

Natural Metric-Affine Inflation

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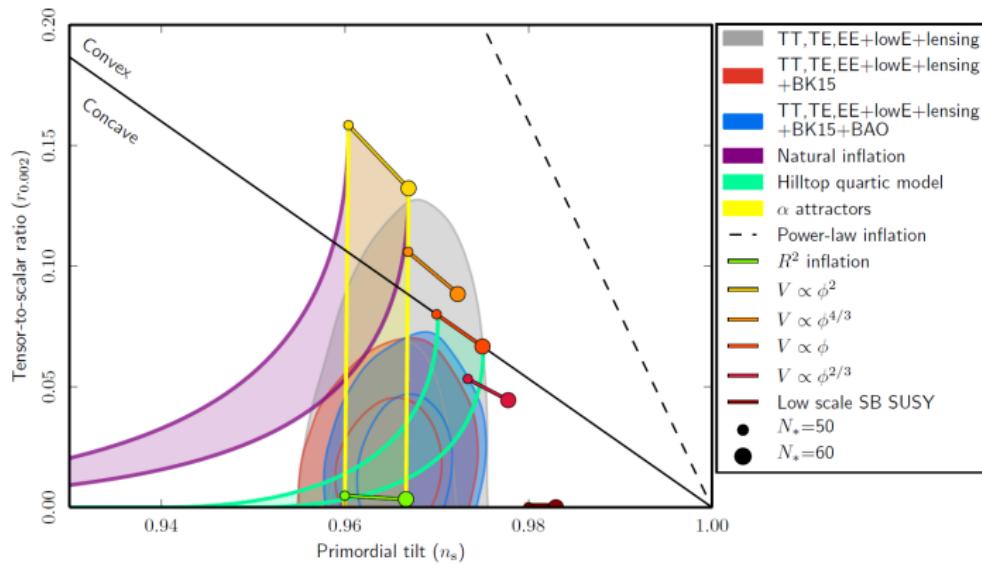
based on

arXiv:2403.18004 (JCAP 06 (2024) 033)

