



# MAGNETOGENESIS FROM AN ANISOTROPIC UNIVERSE

*BASED ON ARXIV:2311.07685 (PHYS.REV.D 109 (2024) 8, 083507)*

---

**Sourav Pal**

with D. Maity and T.Q. Do

10th July, 2024

PASCOS 2024, 7-13 July, ICISE, Quy Nhon, Vietnam

1. Introduction
2. Inflationary magnetogenesis
3. Post inflationary evolution
4. Conclusion

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [[AN, IV, Science 328, 73-75 \(2010\)](#)]

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [AN, IV, Science 328, 73-75 (2010)]
- The strength measured is  $10^{-10} < B_0 < 10^{-22} \text{ G}$ .

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [AN, IV, Science 328, 73-75 (2010)]
- The strength measured is  $10^{-10} < B_0 < 10^{-22}$  G.
- Inflationary magnetogenesis is the most widely accepted theory to explain the existence of large-scale fields. [RS et al. Phys. Rev. D 96, 083511 (2017)]

$$\mathcal{L}_{em} = -\frac{1}{4}l^2(\phi)F_{\mu\nu}F^{\mu\nu}$$

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [AN, IV, *Science* 328, 73-75 (2010)]
- The strength measured is  $10^{-10} < B_0 < 10^{-22}$  G.
- Inflationary magnetogenesis is the most widely accepted theory to explain the existence of large-scale fields. [RS et al. *Phys. Rev. D* 96, 083511 (2017)]
- It suffers from strong coupling and backreaction problems.

## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [AN, IV, *Science* 328, 73-75 (2010)]
- The strength measured is  $10^{-10} < B_0 < 10^{-22}$  G.
- Inflationary magnetogenesis is the most widely accepted theory to explain the existence of large-scale fields. [RS et al. *Phys. Rev. D* 96, 083511 (2017)]
- It suffers from strong coupling and backreaction problems.
- Propose a novel mechanism to produce a large-scale magnetic field without an explicit coupling.



## INTRODUCTION AND MOTIVATION

- Magnetic fields are observed in the universe on different coherence length scales.
- Faraday rotation measurements and Gamma rays from blazars put a bound on large-scale magnetic fields. [AN, IV, *Science* 328, 73-75 (2010)]
- The strength measured is  $10^{-10} < B_0 < 10^{-22}$  G.
- Inflationary magnetogenesis is the most widely accepted theory to explain the existence of large-scale fields. [RS et al. *Phys. Rev. D* 96, 083511 (2017)]
- It suffers from strong coupling and backreaction problems.
- Propose a novel mechanism to produce a large-scale magnetic field without an explicit coupling.
- Anisotropic inflation has been proposed in the literature to explain the cold spots in CMB.

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- Bound on the anisotropic parameter:

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- Bound on the anisotropic parameter:
  1. No backreaction on background evolution of inflaton field.

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- Bound on the anisotropic parameter:
  1. No backreaction on background evolution of inflaton field.
  2. The produced EM field does not violate inflation.

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- Bound on the anisotropic parameter:
  1. No backreaction on background evolution of inflaton field.
  2. The produced EM field does not violate inflation.
  3. Reduces to the conformally flat background in the infinite past.

- We work with a Bianchi-I-type metric

$$ds^2 = a^2(\eta) [-d\eta^2 + b^2(\eta)dx^2 + dy^2 + dz^2] \quad (1)$$

$\eta \rightarrow$  Conformal time ;  $a(\eta) \rightarrow$  Scale factor

$b(\eta) \rightarrow$  Anisotropic parameter

- Bound on the anisotropic parameter:
  1. No backreaction on background evolution of inflaton field.
  2. The produced EM field does not violate inflation.
  3. Reduces to the conformally flat background in the infinite past.
- No explicit coupling of the inflaton field to the gauge field, free action:

$$S = -\frac{1}{4} \int d^4x \sqrt{-g} F_{\mu\nu} F^{\mu\nu}$$

$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$  ,  $A_\mu \rightarrow$  Magnetic vector potential.

