Baryogenesis, Dark Matter, and PTA signal from a Dark Conformal Phase Transition

Sudhakantha Girmohanta

July 10, 2024

Based on:

1. Fujikura, Girmohanta, Nakai and Suzuki [PLB 846, 138203 (2023)]

2. Fujikura, Girmohanta, Nakai and Zhang [arXiv:2406.12956]





I. Introduction

Pulsar Timing Arrays, nano-Hz Gravitational waves, and supermassive black hole binaries (SMBHB)

The spectrum of gravitational waves

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Earth-pulsar system as gravitational wave antenna. Gravitational waves change the arrival time of pulses. Estabrook, Wahlquist '75; Sazhin '77; Detweiler '79



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GW: Distinctive quadrupolar inter-pulsar correlation. Hellings, Downs '82

$$\frac{\Delta \nu_i(t)}{\nu_i} = \alpha_i h(t) + n_i(t),$$

$$\alpha_{ij} \equiv \frac{1}{4\pi} \int \alpha_i \alpha_j \, d\Omega = \frac{1 - \cos \gamma_{ij}}{2} \ln\left(\frac{1 - \cos \gamma_{ij}}{2}\right)$$

$$-\frac{1}{6} \frac{1 - \cos \gamma_{ij}}{2} + \frac{1}{3}, \qquad (5)$$

where γ_{ij} is the angle between the two pulsars.

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Separation Angle Between Pulsars, ξ_{ab} [degrees]



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Possible sources:



Cosmological phase transitions.

Defects: Cosmic strings, domain walls...



EPTA+InPTA





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РРТА



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But...one has to first get to a orbital separation of ~ 0.01 pc.

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Do SMBHBs merge in the lifetime of the Universe?



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II. Phase Transition Interpretation

First-order phase transition in a nearly conformal dark sector and the production of gravitational waves

Fujikura, Girmohanta, Nakai and Suzuki [PLB 846, 138203 (2023)]

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1st order phase transition proceeds via nucleation, expansion and merger of bubbles of the true ground state.

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$$f_0 = f_* \frac{a(\tau_*)}{a(\tau_0)}$$
 $f_0 \simeq 10^{-8} \text{ Hz} \left(\frac{T_*}{1 \text{ GeV}}\right)$

Observed frequency f_0 is redshifted and is associated with the epoch when GW was produced.

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From the underlying field theory to GW spectra

Calculate the bubble nucleation rate from the bounce action $S_{\rm B}$ $\Gamma \sim T^4 e^{-S_{\rm B}}$ Calculable Effective parameters α : Latent heat released β : Bubble nucleation rate v_w : Bubble wall speed

Extensive numerical simulations have given approximate analytical fit for the resulting GW spectrum

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Utilize weakly coupled description of deconfining phase transition via AdS/CFT.

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4D conformal dark sector with large N + dark pure $SU(N_H)$ Yang-Mills

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For eg: dark radiation final state.

Contributes to ΔN_{eff} and may alleviate the Hubble tension.

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Rattazzi+ 2002 ; Servant+ 2017 Dual 5D description

IR brane replaced by an event horizon + $SU(N_H)$ in the bulk

Dark $SU(N_H)$ confinement generates radion potential

Below $T_c^{(D)}$: IR brane configuration has lower free-energy

IR brane bubbles appear and a strong first-order phase transition proceeds.

 $SU(N_H)$ confinement generates $V_{\rm eff}(\varphi)$

$$\frac{1}{g_{\rm H}^2(Q,\varphi)} = -\frac{b_{\rm CFT}}{8\pi^2} \ln\left(\frac{k}{\varphi}\right) - \frac{b_{\rm H}}{8\pi^2} \ln\left(\frac{k}{Q}\right)$$

Running of $SU(N_{\rm H})$ coupling $g_{\rm H}$ from UV scale k to $Q \leq \varphi . b_{\rm CFT} = -\xi N$, $b_{\rm H} = 11 N_{\rm H}/3$.

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III. Cold Darkogenesis

Theoretical challenges to the phase transition interpretation and possible resolution: *Dark Matter and Baryon Asymmetry*

Fujikura, Girmohanta, Nakai and Zhang, arXiv:2406.12956

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Generating dark asymmetry
$$V(\varphi, H_{\rm D}) = V_{\rm eff}(\varphi) + \frac{\lambda}{4} \left[H_{\rm D}^{\dagger} H_{\rm D} - \frac{v_{\rm D}^2}{2} \left(\frac{\varphi}{\varphi_{\rm min}} \right)^2 \right]^2.$$

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- The later induces dark lepton number violation via anomaly.

$$\partial_{\mu}j^{\mu}_{\mathrm{D}_{\mathrm{L}}} = N_{\mathrm{D}_{\mathrm{L}}}\frac{g_{\mathrm{D}}^{2}}{32\pi^{2}}\mathrm{Tr}\left(W_{\mathrm{D}}^{\mu\nu}\widetilde{W}_{\mathrm{D},\mu\nu}\right) ; \quad \mathcal{O}_{\mathrm{CPV}} = \delta_{\mathrm{CP}}\frac{H_{\mathrm{D}}^{\dagger}H_{\mathrm{D}}}{\Lambda_{\mathrm{CP}}^{2}}\frac{g_{\mathrm{D}}^{2}}{32\pi^{2}}\mathrm{Tr}\left(W_{\mathrm{D}}^{\mu\nu}\widetilde{W}_{\mathrm{D},\mu\nu}\right)$$

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• With C & CP violation, $\delta N \equiv N_{\text{CS}} - N_{\text{H}} > 0$ and $\delta N < 0$ winding configurations evolve differently, generating a net dark lepton number $\mathcal{D}_{\text{L,in}}$.