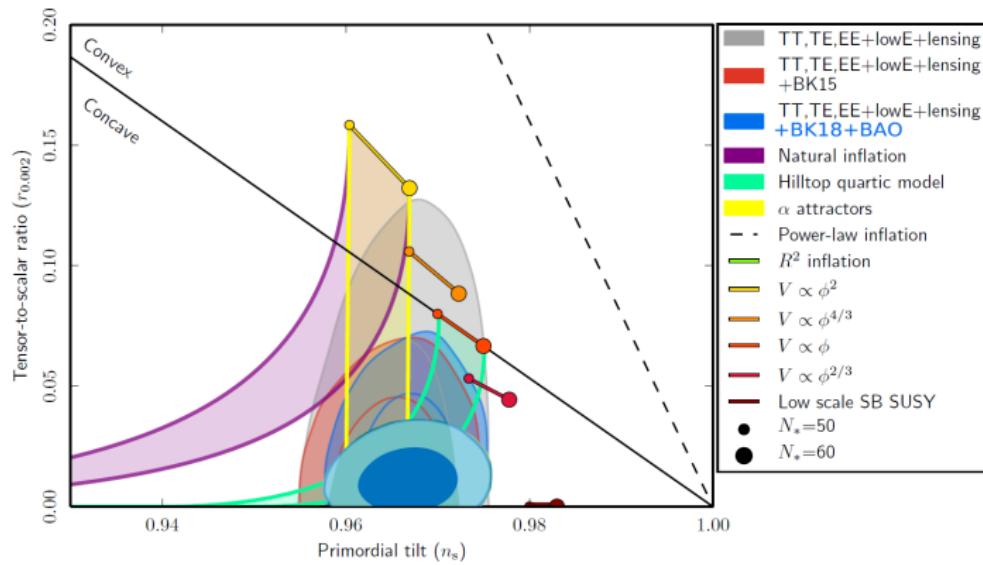
with

A. Salvio (Univ. Rome Tor Vergata & INFN)

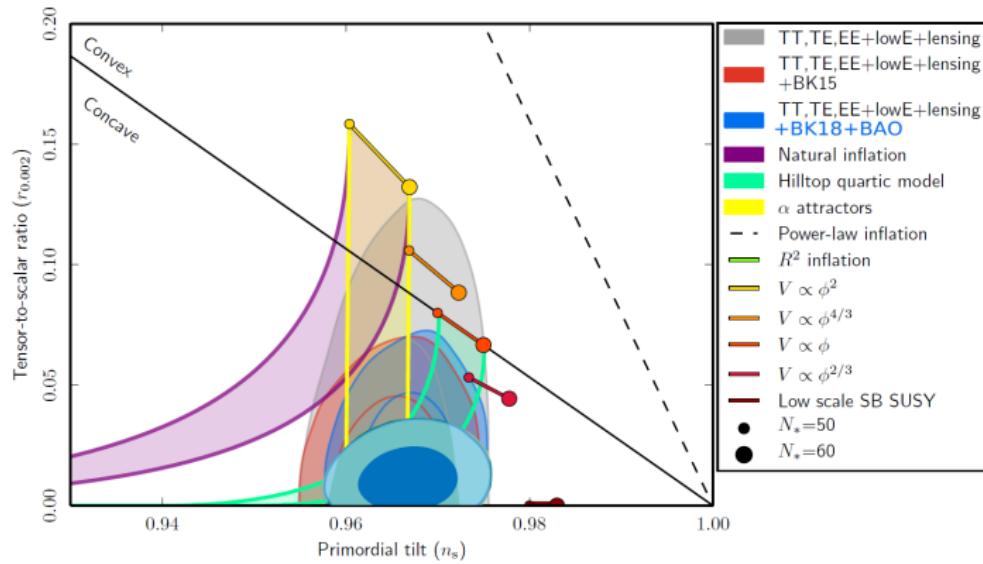




- NI became strongly disfavored after BICEP/Keck 2018 data
- several proposals to save it by modifying gravity:
 - $\xi[1 + \cos(\phi)]R \rightarrow$ OK only at 2σ (Ferreira et al. 1806.05511)
 - $\xi\phi^n R \rightarrow$ OK only at 2σ (Bostan, 2209.02434; dos Santos et al., 2312.12286)
 - Palatini $R^2 \rightarrow$ OK! but $(\partial\phi)^4$ (Antoniadis et al., 1812.00847)
 - probably more?



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The properties of torsion-free spacetime are essentially described by:

- the affine connection: $\mathcal{A}_{\alpha\beta}^{\lambda} \rightarrow$ parallel transport
- the metric tensor: $g_{\mu\nu} \rightarrow$ distance

The connection coefficients and metric tensor are fundamentally independent quantities. They exhibit no *a priori* known relationship. If they are to have any relationship, it must derive from

- additional constraints (metric formalism $\mathcal{A}_{\mu\sigma}^{\rho} = \Gamma_{\mu\sigma}^{\rho}$)
- EoM for \mathcal{A} & g (Palatini formalism)
 - minimal theories $\Rightarrow \mathcal{A}_{\mu\sigma}^{\rho} = \Gamma_{\mu\sigma}^{\rho}$
 \Rightarrow metric \sim Palatini
 - non-minimal theories $\Rightarrow \mathcal{A}_{\mu\sigma}^{\rho} \neq \Gamma_{\mu\sigma}^{\rho}$
 \Rightarrow metric \neq Palatini (e.g. Koivisto & Kurki-Suonio: 0509422)

• Palatini NI action •

$$S_{\text{NI}} = \int d^4x \sqrt{-g_J} \left[\alpha(\phi) \mathcal{R}_J - \frac{\partial_\mu \phi \partial^\mu \phi}{2} - V(\phi) \right]$$

$$V(\phi) = \Lambda^4 \left[1 + \cos \left(\frac{\phi}{f} \right) \right]$$

$$\alpha(\phi) = \frac{M_P^2}{2} \left[1 + \xi \left(1 + \cos \left(\frac{\phi}{f} \right) \right) \right] > 0$$

