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New detector configuration for next-generation accelerator-based long-baseline neutrino experiments

Ankur Nathª, Cao Van Son^b, Jennifer Thomas^c, Quyen Phan To^b

[∗]Namrup College, Assam, India, ZIP 786623 b Institute For Interdisciplinary Research in Science and Education (IFIRSE), ICISE, Vietnam ^c University College London, London

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Neutrino Oscillations

Talks on July 8 by Prof. J.W.F. Valle and on July 10 by Prof. S.K. Agarwalla, PASCOS 2024

$$
P(\nu_{\alpha} \to \nu_{\beta}) = \begin{vmatrix} \sum_{j} U_{\alpha j}^{*} U_{\beta j} e^{-\frac{m_{\tau}^{2}}{2E}L} \end{vmatrix}^2 \text{ such that } U_{ij} = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta_{CP}} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta_{CP}} & c_{13}c_{23} \end{bmatrix}
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where, $c_{ij}=\cos\theta_{ij}$, $s_{ij}=\sin\theta_{ij}$ (for $i,j=1,2,3$ and $\alpha,\beta=e,\mu,\tau)^\dagger$.

Table: Global Fit of the oscillation parameters, assuming normal ordering (NO) (*ν*Fit 5.3).

∙arameter	sin ² H_{12} --	$-$ sınf $H_{\rm max}$	sın" けっっ --	\sim $-$ n .,	$(10^{-7}$ eV^2 \rightarrow m $^{\sim}$	д \cdot \cdot -10 - 7 'eV" ∸ $\overline{}$ -
Best tıt	0.307	2.201	$- -$ U.JIL	197	Δ .	4.JII

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Unsolved Problems in Neutrino Oscillation Physics

Leptonic CP Violation

• We ask whether sin $\delta_{CP} = 0$ or not?

Neutrino Mass Hierarchy (MH)

• Whether neutrino MH follows normal ordering (NO) *i.e.* $m_1 < m_2 < m_3$ or inverted ordering (IO) *i.e.* $m_3 < m_1 < m_2$ is still a question.

Mixing angle $θ_{23}$

• Whether *^θ*²³ follows maximal mixing i.e. ⁴⁵◦, or prefers lower octant (LO, *^θ*²³ *<* ⁴⁵◦) to higher octant (HO,

 θ_{23} > 45[°]) is of interest to pursue. Talk by Quyen Phan To, July 9, PASCOS 2024

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MOTIVATION

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016) 6/14

Conventional approach:

- Uses near-site detector (ND) $(<1km)$ to put constraint on the unoscillated neutrino energy spectra.
- Predicts the far-detector (FD) spectra for a specific set of oscillation parameters.
- The use of an ND typically reduces systematics due to $-flux \times cross-section)$ convolution from (10 \sim 15)% to \sim 5% on the event rate.

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New Set-up:

- **Uses no Near Detector**. Instead a 2nd FD at different baseline of same detector volume as that of the 1st FD is considered.
- The combined fiducial mass is kept intact as that of the single FD as in the conventional approach.
- The role of 2nd detector is not to directly constraint the un-oscillated spectra but to measure the oscillated spectra at a different oscillation length.

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With the proposed set-up, relatively large statistics with high power beam (eg. ESS*ν***SB 5MW, T2HK and DUNE** ≈ **MW) can be achieved. Even with large uncertainty in the flux, the correlation in the flux, cross-section and detector response may constrain themselves via the multi-detector data fit.**

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016) 9/14

Experiments' Specifications: ESS*ν***SB**^a

 \mathbf{GLoBES}^{b} is used to simulate the statistical significance of the experiment.

Figure: $P_{\mu e}$ ($P_{\mu \bar{e}}$) vs E for NO (solid) and IO (dashed) for $\delta_{CP} = 0$ (blue) and 270° (red) values is shown in the *left (right)* plot. The fluxes are given by black solid curves.

^a Eur.Phys.J.ST 231 (2022) 21, 3779-3955. ^b Comput. Phys. Commun. 177 (2007) 432. **N.B.** We acknowledge Dr. M. Ghosh of ESS*ν*SB Collab. for sharing the GLoBES-AEDL file.

 ${\sf Figure:}~~ {\sf P}_{\nu_{\,\,\mu}\rightarrow\nu_{\,\rm e}}~({\sf P}_{\nu_{\,\,\mu}^-\rightarrow\nu_{\,\rm e}})$ as a function of baseline (in km).

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- Oscillation maxima is observed for the case

$$
\frac{\Delta m_{32}^2 L}{4E_n} \sim (2n-1) \times 500 \frac{km}{GeV}
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where E_n (GeV) is the energy of the nth oscillation peak/dip for a fixed baseline L.

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- At fixed $E=0.24$ GeV: 2nd oscillation maxima for $L=360$ km and 1st oscillation maxima for L=120 km are observed.
- Second FD placed at relatively lower baseline means higher statistics.

Appearance Event Spectra, ESSνSB Flux, L = 120 km (269 kton)

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Figure: Left: Event Spectra for ν_e (solid) and $\bar{\nu}_e$ (dashed) events for $\delta_{CP} = 0$ (red) and 270° (black) for $L = 120$ km, using $ESSvSB$ flux profiles. **Right:** Mass Hierarchy sensitivity as a function of true δ_{CP} .

CPV sensitivity plots

Top left: Array of 2 FDs with $L2 > 360$ km and 10% errors on signal event rates. Top right: Array of 2 FDs with $L2 > 360$ km and 5% errors on signal event rates. Bottom left: Array of 2 FDs with $L2 \le 360$ km and 10% errors on signal event rates. Bottom rent. Array of 2 FDs with $L2 \le 360$ km and 5% errors on signal event rates. $18/14$

Figure: Uncertanity in 1σ precision measurement of δ _{CP}, considering a single detector FD of fiducial mass of 538kton with a ND **(black solid line)**, compared to that of the cases without a ND with systematic errors of 10% for different FD baseline combinations.

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(i) 540 km (red solid curve), (ii) 120 km (pink), (iii) 60 km (orange), (iv) 40km (green), and (v) 20 km (cyan).

Table: Fraction of true δ _{CP} values for which the proposed set-up **without a ND** gives better 1σ precision of δ_{CP} .

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Figure: Allowed region of sin² θ_{23} at a 3 σ C.L. for: (i) ESS*ν*SB with a ND (black solid line) (ii) 2nd FD placed at 120 km and without a ND (**blue line**). (iii) 2nd FD placed at 60 km and without a ND (**orange line**).

Discussion & Future Scope

- We investigate the new detector configuration with no near-site detector but multiple far-detectors placed at different baselines, and it looks promising.
- 2nd far-detector at 120 km and 60 km (**shorter baseline**) give better precision in the measurement of δ_{CP} for 1*σ* C.L, for 60% of the true values. Preliminary result shows that θ_{23} can be measured with better precision if the 2nd detector is placed at 120 km.
- Next, we shall quantitatively estmiate the precision on θ_{23} and θ_{13} and \mathcal{J} .
	- Precision measurements of δ_{CP} , sin² θ_{23} and sin² $2\theta_{13}$ for discrimination of lepton flavor models.
	- The magnitude of leptonic CPV violation doesn't only depend on δ_{CP} , but is given by the parameterization independent parameter, Jarlskog invariant J :

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\mathcal{J}=\sin\theta_{12}\cos\theta_{12}\sin\theta_{23}\cos\theta_{23}\sin\theta_{13}\cos^2\theta_{13}\sin\delta_{CP}
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. Thus, the precision measurement of J also forms the future scope of this work.

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OPEN TO QUESTIONS & FEEDBACK...