### PASCOS 2024

# New detector configuration for next-generation accelerator-based long-baseline neutrino experiments

Ankur Nath<sup>a</sup>, Cao Van Son<sup>b</sup>, Jennifer Thomas<sup>c</sup>, Quyen Phan To<sup>b</sup>

\*Namrup College, Assam, India, ZIP 786623 <sup>b</sup>Institute For Interdisciplinary Research in Science and Education (IFIRSE), ICISE, Vietnam <sup>c</sup> University College London, London



29TH INTERNATIONAL SYMPOSIUM ON PARTICLES, STRINGS AND COSMOLOGY organised by International Centre for Interdisciplinary Science Education (ICISE), Quy Nhon, Vietnam  $10^{th}$  June, 2024

### Outline

### Section 1 Introduction

Section 2 Motivation and Approach

Section 3 Experiment Specifications

### Section 4

Results

### Section 5

Discussion and Future Scope

### **Neutrino Oscillations**

Talks on July 8 by Prof. J.W.F. Valle and on July 10 by Prof. S.K. Agarwalla, PASCOS 2024

$$P(\nu_{\alpha} \to \nu_{\beta}) = \left| \sum_{j} U_{\alpha j}^{*} U_{\beta j} e^{-\frac{m_{f}^{2}}{2E}L} \right|^{2} \text{ such that } U_{ij} = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta_{CP}} & -c_{12}s_{13}c_{23}e^{i\delta_{CP}} & c_{13}c_{23} \end{bmatrix}$$

where,  $c_{ij} = \cos \theta_{ij}$ ,  $s_{ij} = \sin \theta_{ij}$  (for i, j = 1, 2, 3 and  $\alpha, \beta = e, \mu, \tau$ )<sup>†</sup>.

Table: Global Fit of the oscillation parameters, assuming normal ordering (NO) ( $\nu$ Fit 5.3).

Parameter	$\sin^2 \theta_{12}$	$\sin^2 \theta_{13}(\times 10^{-2})$	$\sin^2 \theta_{23}$	$\delta_{CP}(^{\circ})$	$\Delta m_{21}^2 (10^{-5} \text{eV}^2/c^4)$	$\Delta m_{31}^2 (10^{-3} \text{eV}^2/c^4)$
Best fit	0.307	2.201	0.572	197	7.41	2.511

<sup>&</sup>lt;sup>†</sup>Prog. Theor. Phys. 28, 870 (1962); Zhur. Eksptl'. i Teoret. Fiz. 34 (1958)

### Neutrino Oscillations

Talks on July 8 by Prof. J.W.F. Valle and on July 10 by Prof. S.K. Agarwalla, PASCOS 2024

$$P(\nu_{\alpha} \to \nu_{\beta}) = \left| \sum_{j} U_{\alpha j}^{*} U_{\beta j} e^{-\frac{m_{f}^{2}}{2E}L} \right|^{2} \text{ such that } U_{ij} = \begin{bmatrix} c_{12}c_{13} & s_{13}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta_{CP}} & -c_{12}s_{13}c_{23}e^{i\delta_{CP}} & c_{13}c_{23} \end{bmatrix}$$

where,  $c_{ij} = \cos \theta_{ij}$ ,  $s_{ij} = \sin \theta_{ij}$  (for i, j = 1, 2, 3 and  $\alpha, \beta = e, \mu, \tau$ )<sup>†</sup>.

Table: Global Fit of the oscillation parameters, assuming normal ordering (NO) ( $\nu$ Fit 5.3).

Parameter	$\sin^2 \theta_{12}$	$\sin^2 \theta_{13}(\times 10^{-2})$	$\sin^2 \theta_{23}$	$\delta_{CP}(^{\circ})$	$\Delta m_{21}^2 (10^{-5} \text{eV}^2/c^4)$	$\Delta m_{31}^2 (10^{-3} \text{eV}^2/c^4)$
Best fit	0.307	2.201	0.572	197	7.41	2.511

### **Unsolved Problems in Neutrino Oscillation Physics**

### Leptonic CP Violation

We ask whether sin δ<sub>CP</sub> = 0 or not?

### Neutrino Mass Hierarchy (MH)

• Whether neutrino MH follows normal ordering (NO) *i.e.*  $m_1 < m_2 < m_3$  or inverted ordering (IO) *i.e.*  $m_3 < m_1 < m_2$  is still a question.