- notation: $\begin{cases} \mathcal{R} \rightarrow \text{curvature from generic } \mathcal{A}_\mu{}^\rho{}_\sigma \\ R \rightarrow \text{curvature from Levi-Civita } \Gamma_\mu{}^\rho{}_\sigma \end{cases}$
- Einstein frame: $g_{\mu\nu}^E = F(\phi) g_{\mu\nu}^J, F \equiv \frac{2\alpha}{M_P^2}$ N.B. Palatini $\Rightarrow \mathcal{R}_J = F R_E$

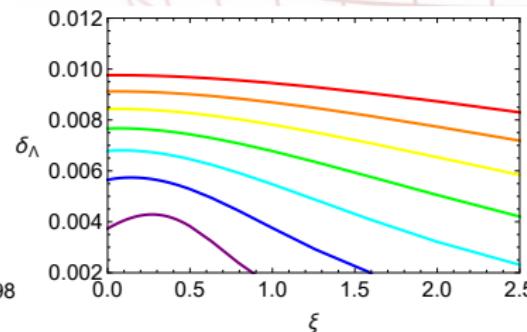
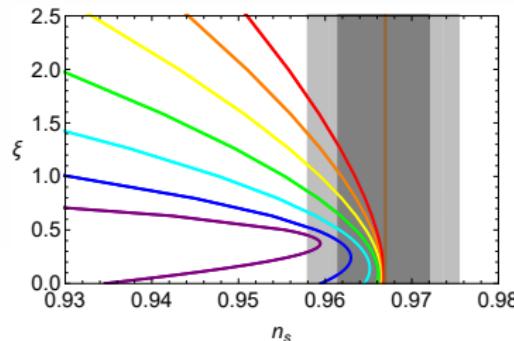
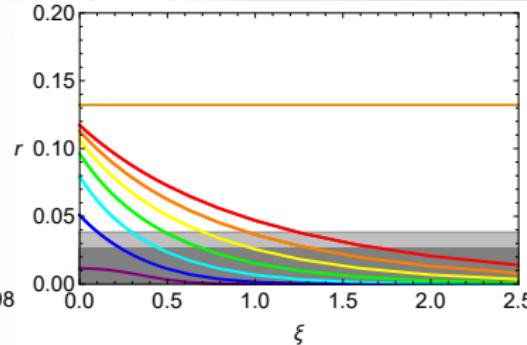
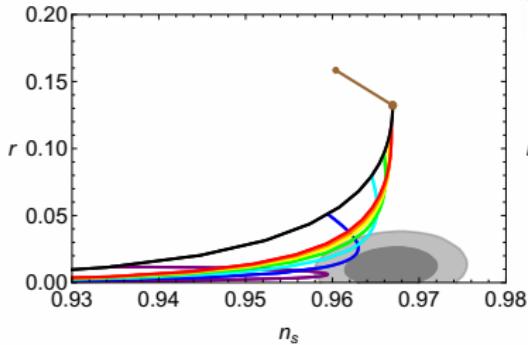
$$S_{\text{NI}} = \int d^4x \sqrt{-g_E} \left[\frac{M_P^2}{2} R_E - \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) \right]$$

$$\frac{d\chi}{d\phi} = \sqrt{\frac{1}{F(\phi)}} = \sqrt{\frac{M_P^2}{2\alpha(\phi)}} \quad \leftarrow \text{no } \frac{3}{2} \left(\frac{F'}{F} \right)^2 \text{ like in metric}$$

$$U(\chi) = \frac{V(\phi(\chi))}{F^2(\phi(\chi))}$$

• Palatini results •

$$N_e = 60 \quad \delta_f = f/M_P \quad \delta_\Lambda = \Lambda/M_P$$



- $\delta_f = 4$
- $\delta_f = 6$
- $\delta_f = 8$
- $\delta_f = 10$
- $\delta_f = 12$
- $\delta_f = 14$
- $\delta_f = 16$

- quadratic
- natural

- BICEP & Planck

- any way to get better results?
- YES! allow for torsion \rightarrow MAG

but in a way that new dof's are not generated!