## ANISOTROPIC MODEL

- The anisotropic parameter is modeled as,

$$b(\eta) = 1 + \alpha e^{-\left(\frac{\eta}{\eta_m}\right)^2} \quad (2)$$



## ANISOTROPIC MODEL

- The anisotropic parameter is modeled as,

$$b(\eta) = 1 + \alpha e^{-\left(\frac{\eta}{\eta_m}\right)^2} \quad (2)$$

$\alpha \rightarrow$  Strength of the anisotropy ,  $\eta_m \rightarrow$  Duration of the anisotropy

## ANISOTROPIC MODEL

- The anisotropic parameter is modeled as,

$$b(\eta) = 1 + \alpha e^{-\left(\frac{\eta}{\eta_m}\right)^2} \quad (2)$$

$\alpha \rightarrow$  Strength of the anisotropy ,  $\eta_m \rightarrow$  Duration of the anisotropy

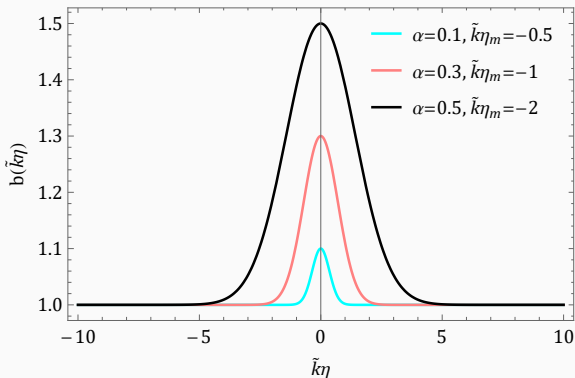


Figure 1: Variation of anisotropic factor  $b$  with time.

- Gauge choice:  $A_0 = 0$ .

## EVOLUTION EQUATIONS

- Gauge choice:  $A_0 = 0$ .
- Constraint equation:

$$\tilde{g}^{im} \partial_i A'_m = 0 ; (\tilde{g}^{ij} = a^2 g^{ij}) \quad (3)$$

# EVOLUTION EQUATIONS

- Gauge choice:  $A_0 = 0$ .
- Constraint equation:

$$\tilde{g}^{im} \partial_i A'_m = 0 ; (\tilde{g}^{ij} = a^2 g^{ij}) \quad (3)$$

- Time evolution of vector potential:

$$A''_n + \frac{b'}{b} A'_n + \tilde{g}_{jn} \tilde{g}'^{jk} A'_k - \tilde{g}^{im} \partial_i F_{mn} = 0 \quad (4)$$

# EVOLUTION EQUATIONS

- Gauge choice:  $A_0 = 0$ .
- Constraint equation:

$$\tilde{g}^{im} \partial_i A'_m = 0 ; (\tilde{g}^{ij} = a^2 g^{ij}) \quad (3)$$

- Time evolution of vector potential:

$$A''_n + \frac{b'}{b} A'_n + \tilde{g}_{jn} \tilde{g}'^{jk} A'_k - \tilde{g}^{im} \partial_i F_{mn} = 0 \quad (4)$$

•

$$\left[ A_i(\tau, \vec{x}), \Pi^j(\tau, \vec{y}) \right] = \frac{i}{\sqrt{-g}} \int \frac{d^3k}{(2\pi)^3} e^{ik_n(x^n - y^n)} \left( \delta_i^j - \frac{k_i k^j}{k_n k^n} \right) \quad (5)$$

# EVOLUTION EQUATIONS

- Gauge choice:  $A_0 = 0$ .
- Constraint equation:

$$\tilde{g}^{im} \partial_i A'_m = 0 ; (\tilde{g}^{ij} = a^2 g^{ij}) \quad (3)$$

- Time evolution of vector potential:

$$A''_n + \frac{b'}{b} A'_n + \tilde{g}_{jn} \tilde{g}'^{jk} A'_k - \tilde{g}^{im} \partial_i F_{mn} = 0 \quad (4)$$

•

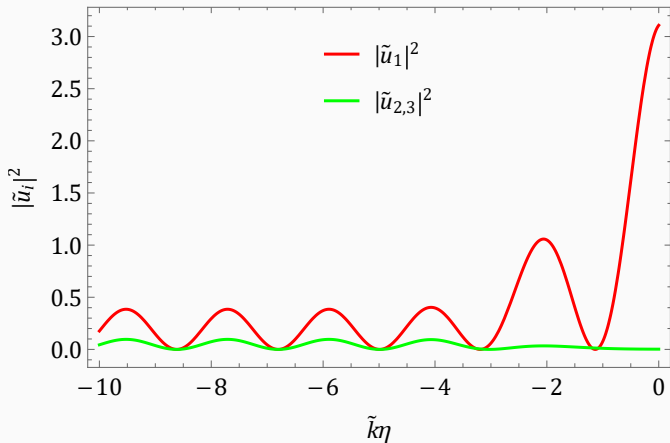
$$[A_i(\tau, \vec{x}), \Pi^j(\tau, \vec{y})] = \frac{i}{\sqrt{-g}} \int \frac{d^3k}{(2\pi)^3} e^{ik_n(x^n - y^n)} \left( \delta_i^j - \frac{k_i k^j}{k_n k^n} \right) \quad (5)$$