### Mixing angle $\theta_{23}$

• Whether  $\theta_{23}$  follows maximal mixing *i.e.* 45°, or prefers lower octant (LO,  $\theta_{23} < 45^{\circ}$ ) to higher octant (HO,

 $\theta_{23}$  > 45°) is of interest to pursue. Talk by Quyen Phan To, July 9, PASCOS 2024

<sup>†</sup>Prog. Theor. Phys. 28, 870 (1962); Zhur. Eksptl'. i Teoret. Fiz. 34 (1958)

## **MOTIVATION**

$$\mathcal{N}_{\nu_{\beta}} \propto \underbrace{\Phi_{\alpha}(E)}_{\mathsf{f}(E)} \times \underbrace{\frac{1}{L^{2}} P_{(\nu_{\alpha} \rightarrow \nu_{\beta})}(E, L, \rho; \theta_{12}, \theta_{13}, \theta_{23}, \Delta m_{21}^{2}, \Delta m_{31}^{2}, \delta_{CP})}_{Propagation} \times \underbrace{\sigma(E)}_{\mathsf{f}(E)} \times \underbrace{\epsilon}_{\mathsf{Detection}} \times \underbrace{\epsilon}_{\mathsf{Detection}} \times \underbrace{\epsilon}_{\mathsf{Detection}} \times \underbrace{\epsilon}_{\mathsf{f}(E)} \times \underbrace{\epsilon}_{\mathsf{f}(E)}$$

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016)

$$N_{\nu_{\beta}} \propto \Phi_{\alpha}(E) \times \frac{1}{L^{2}} P_{(\nu_{\alpha} \rightarrow \nu_{\beta})}(E, L, \rho; \theta_{12}, \theta_{13}, \theta_{23}, \Delta m_{21}^{2}, \Delta m_{31}^{2}, \delta_{CP}) \times \underbrace{\sigma(E)}_{Detection} \times \underbrace{\epsilon}_{Detection}$$
Propagation
Proposed Set-Up
Proton accelerator
$$\underbrace{r_{u} \text{ beam diverges}}_{Earth} \xrightarrow{r_{u} \text{ detector}}_{Earth} \xrightarrow{r_{u} \text{ detector}}_{Earth}$$

#### **Conventional approach:**

- Uses near-site detector (ND) (<1km) to put constraint on the unoscillated neutrino energy spectra.
- Predicts the far-detector (FD) spectra for a specific set of oscillation parameters.
- The use of an ND typically reduces systematics due to ( $flux \times cross$ -section) convolution from ( $10 \sim 15$ )% to  $\sim 5\%$  on the event rate.

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016)

$$\mathcal{N}_{\nu_{\beta}} \propto \underbrace{\Phi_{\alpha}(E)}_{\mathsf{F}} \times \underbrace{\frac{1}{L^{2}} P_{(\nu_{\alpha} \rightarrow \nu_{\beta})}(E, L, \rho; \theta_{12}, \theta_{13}, \theta_{23}, \Delta m_{21}^{2}, \Delta m_{31}^{2}, \delta_{CP})}_{Propagation} \times \underbrace{\underbrace{\Phi_{C}(E)}_{\mathsf{Detection}} \times \underbrace{\Phi_{C}(E)}_{\mathsf{Detection}} \times \underbrace$$

### **Conventional approach:**

- Uses near-site detector (ND) (<1km) to put constraint on the unoscillated neutrino energy spectra.
- Predicts the far-detector (FD) spectra for a specific set of oscillation parameters.
- The use of an ND typically reduces systematics due to ( $flux \times cross$ -section) convolution from ( $10 \sim 15$ )% to  $\sim 5\%$  on the event rate.

### New Set-up:

- Uses no Near Detector. Instead a 2nd FD at different baseline of same detector volume as that of the 1st FD is considered.
- The combined fiducial mass is kept intact as that of the single FD as in the conventional approach.
- The role of 2nd detector is not to directly constraint the un-oscillated spectra but to measure the oscillated spectra at a different oscillation length.

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016)

$$N_{\nu_{\beta}} \propto \underbrace{\Phi_{\alpha}(E)}_{\text{for }} \times \underbrace{\frac{1}{L^{2}} P_{(\nu_{\alpha} \rightarrow \nu_{\beta})}(E, L, \rho; \theta_{12}, \theta_{13}, \theta_{23}, \Delta m_{21}^{2}, \Delta m_{31}^{2}, \delta_{CP})}_{Propagation} \times \underbrace{\sigma(E)}_{\text{for }} \times \underbrace{\epsilon}_{\text{Detection}}$$

### **Conventional approach:**

- Uses near-site detector (ND) (<1km) to put constraint on the unoscillated neutrino energy spectra.
- Predicts the far-detector (FD) spectra for a specific set of oscillation parameters.
- The use of an ND typically reduces systematics due to ( $flux \times cross$ -section) convolution from ( $10 \sim 15$ )% to  $\sim 5\%$  on the event rate.