- $A_{\alpha\beta}^{\lambda} \neq A_{\beta\alpha}^{\lambda}$

$$\mathcal{F}_{\mu\nu}{}^{\rho}{}_{\sigma} \equiv \partial_{\mu}A_{\nu}{}^{\rho}{}_{\sigma} - \partial_{\nu}A_{\mu}{}^{\rho}{}_{\sigma} + A_{\mu}{}^{\rho}{}_{\lambda}A_{\nu}{}^{\lambda}{}_{\sigma} - A_{\nu}{}^{\rho}{}_{\lambda}A_{\mu}{}^{\lambda}{}_{\sigma}$$

$$\text{Holst inv. } \rightarrow \tilde{\mathcal{R}} \equiv \frac{1}{\sqrt{-g}} \epsilon^{\mu\nu\rho\sigma} \mathcal{F}_{\mu\nu\rho\sigma}$$

$$\text{cf. } \mathcal{R} \equiv \mathcal{F}_{\mu\nu}{}^{\mu\nu}$$

where $\epsilon^{\mu\nu\rho\sigma}$ is the totally antisymmetric Levi-Civita symbol with $\epsilon^{0123} = 1$

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The properties of ~~torsion-free~~ spacetime are essentially described by:

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$$S_{\text{NI}} = \int d^4x \sqrt{-g_J} \left[\alpha(\phi) \mathcal{R}_J + \beta(\phi) \tilde{\mathcal{R}}_J - \frac{\partial_\mu \phi \partial^\mu \phi}{2} - V(\phi) \right]$$

$$V(\phi) = \Lambda^4 \left[1 + \cos \left(\frac{\phi}{f} \right) \right] \quad \alpha(\phi) = \frac{M_P^2}{2} \left[1 + \xi \left(1 + \cos \left(\frac{\phi}{f} \right) \right) \right] > 0$$

$$\beta(\phi) = \beta_0 + \frac{M_P^2}{2} \tilde{\xi} \left(\cos \left(\frac{\phi}{f} \right) + 1 \right)$$

$\frac{M_P^2}{4\beta_0}$ → Barbero-Irmizzi par.

- it is possible to integrate out the \tilde{R} term
- performing all the computations ...

$$S_{\text{NI}} = \int d^4x \sqrt{-g} \left[\alpha \mathcal{R}_J - \left[1 + \frac{12(\alpha' \beta + \alpha \beta')^2}{\alpha(\alpha^2 + 4\beta^2)} \right] \frac{\partial_\mu \phi \partial^\mu \phi}{2} - V \right]$$

- allowing torsion changes the inflaton kinetic term

Einstein frame:

