

Testing the neutrino mass sum rules with present data and future prospects

Quang Anh Nguyen¹ Dr. Son Cao²

¹University of Science and Technology of Hanoi

²Institute for Interdisciplinary Research in Science and Education

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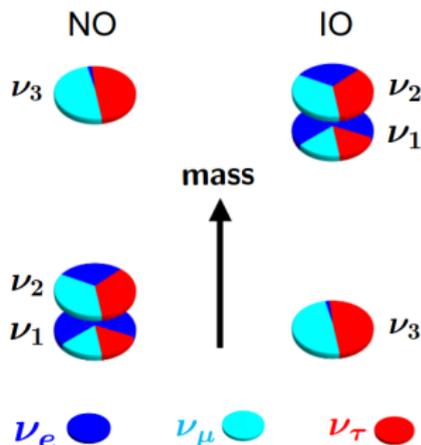
Neutrino Mass Ordering

- Neutrinos have non-zero mass.
- According to neutrino oscillation data, the mass squared difference:

$$|\Delta m_{31}^2| = |m_3^2 - m_1^2| \approx 2.5 \times 10^{-3} \text{eV}^2$$

$$\Delta m_{21}^2 = m_2^2 - m_1^2 \approx 7.5 \times 10^{-5} \text{eV}^2$$

- Only the mass scale is known, the exact mass value yet be able to be measured.



PMNS matrix parameterization

- The matrix can be parameterized with three mixing angles $\theta_{13}, \theta_{12}, \theta_{23}$ and one irreducible Dirac CP-violation phase δ_{CP}
- $c_{ij} \equiv \cos \theta_{ij}$ and $s_{ij} \equiv \sin \theta_{ij}$
- If the neutrino is a Majorana particle, there are two extra CP-violation phases, ρ_1 and ρ_2 , which do not affect neutrino oscillations.

$$U = U(\theta_{23})U(\theta_{13}, \delta_{CP})U(\theta_{12})U_M(\alpha_1, \alpha_2)$$
$$= \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\alpha_1} & 0 & 0 \\ 0 & e^{i\alpha_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Constraints on neutrino mass

- From cosmology: arxiv.org/abs/1807.06209
(Results from Planck and Baryon Acoustics Oscillation)

$$\sum_{N_\nu} m_{\nu_i} = m_{\nu_1} + m_{\nu_2} + m_{\nu_3} < 0.12\text{eV at } 90\% \text{ C.L}$$

where $m_{\nu_1}, m_{\nu_2}, m_{\nu_3}$ are three mass eigenstates of neutrino

- Beta decay (KATRIN experiment): arxiv.org/abs/2404.08733v1

$$(A, Z) \rightarrow (A, Z + 1) + e^- + \bar{\nu}_i$$

$$\langle m_{\nu_\beta} \rangle = \sqrt{|U_{e1}|^2 m_{\nu_1}^2 + |U_{e2}|^2 m_{\nu_2}^2 + |U_{e3}|^2 m_{\nu_3}^2}$$
$$\langle m_{\nu_\beta} \rangle < 0.8\text{eV at } 90\% \text{ C.L}$$

where U_{e1}, U_{e2}, U_{e3} are elements of the PMNS matrix

Constraints on neutrino mass

- Neutrinoless double beta decay: arxiv.org/abs/2406.11438v1
(KAMLand-Zen experiment, only if neutrinos are Majorana particles)

$$(A, Z) \rightarrow (A, Z + 2) + 2e$$

$$\begin{aligned} \langle m_{\beta\beta}^{0\nu} \rangle &= \sum_i U_{ei}^2 m_{\nu_i} < 28 - 122 \text{meV at } 90\% \text{ C.L} \\ &= |m_{\nu_1} \cos^2 \theta_{13} \cos^2 \theta_{12} + m_{\nu_2} \cos^2 \theta_{13} \sin^2 \theta_{12} e^{i\rho_1} + m_{\nu_3} \cos^2 \theta_{13} e^{i\rho_2}| \end{aligned}$$