### New Set-up:

- Uses no Near Detector. Instead a 2nd FD at different baseline of same detector volume as that of the 1st FD is considered.
- The combined fiducial mass is kept intact as that of the single FD as in the conventional approach.
- The role of 2nd detector is not to directly constraint the un-oscillated spectra but to measure the oscillated spectra at a different oscillation length.

With the proposed set-up, relatively large statistics with high power beam (eg. ESS $\nu$ SB 5MW, T2HK and DUNE  $\approx$  MW) can be achieved. Even with large uncertainty in the flux, the correlation in the flux, cross-section and detector response may constrain themselves via the multi-detector data fit.

Image Courtesy: Annual Review of Nuclear and Particle Science, 66, (2016)

### **Experiments' Specifications: ESS** $\nu$ **SB**<sup>a</sup>

**GLoBES**<sup>*b*</sup> is used to simulate the statistical significance of the experiment.



**Figure:**  $P_{\mu e} (P_{\bar{\mu} e})$  vs *E* for NO (*solid*) and IO (dashed) for  $\delta_{CP} = 0$  (blue) and 270° (red) values is shown in the *left* (*right*) plot. The fluxes are given by black solid curves.

<sup>&</sup>lt;sup>a</sup> Eur.Phys.J.ST 231 (2022) 21, 3779-3955. <sup>b</sup> Comput. Phys. Commun. 177 (2007) 432. **N.B.** We acknowledge Dr. M. Ghosh of ESS*ν*SB Collab. for sharing the GLoBES-AEDL file.







#### In search of the second baseline...



Figure:  $P_{\nu_{\mu} \rightarrow \nu_{e}} (P_{\bar{\nu_{\mu}} \rightarrow \bar{\nu_{e}}})$  as a function of baseline (in km).



#### In search of the second baseline...



- Figure:  $P_{\nu_{\mu} \to \nu_{e}} (P_{\bar{\nu_{\mu}} \to \bar{\nu_{e}}})$  as a function of baseline (in km).
- Oscillation maxima is observed for the case

$$\frac{\Delta m_{32}^2 L}{4E_n} \sim (2n-1) \times 500 \frac{km}{GeV}$$

where  $E_n$  (GeV) is the energy of the nth oscillation peak/dip for a fixed baseline L.



#### In search of the second baseline...



- Figure:  $P_{\nu_{\mu} \to \nu_{e}} (P_{\bar{\nu_{\mu}} \to \bar{\nu_{e}}})$  as a function of baseline (in km).
- Oscillation maxima is observed for the case

$$rac{\Delta m_{32}^2 L}{4E_n} \sim (2n-1) imes 500 rac{km}{GeV}$$

where  ${\it E}_n$  (GeV) is the energy of the nth oscillation peak/dip for a fixed baseline L.

- At fixed E=0.24 GeV; 2nd oscillation maxima for L=360km and 1st oscillation maxima for L=120 km are observed.
- Second FD placed at relatively lower baseline means higher statistics.



Appearance Event Spectra, ESSvSB Flux, L = 120 km (269 kton)



Appearance Event Spectra, ESSvSB Flux, L = 120 km (269 kton)

**Figure:** Left: Event Spectra for  $\nu_e$  (solid) and  $\bar{\nu}_e$  (dashed) events for  $\delta_{CP} = 0$  (red) and  $270^{\circ}$  (black) for L = 120km, using ESS $\nu$ SB flux profiles. **Right:** Mass Hierarchy sensitivity as a function of true  $\delta_{CP}$ .

### **CPV** sensitivity plots

The use of far-site detector typically reduce systematics due to flux×cross section convolution from (10 - 15)% to ~ 5% on the event rate.

In these plots, Systematic errors: 5% represents a set-up of 2FDs with a near detector. Systematic errors: 10% represents a set-up of 2FDs without a near detector.



 $\begin{array}{c} \mbox{Figure: CPV sensitivity plots.} \\ \mbox{Top right: Array of 2 FDs with $L2$ > 360km and 10\% errors on signal event rates.} \\ \mbox{Top right: Array of 2 FDs with $L2$ > 360km and 5\% errors on signal event rates.} \\ \mbox{Bottom left: Array of 2 FDs with $L2$ < 360km and 10\% errors on signal event rates.} \\ \mbox{Bottom right: Array of 2 FDs with $L2$ < 360km and 5\% errors on signal event rates.} \\ \mbox{18} 18/14 \end{array}$ 



Figure: Uncertanity in 1 $\sigma$  precision measurement of  $\delta_{CP}$ , considering a single detector FD of fiducial mass of 538kton with a ND (black solid line), compared to that of the cases without a ND with systematic errors of 10% for different FD baseline combinations.