- $g_{\mu\nu}^E = F(\phi)g_{\mu\nu}^J$, $F \equiv \frac{2\alpha}{M_P^2}$

N.B. MAG $\Rightarrow \mathcal{R}_J = F R_E$

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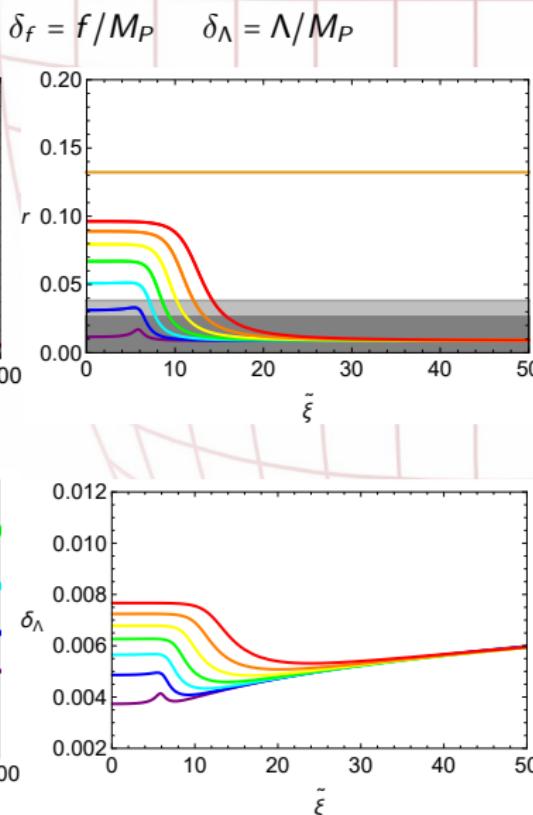
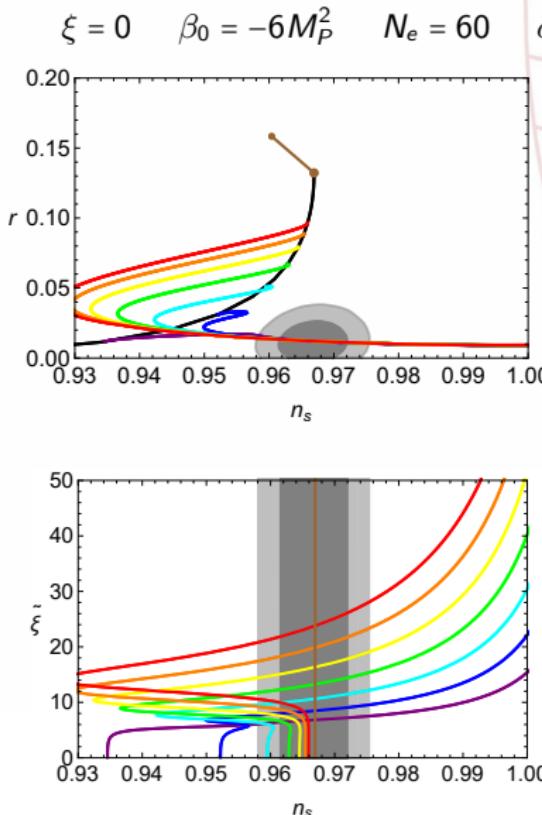
$$\frac{d\chi}{d\phi} = M_P \sqrt{\frac{1}{2\alpha} \left[1 + \frac{12(\alpha'\beta + \alpha\beta')^2}{\alpha(\alpha^2 + 4\beta^2)} \right]} \quad \leftarrow \text{new term from } \beta \tilde{R}$$

$$U(\chi) = \frac{V(\phi(\chi))}{F^2(\phi(\chi))} \quad \leftarrow \text{same as before}$$

N.B.

- $\beta < 0$ allowed
- symmetry: $\beta \rightarrow -\beta \quad \Rightarrow \quad \tilde{\xi} > 0, \beta_0 \geq 0$

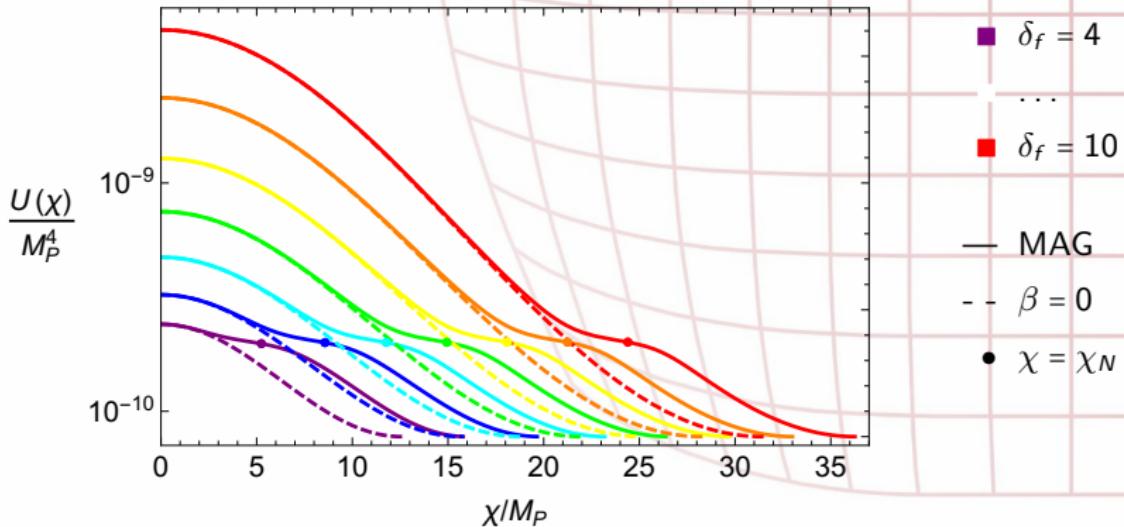
$$\bullet \xi = 0 \text{ } \& \text{ } \tilde{\xi} > 0 \text{ } \& \text{ } \beta_0 < 0 \bullet$$



