

Broken scale invariant unparticle physics and its prospective effect on the MuonE experiment

arxiv:2304.04439

Van Dung Le^(a), Duc Ninh Le^(b), Duc Truyen Le^(c), Van Cuong Le^(a)

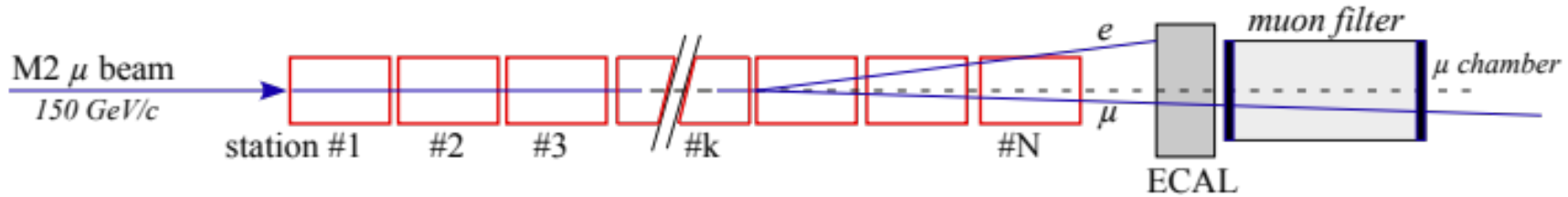
(a) Ho Chi Minh University of Science - Vietnam National University

(b) PHENIKAA University, Hanoi 12116, Vietnam

(c) National Tsing Hua University, Hsinchu, Taiwan (NTHU)