where: θ_{13}, θ_{12} are the mixing angles

ρ_1, ρ_2 are two Majorana phases, from 0 to π

Emergence of ν mass sum rule from flavor symmetries

- Neutrino mass matrix M_ν , in general, depends on six complex parameters
- Unitary transformation diagonalises M_ν and leads to three complex mass eigenvalues, which in principle, is uncorrelated.
- When a flavor symmetry is assumed to govern the gauge transformation of the leptons, electroweak symmetry breaking will result in an effective neutrino mass matrix M_ν^{eff} , which is parameterized with fewer parameters
- Diagonalising the M_ν^{eff} will lead to three complex mass eigenstates $(\tilde{m}_1, \tilde{m}_2, \tilde{m}_3)$ and additional relationship(s) among these complex mass eigenstates \rightarrow neutrino mass sum rules

Sum Rule	Group	Seesaw Type
$\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$	$A_4[158, 247-251]; S_4[252]; A_5[75]$	Weinberg
$\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$	$\Delta(54)[253]; S_4[254]$	Type II
$\tilde{m}_1 + 2\tilde{m}_2 = \tilde{m}_3$	$S_4[255]$	Type II
$2\tilde{m}_2 + \tilde{m}_3 = \tilde{m}_1$	$A_4 [117, 118, 158, 237, 248-251, 256-262]$ $S_4[133, 263]; T'[156, 246, 264-267]; T_7[268]$	Weinberg
$2\tilde{m}_2 + \tilde{m}_3 = \tilde{m}_1$	$A_4[269]$	Type II
$\tilde{m}_1 + \tilde{m}_2 = 2\tilde{m}_3$	$S_4[270]$	Dirac
$\tilde{m}_1 + \tilde{m}_2 = 2\tilde{m}_3$	$L_e - L_\mu - L_\tau[271]$	Type II
$\tilde{m}_1 + \frac{\sqrt{3}+1}{2}\tilde{m}_3 = \frac{\sqrt{3}-1}{2}\tilde{m}_2$	$A_5'[272]$	Weinberg
$\tilde{m}_1^{-1} + \tilde{m}_2^{-1} = \tilde{m}_3^{-1}$	$A_4[158]; S_4[247, 254]; A_5[78, 166]$	Type I
$\tilde{m}_1^{-1} + \tilde{m}_2^{-1} = \tilde{m}_3^{-1}$	$S_4[254]$	Type III
$2\tilde{m}_2^{-1} + \tilde{m}_3^{-1} = \tilde{m}_1^{-1}$	$A_4[118, 158, 235, 237, 238, 245, 273-281]; T'[246]$	Type I
$\tilde{m}_1^{-1} + \tilde{m}_3^{-1} = 2\tilde{m}_2^{-1}$	$A_4[282-284]; T'[285]$	Type I
$\tilde{m}_3^{-1} \pm 2i\tilde{m}_2^{-1} = \tilde{m}_1^{-1}$	$\Delta(96)[286]$	Type I
$\tilde{m}_1^{1/2} - \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$	$A_4[236]$	Type I
$\tilde{m}_1^{1/2} + \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$	$A_4[287]$	Scotogenic
$\tilde{m}_1^{-1/2} + \tilde{m}_2^{-1/2} = 2\tilde{m}_3^{-1/2}$	$S_4[240]$	Inverse

Table 1. Sum rules for complex light neutrino mass eigenvalues (\tilde{m}_i) derived from combinations of neutrino mass generation mechanisms and discrete flavor symmetries (arxiv.org/pdf/2310.20681)