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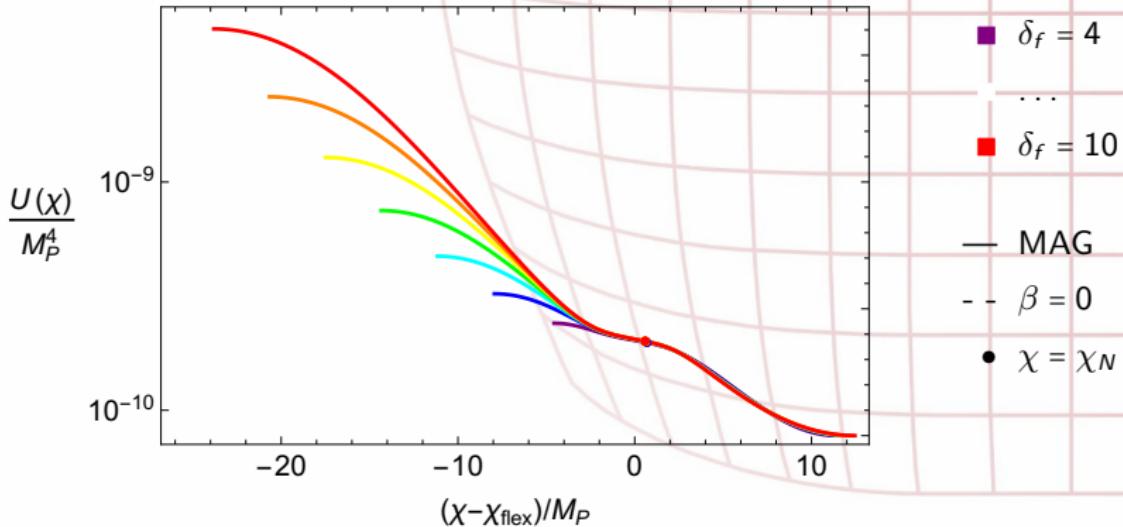
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$$\xi = 0 \quad \beta_0 = -6M_P^2 \quad N_e = 60 \quad n_s \simeq 0.97 \quad \delta_f = f/M_P \quad \delta_\Lambda = \Lambda/M_P$$