MuonE Experiment



Proposed: 2017

Location: CERN

Experiment: $e\mu \rightarrow e\mu$

Purpose: evaluate $\Delta\alpha_{had}(t)$, $t < 0$

$$\sqrt{s} \approx 0.4\text{GeV}, t = (p_1 - p_3)^2$$

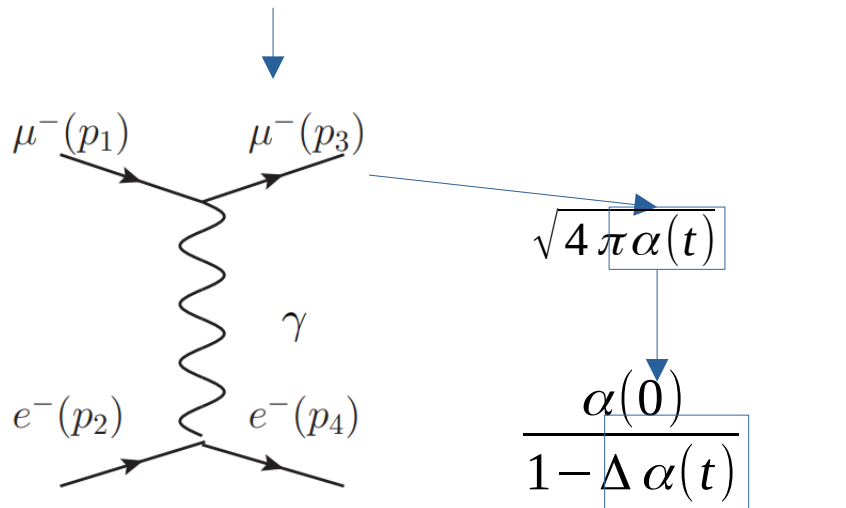
Systematic accuracy: 10 ppm

→ Competitive: Sensitive to small new physics effects

→ Offer an independent measurement of a_μ^{had}

Measure cross

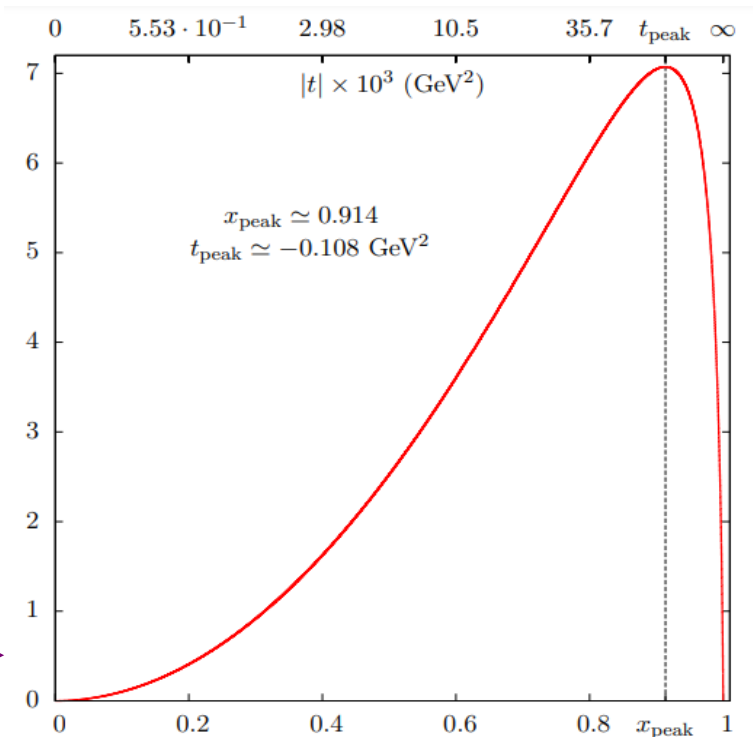
$e\mu \rightarrow e\mu$



pQCD + Exp.
data

$$\Delta\alpha_{had}(t) + \Delta\alpha_{lep}(t)$$

$$a_{\mu}^{had} = \frac{\alpha}{\pi} \int dx (1-x) \Delta\alpha_{had}[t(x)]$$



Unparticle

- Unparticle was proposed by Howard Georgi in 2007 to discuss some simple aspects of the low-energy physics of a nontrivial scale invariant sector of an effective field theory.
- Can be directly detected at colliders as missing energy, e.g.

$$e^+ + e^- \rightarrow \gamma + \textit{unparticle}$$

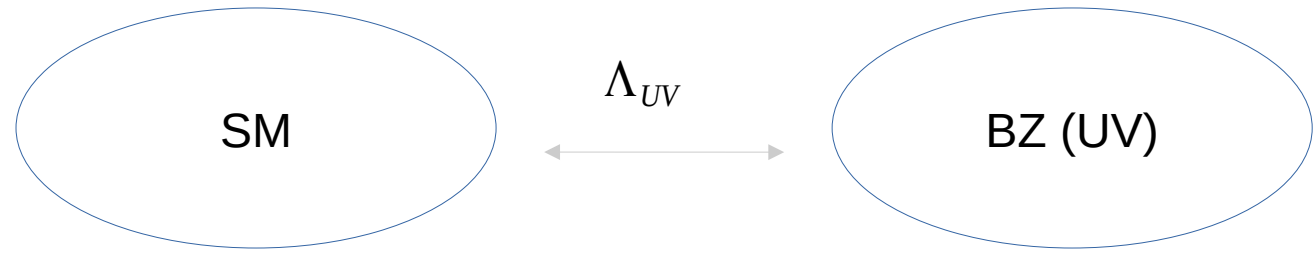
$$p + p \rightarrow \textit{jet} + \textit{unparticle}$$

- We only consider here the indirect effects of unparticle at the MuonE experiment

Unparticle scheme

E

Very high energy



Scale M_{UV}

$$\frac{c_{UV}}{\Lambda_{UV}^{d_{SM}+d_{UV}-4}} O_{SM} O_{UV}$$

Scale Λ_U

$$\frac{c_U}{\Lambda_{UV}^{d_{SM}+d_U-4}} O_{SM} O_U$$

Scale invariant

non-integer dimension $d=1.1, 1.2, \dots$

Below scale μ

Scale invariant is broken

Effective interactions with fermions

$$\frac{c_U}{\Lambda_{UV}^{d_{SM}+d_U-4}} O_{SM} O_U \longrightarrow \frac{\lambda_u}{M_Z^{d-1}} O_{SM} O_U \quad \text{parameter : } d, \lambda_i, \mu$$

Our convention

$$\frac{\lambda_S}{M_Z^{d-1}} \bar{f} f O_U, \quad \frac{\lambda_P}{M_Z^{d-1}} \bar{f} i \gamma^5 f O_U, \quad \frac{\lambda_V}{M_Z^{d-1}} \bar{f} \gamma_\mu f O_U^\mu, \quad \frac{\lambda_A}{M_Z^{d-1}} \bar{f} \gamma_\mu \gamma^5 f O_U^\mu,$$