• Generated dark asymmetry $\mathcal{D}_{L,in}$ is stored in L_{χ} , χ :

$$\mathscr{D}_{\rm L,in} \simeq 10^{-10} \left(\frac{N_{D_{\rm L}}}{2}\right) \left(\frac{\delta_{\rm CP}}{10^{-4}} \frac{\varphi_{\rm min}^2}{\Lambda_{\rm CP}^2}\right) \left(\frac{\alpha_{\rm D}}{1.5 \times 10^{-2}}\right)^4 \left(\frac{\lambda}{10^{-4}}\right)^{-3/2} \left(\frac{5v_{\rm D}}{\varphi_{\rm min}}\right)^{1/2} \left(\frac{\varphi_{\rm min}}{2T_{\rm RH}}\right)^3$$

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♦ Asymmetries in the visible sector \mathscr{B}_f and dark baryon sector \mathscr{D}_B can be related:

$$\mathscr{B}_{f} = \left[\frac{2+4N_{D_{L}}}{4N_{D_{L}}+N_{D_{B}}+2}\right] \mathscr{D}_{L,in} \quad ; \quad \mathscr{D}_{B} = \left[\frac{N_{D_{B}}}{4N_{D_{L}}+N_{D_{B}}+2}\right] \mathscr{D}_{L,in} \quad ; \quad m_{p_{D}} \simeq 5 \left|\frac{\mathscr{B}_{f}}{\mathscr{D}_{B}}\right| \text{GeV}$$

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Baryonic DM composed of $f(\mathbb{Z}_2 \text{ odd})$ $\mathcal{O}_{\mathrm{D}} \sim \frac{1}{\Lambda_{\mathrm{D}}^2} p_{\mathrm{D}} p_{\mathrm{D}} \chi \chi$; $\mathcal{O}_n \sim \frac{1}{\Lambda_n^2} \chi u_{\mathrm{R}} d_{\mathrm{R}} d_{\mathrm{R}}$ $\mathcal{O}_n \sim \frac{1}{\Lambda_n^2} \chi u_{\mathrm{R}} d_{\mathrm{R}} d_{\mathrm{R}}$ For equilibrium at GeV $\Lambda_n \lesssim 15$ TeV.

♦ Asymmetries in the visible sector \mathscr{B}_f and dark baryon sector \mathscr{D}_B can be related:

$$\mathscr{B}_{f} = \left[\frac{2+4N_{D_{L}}}{4N_{D_{L}}+N_{D_{B}}+2}\right]\mathscr{D}_{L,\text{in}} \quad ; \quad \mathscr{D}_{B} = \left[\frac{N_{D_{B}}}{4N_{D_{L}}+N_{D_{B}}+2}\right]\mathscr{D}_{L,\text{in}} \quad ; \quad m_{p_{D}} \simeq 5\left|\frac{\mathscr{B}_{f}}{\mathscr{D}_{B}}\right| \text{GeV}$$

★ The DM is self-interacting via the mediation of dark pions π_D with cross-section: $\frac{\sigma_{p_{\rm D}p_{\rm D}}}{m_{p_{\rm D}}} \sim 1 \text{ cm}^2/g \left(\frac{\Lambda_{\rm H,0}}{m_{p_{\rm D}}}\right) \left(\frac{\Lambda_{\rm H,0}}{a_{\rm D}^{-1}}\right)^2 \left(\frac{150 \text{ MeV}}{\Lambda_{\rm H,0}}\right)^3; \quad a_{\rm D}: \text{ scattering length}.$ Tulin Yu (2017); Kribs (2016)

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 $\lambda_h \lesssim 0.1$ from Higgs invisible decay.

Lower bound from BBN, upper bound from DM direct detection.

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PTA signal explanation together with DM and baryon asymmetry



Conclusions

- ✓ <u>Dark first-order phase transition</u> is a promising interpretation of the observed PTA signal.
- ✓ Confining nearly conformal phase transition can realize a supercooled phase transition to explain the data. We analyzed it using the dilaton effective potential.
- ✓ Both secluded dark sector (together with SMBHB) and decaying dark sector can explain the observed signal.
- ✓ The strong supercooling exponentially dilutes away pre-existing baryon asymmetry and DM, posing a challenge to this scenario.
- ✓ We provide a concrete scenario of cold darkogenesis where the baryon asymmetry and DM are produced utilizing the phase transition.
- ✓ Future direct detection searches for DM and mono-jet searches at colliders will probe this model further.

Thank you for your time! Questions?