
@TianQin

cosmic symmetry  
breaking relics

New DM formation mechanism;  
New approach to explore DM

Thanks for  
your attention !  
Comments and  
collaborations are  
welcome.

[huangfp8@sysu.edu.cn](mailto:huangfp8@sysu.edu.cn)