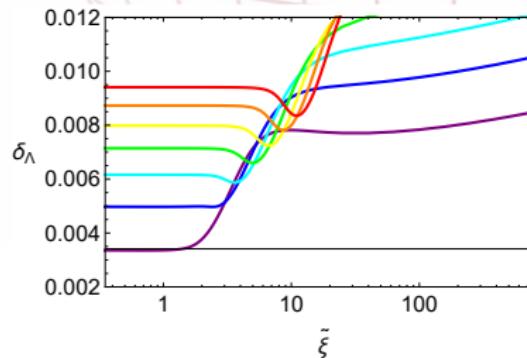
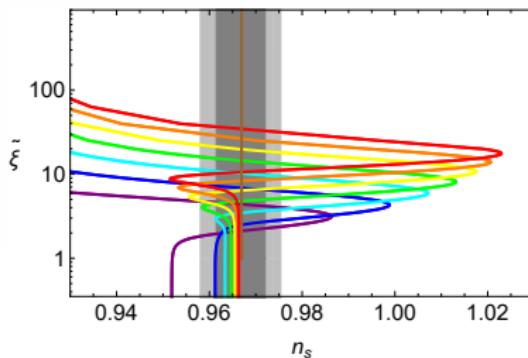
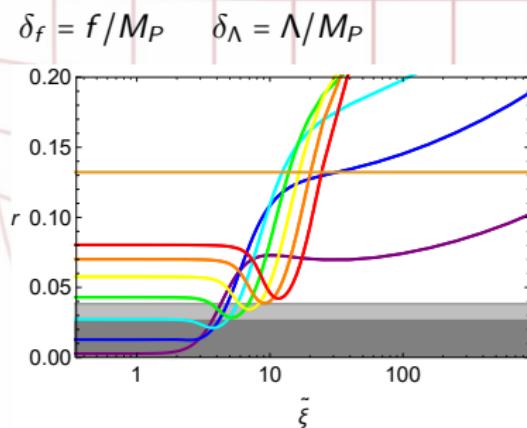
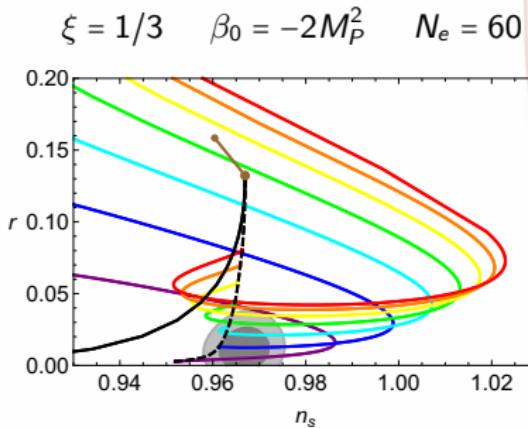


- inflection point inflation!!!

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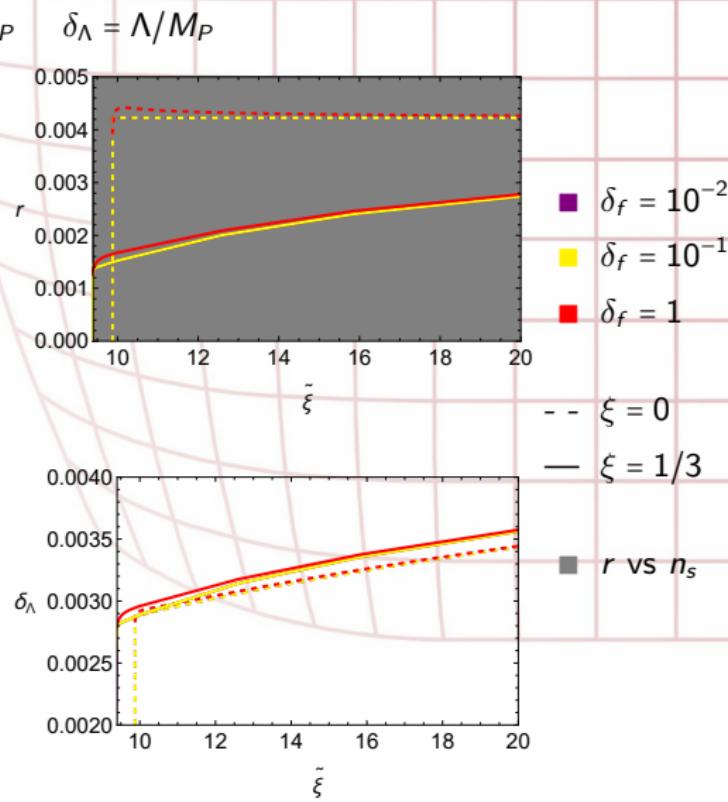
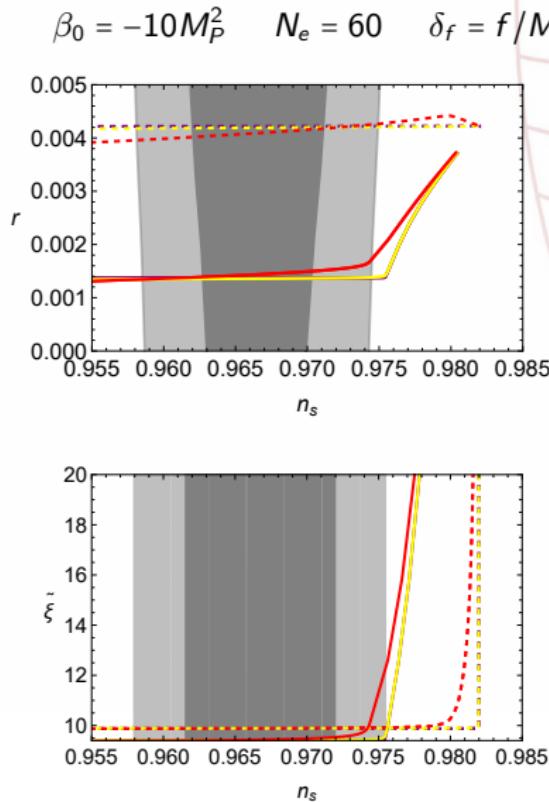