Figure: Uncertanity in  $1\sigma$  precision measurement of  $\delta_{CP}$ , considering a single detector FD of fiducial mass of 538kton with a ND (black solid line), compared to that of the cases without a ND with systematic errors of 10% for different FD baseline combinations. Uncertanity in  $1\sigma$  precision measurement of  $\delta_{CP}$ , without a ND when the 2nd FD is placed at

(i) 540 km (red solid curve), (ii) 120 km (pink), (iii) 60 km (orange), (iv) 40km (green), and (v) 20km (cyan).



**Table:** Fraction of true  $\delta_{CP}$  values for which the proposed set-up **without a ND** gives better  $1\sigma$  precision of  $\delta_{CP}$ .

FD Array	Coverage		
(in km)	(in %)		
(360,120)	60		
(360, 60)	60		
(360, 40)	56		
(360, 20)	43.8		

Figure: Uncertanity in  $1\sigma$  precision measurement of  $\delta_{CP}$ , considering a single detector FD of fiducial mass of 538kton with a ND (black solid line), compared to that of the cases without a ND with systematic errors of 10% for different FD baseline combinations. Uncertanity in  $1\sigma$  precision measurement of  $\delta_{CP}$ , without a ND when the 2nd FD is placed at

(i) 540 km (red solid curve), (ii) 120 km (pink), (iii) 60 km (orange), (iv) 40km (green), and (v) 20km (cyan).



**Figure:** Allowed region of  $\sin^2 \theta_{23}$  at a  $3\sigma$  C.L. for: (i) ESS $\nu$ SB with a ND (black solid line) (ii) 2nd FD placed at 120 km and without a ND (blue line). (iii) 2nd FD placed at 60 km and without a ND (orange line).

### **Discussion & Future Scope**

- We investigate the new detector configuration with no near-site detector but multiple far-detectors placed at different baselines, and *it looks promising*.
- 2nd far-detector at 120 km and 60 km (shorter baseline) give better precision in the measurement of  $\delta_{CP}$  for  $1\sigma$  C.L, for 60% of the true values. Preliminary result shows that  $\theta_{23}$  can be measured with better precision if the 2nd detector is placed at 120 km.
- Next, we shall quantitatively estmiate the precision on  $\theta_{23}$  and  $\theta_{13}$  and  $\mathcal{J}$ .
  - Precision measurements of  $\delta_{CP}$ ,  $\sin^2 \theta_{23}$  and  $\sin^2 2\theta_{13}$  for discrimination of lepton flavor models.
  - The magnitude of leptonic CPV violation doesn't only depend on δ<sub>CP</sub>, but is given by the parameterization independent parameter, Jarlskog invariant *J*:

$$\mathcal{J} = \sin \theta_{12} \cos \theta_{12} \sin \theta_{23} \cos \theta_{23} \sin \theta_{13} \cos^2 \theta_{13} \sin \delta_{CP}$$

. Thus, the precision measurement of  ${\mathcal J}$  also forms the future scope of this work.

### **Discussion & Future Scope**

- We investigate the new detector configuration with no near-site detector but multiple far-detectors placed at different baselines, and *it looks promising*.
- 2nd far-detector at 120 km and 60 km (shorter baseline) give better precision in the measurement of  $\delta_{CP}$  for  $1\sigma$  C.L, for 60% of the true values. Preliminary result shows that  $\theta_{23}$  can be measured with better precision if the 2nd detector is placed at 120 km.
- Next, we shall quantitatively estmiate the precision on  $\theta_{23}$  and  $\theta_{13}$  and  $\mathcal{J}$ .
  - Precision measurements of  $\delta_{CP},\,\sin^2\theta_{23}$  and  $\sin^22\theta_{13}$  for discrimination of lepton flavor models.
  - The magnitude of leptonic CPV violation doesn't only depend on δ<sub>CP</sub>, but is given by the parameterization independent parameter, Jarlskog invariant *J*:

$$\mathcal{J} = \sin \theta_{12} \cos \theta_{12} \sin \theta_{23} \cos \theta_{23} \sin \theta_{13} \cos^2 \theta_{13} \sin \delta_{CP}$$

. Thus, the precision measurement of  ${\mathcal J}$  also forms the future scope of this work.

### **OPEN TO QUESTIONS & FEEDBACK...**