Complex mass in mass rules

- Complex mass in the polar form $\tilde{m}_i = m_i e^{i\phi_i}$, where m_i is the physical mass eigenstate; in general form, $\phi_i \in [0; 2\pi]$ but α_{21} and α_{31} is from $[0; \pi]$ following the reference arxiv.org/abs/0804.3627
- It connects to standard parameterization α_{21} and α_{31} by the PMNS matrix
- Assume a sum rule $\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$
 - $\rightarrow m_1 e^{i\phi_1} + m_2 e^{i\phi_2} = m_3 e^{i\phi_3}$
 - $\rightarrow m_1 + m_2 e^{i(\phi_2 - \phi_1)} = m_3 e^{i(\phi_3 - \phi_1)}$
 - $\rightarrow m_1 + m_2(\cos\alpha_{21} + i\sin\alpha_{21}) = m_3(\cos\alpha_{31} + i\sin\alpha_{31})$
- Two relations can be concluded:
 - Real: $m_1 + m_2 \cos\alpha_{21} = m_3 \cos\alpha_{31}$
 - Imaginary: $m_2 \sin\alpha_{21} = m_3 \sin\alpha_{31}$
- $m_2, m_3 = f(m_1, \alpha_{21}, \alpha_{31})$

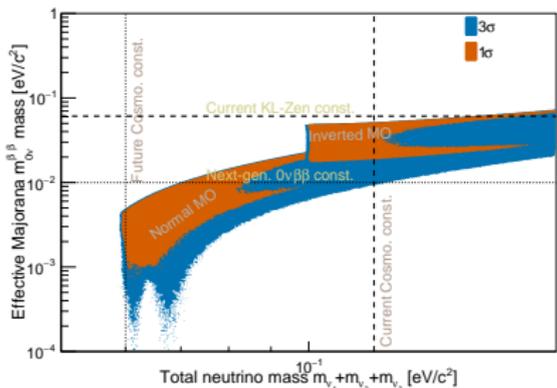
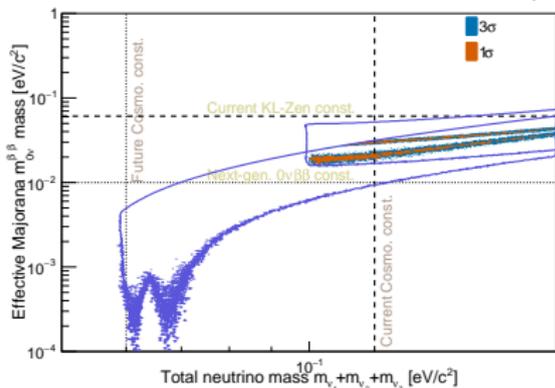
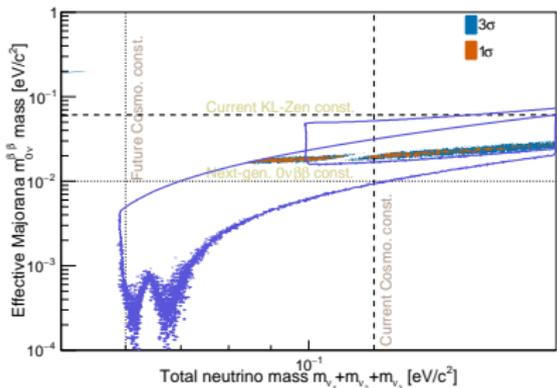
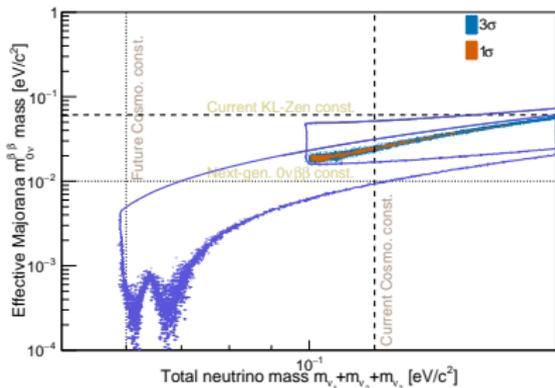
Input constraints

	Current (eV)	Future (eV)
Beta decay (KATRIN)	0.8	0.2
Neutrinoless double beta decay (KAMLand-Zen2016)	0.061 - 0.175	0.02
Cosmology	0.12	0.06