- Evolution in the momentum space:

$$u''_n + \frac{b'}{b} u'_n + \tilde{g}'^{jl} \tilde{g}_{jn} u'_l + \tilde{g}^{im} (k_m k_i u_n - k_n k_i u_m) = 0. \quad (6)$$

# MODE FUNCTION EVOLUTION

- Evolution of mode functions:



**Figure 2:** Evolution of mode functions with time for anisotropic free parameters  $\alpha = 3$ ,  $\tilde{k}\eta_m = -2$  ( $\tilde{u}_i = \sqrt{\tilde{k}}u_i$ ).



- Power spectra of the EM field are defined as:

$$\mathcal{P}_{E/B}(k, \eta) = \frac{\partial \rho_{E/B}}{\partial \ln k}$$

- Power spectra of the EM field are defined as:

$$\mathcal{P}_{E/B}(k, \eta) = \frac{\partial \rho_{E/B}}{\partial \ln k}$$

- In terms of mode functions:

$$\mathcal{P}_E(\eta, \tilde{k}) = \frac{\tilde{k}^3}{2\pi^2 a^4} \left( \frac{|u'_1(\eta)|^2}{b^2} + |u'_2(\eta)|^2 + |u'_3(\eta)|^2 \right) \quad (7)$$

$$\mathcal{P}_B(\eta, \tilde{k}) = \frac{\tilde{k}^5}{2\pi^2 a^4} \left[ \frac{1}{b^2} \left( 2|u_1|^2 + |u_2|^2 + |u_3|^2 - 2\Re(u_1 u_2^*) \right) - 2\Re(u_1 u_3^*) \right] + \left( |u_2|^2 + |u_3|^2 - 2\Re(u_2 u_3^*) \right) \quad (8)$$

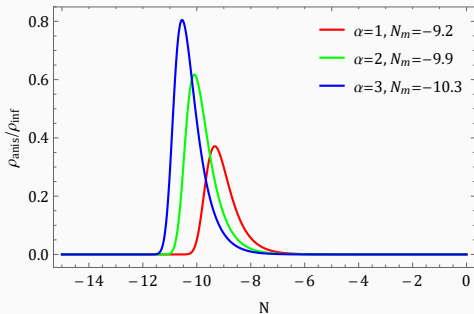
- The total energy of the background:

- The total energy of the background:

$$\rho_{\text{total}} = -T_0^0 = 3H^2 M_{\text{pl}}^2 + 2HM_{\text{pl}}^2 \frac{b'}{ab} \quad (9)$$

- The total energy of the background:

$$\rho_{\text{total}} = -T_0^0 = 3H^2 M_{\text{pl}}^2 + 2HM_{\text{pl}}^2 \frac{b'}{ab} \quad (9)$$



**Figure 3:** Evolution of the ratio of anisotropic energy density to inflationary energy density with e-folding number.

$$N = \ln(a/a_{\text{end}}) ; a = -1/H\eta$$

- We get the upper limit on  $\alpha \leq 1.48$ .  $\left( \frac{\rho_{anis}}{\rho_{inf}} \sim 0.5 \right)$

- We get the upper limit on  $\alpha \leq 1.48$ .  $\left(\frac{\rho_{anis}}{\rho_{inf}} \sim 0.5\right)$
- Checking backreaction from produced EM field:

- We get the upper limit on  $\alpha \leq 1.48$ .  $\left(\frac{\rho_{anis}}{\rho_{inf}} \sim 0.5\right)$
- Checking backreaction from produced EM field:
- For  $\alpha = 1.45$ ,  $\tilde{k}\eta_m = -1$  we have

$$\frac{\rho_E + \rho_B}{\rho_{inf}} \sim 10^{-9};$$



- We get the upper limit on  $\alpha \leq 1.48$ .  $\left(\frac{\rho_{anis}}{\rho_{inf}} \sim 0.5\right)$
- Checking backreaction from produced EM field:
- For  $\alpha = 1.45$ ,  $\tilde{k}\eta_m = -1$  we have

$$\frac{\rho_E + \rho_B}{\rho_{inf}} \sim 10^{-9};$$

- No backreaction from produced EM field during inflation.