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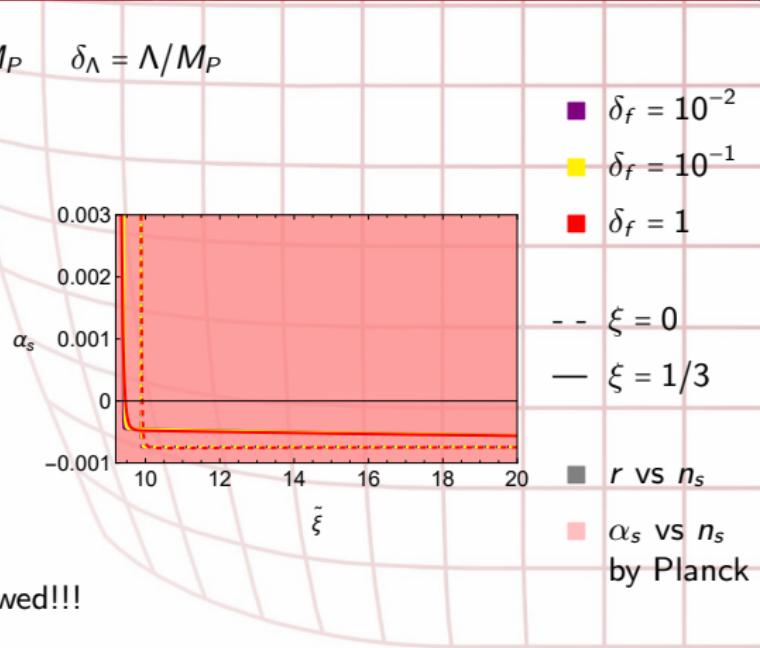
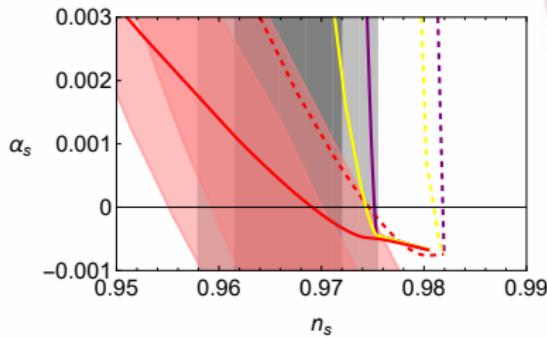
- quadratic
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- - - $\beta = 0$

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$$\beta_0 = -10M_P^2 \quad N_e = 60 \quad \delta_f = f/M_P \quad \delta_\Lambda = \Lambda/M_P$$



- $\xi > 0$ & $\tilde{\xi} > 0 \Rightarrow f \lesssim M_P$ allowed!!!

- NI strongly disfavored after Planck+BICEP 2018 data
- introducing a non-minimal coupling to gravity
 - compatible at 2σ in the Palatini formulation
- allowing for torsion (i.e. $\tilde{\mathcal{R}}$) (MAG formalism)
 - compatible at 1σ with data
 - allows also for subPlanckian f !!!

Grazie! - Thank you! - Aitäh!