Scalar

Pseudo-scalar

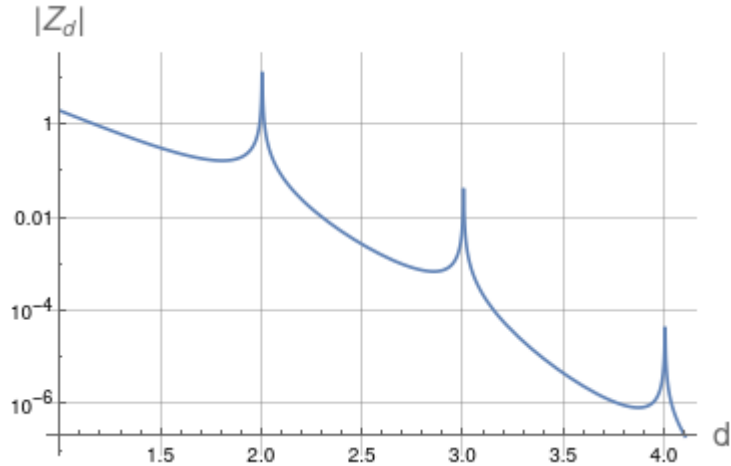
Vector

Axial-vector

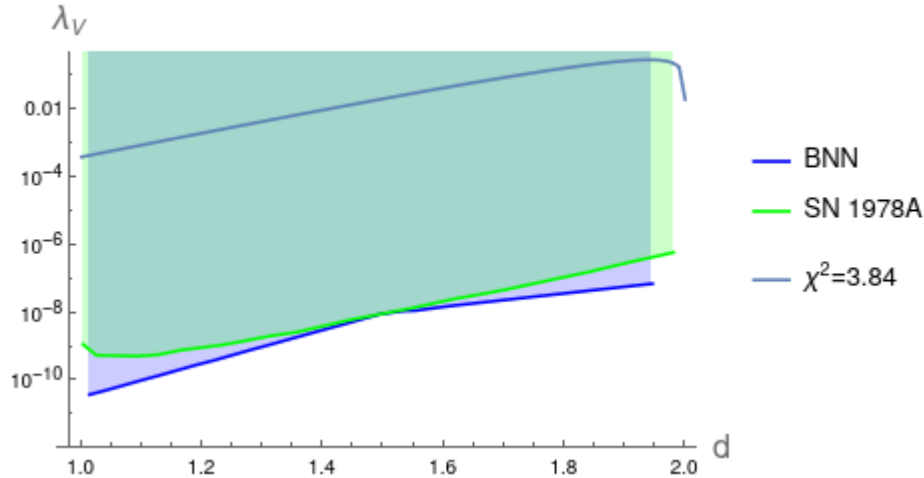
Propagator and implementation broken scale invariant

$$\Delta_F(k) = \frac{iZ_d}{(-k^2 - i\epsilon)^{2-d}}$$

$$\Delta_F^{\mu\nu}(k) = \frac{iZ_d}{(-k^2 - i\epsilon)^{2-d}} \left(-g^{\mu\nu} + a \frac{k^\mu k^\nu}{k^2} \right)$$



V parameter space



$$\Delta_F(k) = \frac{iZ_d}{(-k^2 + \mu^2 - i\epsilon)^{2-d}}$$

$$\Delta_F^{\mu\nu}(k) = \frac{iZ_d}{(-k^2 + \mu^2 - i\epsilon)^{2-d}} \left(-g^{\mu\nu} + a \frac{k^\mu k^\nu}{k^2} \right)$$

We can evade these constraints by simply choosing $\mu \geq 1 \text{ GeV}$

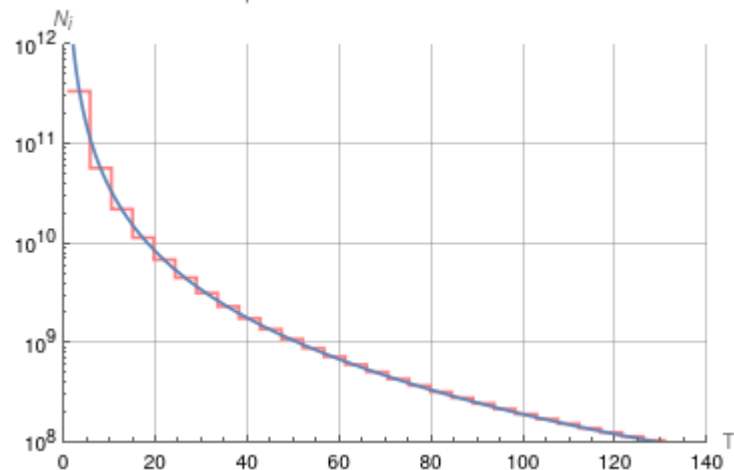
V. Barger, Y. Gao, W.-Y. Keung, D. Marfatia and V. N. Senoguz [arxiv:0801.3771]