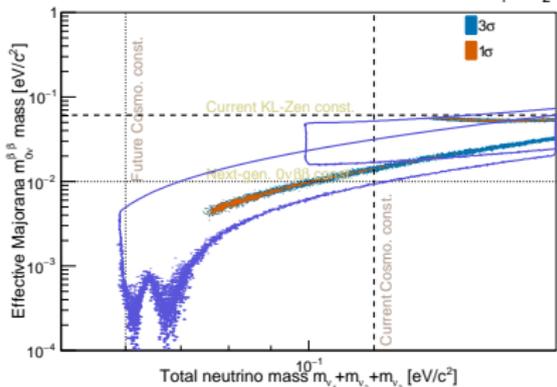
NuFIT 5.0 (2020)

	Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 2.7$)		
	bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range	
without SK atmospheric data	$\sin^2 \theta_{12}$	$0.304^{+0.013}_{-0.012}$	$0.269 \rightarrow 0.343$	$0.304^{+0.013}_{-0.012}$	$0.269 \rightarrow 0.343$
	$\theta_{12}/^\circ$	$33.44^{+0.78}_{-0.75}$	$31.27 \rightarrow 35.86$	$33.45^{+0.78}_{-0.75}$	$31.27 \rightarrow 35.87$
	$\sin^2 \theta_{23}$	$0.570^{+0.018}_{-0.024}$	$0.407 \rightarrow 0.618$	$0.575^{+0.017}_{-0.021}$	$0.411 \rightarrow 0.621$
	$\theta_{23}/^\circ$	$49.0^{+1.1}_{-1.4}$	$39.6 \rightarrow 51.8$	$49.3^{+1.0}_{-1.2}$	$39.9 \rightarrow 52.0$
	$\sin^2 \theta_{13}$	$0.02221^{+0.00068}_{-0.00062}$	$0.02034 \rightarrow 0.02430$	$0.02240^{+0.00062}_{-0.00062}$	$0.02053 \rightarrow 0.02436$
	$\theta_{13}/^\circ$	$8.57^{+0.13}_{-0.12}$	$8.20 \rightarrow 8.97$	$8.61^{+0.12}_{-0.12}$	$8.24 \rightarrow 8.98$
	$\delta_{CP}/^\circ$	195^{+51}_{-25}	$107 \rightarrow 403$	286^{+27}_{-32}	$192 \rightarrow 360$
	$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.42^{+0.21}_{-0.20}$	$6.82 \rightarrow 8.04$	$7.42^{+0.21}_{-0.20}$	$6.82 \rightarrow 8.04$
	$\frac{\Delta m_{3l}^2}{10^{-3} \text{ eV}^2}$	$+2.514^{+0.028}_{-0.027}$	$+2.431 \rightarrow +2.598$	$-2.497^{+0.028}_{-0.028}$	$-2.583 \rightarrow -2.412$

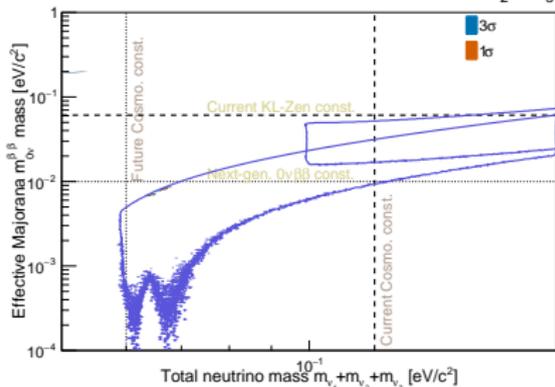
No rule applied

Sum rule: $\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$ Sum rule: $\tilde{m}_1 + 2\tilde{m}_2 = \tilde{m}_3$ Sum rule: $\tilde{m}_1 + \tilde{m}_2 = 2\tilde{m}_3$ 