- We get the upper limit on  $\alpha \leq 1.48$ .  $\left(\frac{\rho_{anis}}{\rho_{inf}} \sim 0.5\right)$
- Checking backreaction from produced EM field:
- For  $\alpha = 1.45$ ,  $\tilde{k}\eta_m = -1$  we have

$$\frac{\rho_E + \rho_B}{\rho_{inf}} \sim 10^{-9};$$

- No backreaction from produced EM field during inflation.
- Anisotropy does not affect inflation.

## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING

- The universe becomes radiation-dominated instantly after inflation.

## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING

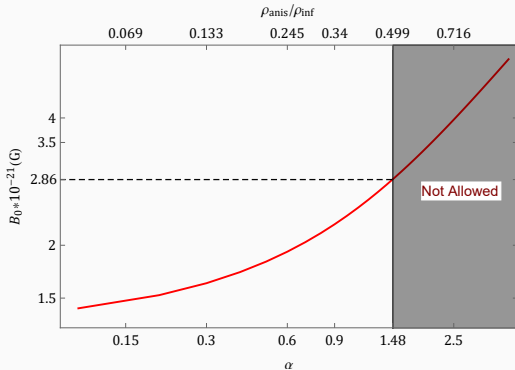
- The universe becomes radiation-dominated instantly after inflation.
- Energy density decays as  $\rho \propto a^{-4}$ .

## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING

- The universe becomes radiation-dominated instantly after inflation.
- Energy density decays as  $\rho \propto a^{-4}$ .
- Magnetic field decays as  $B \propto a^{-2}$ .

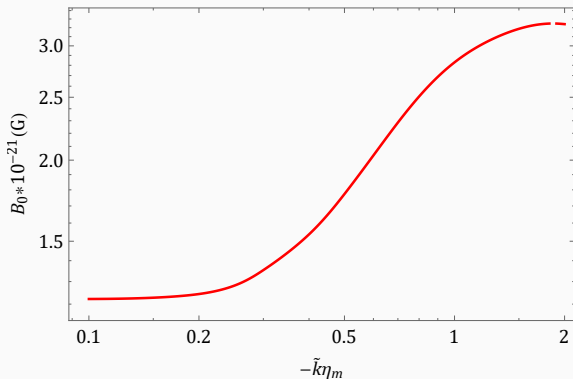
## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING

- The universe becomes radiation-dominated instantly after inflation.
- Energy density decays as  $\rho \propto a^{-4}$ .
- Magnetic field decays as  $B \propto a^{-2}$ .



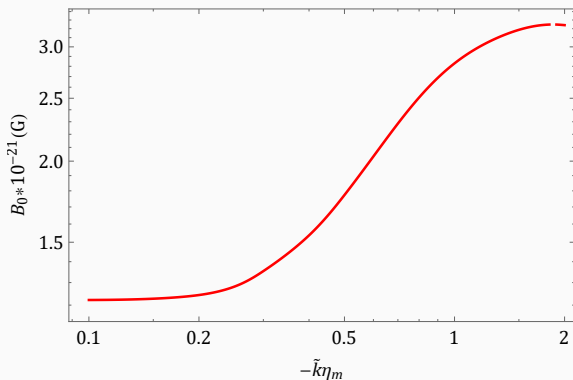
**Figure 4:** Variation of present strength of magnetic field with anisotropic parameter  $\alpha$  for  $\tilde{k}\eta_m = -2$ .

## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING



**Figure 5:** Variation of present strength of magnetic field with  $\tilde{k}\eta_m$  for  $\alpha = 1.45$ .

## POST INFLATIONARY EVOLUTION IN INSTANT REHEATING



**Figure 5:** Variation of present strength of magnetic field with  $\tilde{k}\eta_m$  for  $\alpha = 1.45$ .

- Maximum magnetic field value obtained through instant reheating:  $B_0 \sim 3 \times 10^{-21} \text{ G}$ .



## POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

- The universe goes through an era of negligible conductivity.

## POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

- The universe goes through an era of negligible conductivity.
- The electric field does not vanish instantly.

## POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

- The universe goes through an era of negligible conductivity.
- The electric field does not vanish instantly.
- Leading to Faraday's induction, and the electric field converts into a magnetic field.

## POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

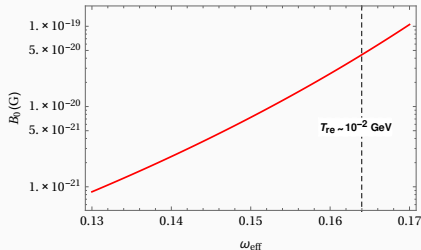
- The universe goes through an era of negligible conductivity.
- The electric field does not vanish instantly.
- Leading to Faraday's induction, and the electric field converts into a magnetic field.
- Magnetic field decays slower, and strength depends on the effective equation of state ( $\omega_{eff}$ ). [ DM, SP, TP, JCAP 05 (2021) 045]

## POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

- The universe goes through an era of negligible conductivity.
- The electric field does not vanish instantly.
- Leading to Faraday's induction, and the electric field converts into a magnetic field.
- Magnetic field decays slower, and strength depends on the effective equation of state ( $\omega_{eff}$ ). [ DM, SP, TP, JCAP 05 (2021) 045]
- Enhance magnetic field strength more than in an instant reheating scenario.

# POST INFLATIONARY EVOLUTION IN PROLONGED REHEATING

- The universe goes through an era of negligible conductivity.
- The electric field does not vanish instantly.
- Leading to Faraday's induction, and the electric field converts into a magnetic field.
- Magnetic field decays slower, and strength depends on the effective equation of state ( $\omega_{eff}$ ). [ DM, SP, TP, JCAP 05 (2021) 045]
- Enhance magnetic field strength more than in an instant reheating scenario.



**Figure 6:** Variation of present-day magnetic field with effective equation of state  $\omega_{eff}$ .

- The anisotropic parameter  $\alpha$  is constrained from back-reaction  
 $0.3 < \alpha < 1.48$

- The anisotropic parameter  $\alpha$  is constrained from back-reaction  
 $0.3 < \alpha < 1.48$
- The produced EM field does not back-react to the inflationary background.









- The anisotropic parameter  $\alpha$  is constrained from back-reaction  
 $0.3 < \alpha < 1.48$
- The produced EM field does not back-react to the inflationary background.
- We can achieve  $B_0 \sim 3 \times 10^{-21} \text{ G}$  with instant reheating.

- The anisotropic parameter  $\alpha$  is constrained from back-reaction  $0.3 < \alpha < 1.48$
- The produced EM field does not back-react to the inflationary background.
- We can achieve  $B_0 \sim 3 \times 10^{-21} \text{ G}$  with instant reheating.
- With the introduction of reheating, we can further strengthen up to  $B_0 \sim 4 \times 10^{-20} \text{ G}$ .

- The anisotropic parameter  $\alpha$  is constrained from back-reaction  $0.3 < \alpha < 1.48$
- The produced EM field does not back-react to the inflationary background.
- We can achieve  $B_0 \sim 3 \times 10^{-21} \text{ G}$  with instant reheating.
- With the introduction of reheating, we can further strengthen up to  $B_0 \sim 4 \times 10^{-20} \text{ G}$ .
- We get a rather tight constraint on  $0.132 < \omega_{eff} < 0.164$ .

## REFERENCES

-  A. Neronov and I. Vovk, Science **328**, 73-75 (2010) [arXiv:1006.3504 [astro-ph.HE]].
-  R. Sharma, S. Jagannathan, T. R. Seshadri and K. Subramanian, Phys. Rev. D **96**, no.8, 083511 (2017) [arXiv:1708.08119 [astro-ph.CO]].
-  C. Pitrou, T. S. Pereira, and J. P. Uzan, JCAP **04**, 004 (2008) [arXiv:0801.3596].
-  T. Kobayashi and M. S. Sloth, Phys. Rev. D **100**, 023524 (2019) [arXiv:1903.02561].
-  L. Dai, M. Kamionkowski, and J. Wang, Phys. Rev. Lett. **113**, 041302 (2014) [arXiv:1404.6704].
-  D. Maity, S. Pal, and T. Paul, JCAP **05**, 045 (2021) [arXiv:2103.02411].

THANK YOU!

*PAL.SOURAV@IITG.AC.IN*