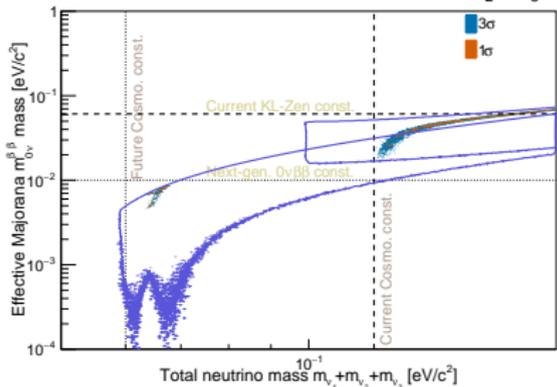
$$\text{Sum rule: } \tilde{m}_1^{-1} + \tilde{m}_2^{-1} = \tilde{m}_3^{-1}$$



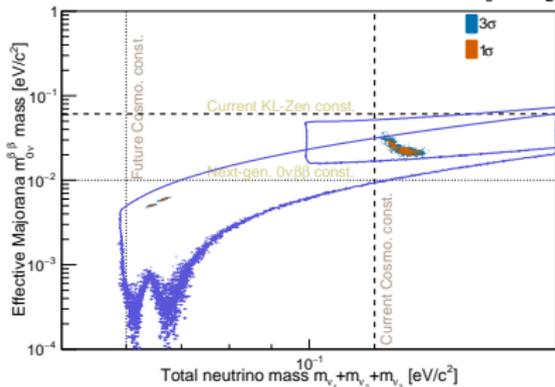
$$\text{Sum rule: } 2\tilde{m}_2^{-1} + \tilde{m}_3^{-1} = \tilde{m}_1^{-1}$$



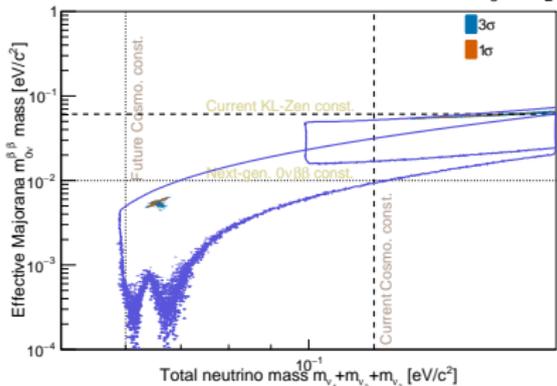
$$\text{Sum rule: } \tilde{m}_2^{-1} + \tilde{m}_3^{-1} = 2\tilde{m}_1^{-1}$$



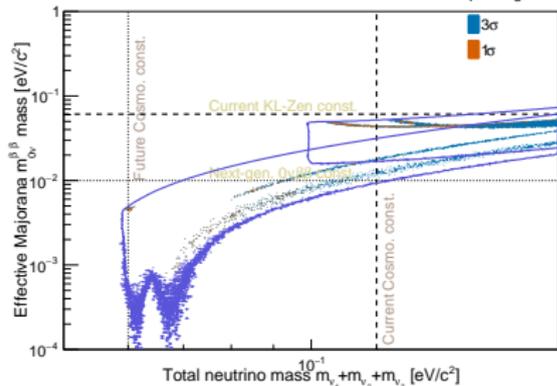
$$\text{Sum rule: } \tilde{m}_3^{-1} + 2\tilde{m}_2^{-1} = \tilde{m}_1^{-1}$$



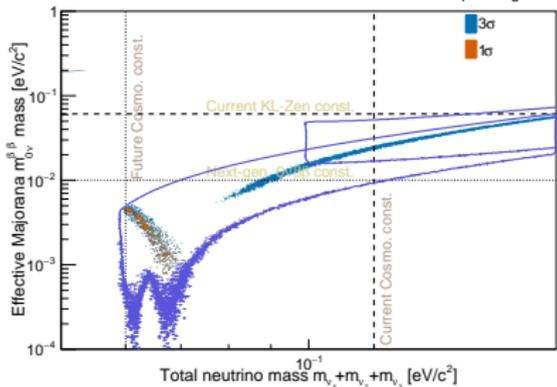
$$\text{Sum rule: } \tilde{m}_3^{-1} - 2\tilde{m}_2^{-1} = \tilde{m}_1^{-1}$$



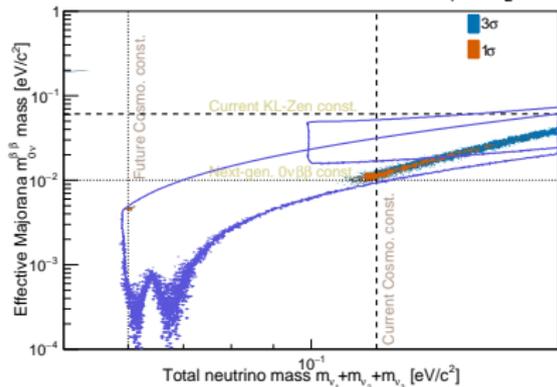
$$\text{Sum rule: } \tilde{m}_1^{1/2} - \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$$



$$\text{Sum rule: } \tilde{m}_1^{1/2} + \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$$



$$\text{Sum rule: } \tilde{m}_1^{-1/2} + \tilde{m}_2^{-1/2} = 2\tilde{m}_3^{-1/2}$$



Conclusion

Sum Rule	Group	Seesaw Type	
$\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$	A_4 [158, 247–251]; S_4 [252]; A_5 [75]	Weinberg	Largely allowed
$\tilde{m}_1 + \tilde{m}_2 = \tilde{m}_3$	$\Delta(54)$ [253]; S_4 [254]	Type II	
$\tilde{m}_1 + 2\tilde{m}_2 = \tilde{m}_3$	S_4 [255]	Type II	Largely allowed
$2\tilde{m}_2 + \tilde{m}_3 = \tilde{m}_1$	A_4 [117, 118, 158, 237, 248–251, 256–262] S_4 [133, 263]; T' [156, 246, 264–267]; T_7 [268]	Weinberg	Disfavored
$2\tilde{m}_2 + \tilde{m}_3 = \tilde{m}_1$	A_4 [269]	Type II	
$\tilde{m}_1 + \tilde{m}_2 = 2\tilde{m}_3$	S_4 [270]	Dirac	Largely allowed
$\tilde{m}_1 + \tilde{m}_2 = 2\tilde{m}_3$	$L_e - L_\mu - L_\tau$ [271]	Type II	
$\tilde{m}_1 + \frac{\sqrt{3}+1}{2}\tilde{m}_3 = \frac{\sqrt{3}-1}{2}\tilde{m}_2$	A_5' [272]	Weinberg	Disfavored
$\tilde{m}_1^{-1} + \tilde{m}_2^{-1} = \tilde{m}_3^{-1}$	A_4 [158]; S_4 [247, 254]; A_5 [78, 166]	Type I	Largely allowed
$\tilde{m}_1^{-1} + \tilde{m}_2^{-1} = \tilde{m}_3^{-1}$	S_4 [254]	Type III	Largely allowed
$2\tilde{m}_2^{-1} + \tilde{m}_3^{-1} = \tilde{m}_1^{-1}$	A_4 [118, 158, 235, 237, 238, 245, 273–281]; T' [246]	Type I	Marginally allowed
$\tilde{m}_1^{-1} + \tilde{m}_3^{-1} = 2\tilde{m}_2^{-1}$	A_4 [282–284]; T' [285]	Type I	
$\tilde{m}_3^{-1} \pm 2i\tilde{m}_2^{-1} = \tilde{m}_1^{-1}$	$\Delta(96)$ [286]	Type I	Marginally allowed
$\tilde{m}_1^{1/2} - \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$	A_4 [236]	Type I	Largely allowed
$\tilde{m}_1^{1/2} + \tilde{m}_3^{1/2} = 2\tilde{m}_2^{1/2}$	A_4 [287]	Scotogenic	Largely allowed
$\tilde{m}_1^{-1/2} + \tilde{m}_2^{-1/2} = 2\tilde{m}_3^{-1/2}$	S_4 [240]	Inverse	Marginally allowed

Here we just look at the sum rules in their most basic form, without any correction.

Thank you for your listening!