Differential cross section

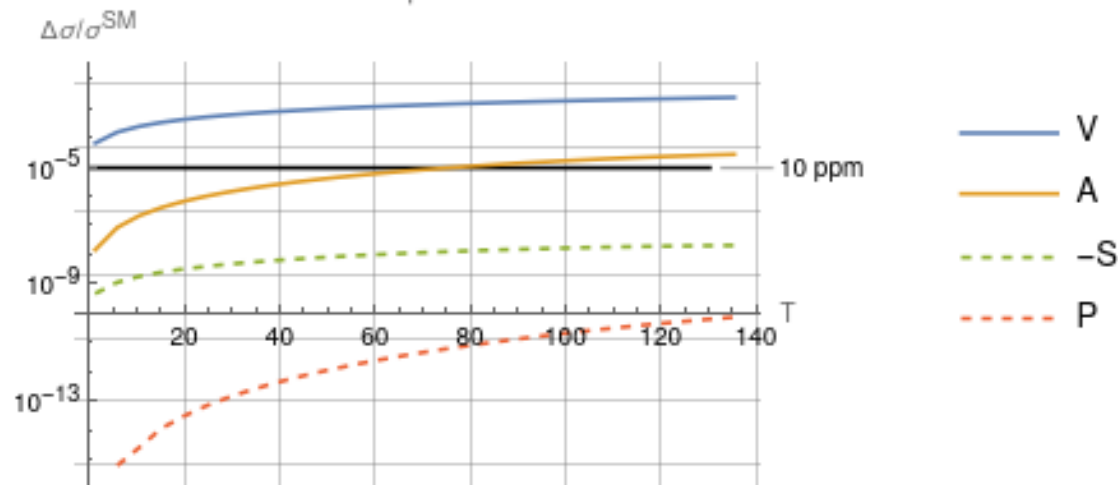
$$\begin{aligned} \frac{d\sigma_X}{dT} = & \text{Tr}[\gamma\gamma] \frac{c_1^2 \alpha^2(t)}{t^2} \\ & + \text{Tr}[\gamma X] \frac{c_1 \alpha(t)}{t} \frac{\lambda_X Z_d}{M_Z^{d-1} |t - \mu^2|^{2-d}} \\ & + \text{Tr}[X X] \left[\frac{\lambda_X Z_d}{M_Z^{d-1} |t - \mu^2|^{2-d}} \right]^2 \end{aligned}$$

| | type | T^0 | T | T^2 |
|--------------------|------|---|--|------------|
| $Tr[\gamma\gamma]$ | | $\underbrace{64E_\mu^2 m_e^2}$ | $-64E_\mu m_e^2 - 32m_e^3 - 32m_e m_\mu^2$ | $32m_e^2$ |
| $Tr[XX]$ | S | $64m_e^2 m_\mu^2$ | $32m_e^3 + 32m_e m_\mu^2$ | $16m_e^2$ |
| | P | 0 | 0 | $16m_e^2$ |
| | V | $\underbrace{64E_\mu^2 m_e^2}$ | $-64E_\mu m_e^2 - 32m_e^3 - 32m_e m_\mu^2$ | $32m_e^2$ |
| | A | $\underbrace{64E_\mu^2 m_e^2} + 128m_e^2 m_\mu^2$ | $-64E_\mu m_e^2 + 32m_e^3 + 32m_e m_\mu^2$ | $32m_e^2$ |
| $Tr[\gamma X]$ | -S | $64E_\mu m_e^2 m_\mu$ | $-32m_e^2 m_\mu$ | 0 |
| | -P | 0 | 0 | 0 |
| | V | $\underbrace{64E_\mu^2 m_e^2}$ | $-64E_\mu m_e^2 - 32m_e^3 - 32m_e m_\mu^2$ | $32m_e^2$ |
| | A | 0 | $64E_\mu m_e^2$ | $-32m_e^2$ |

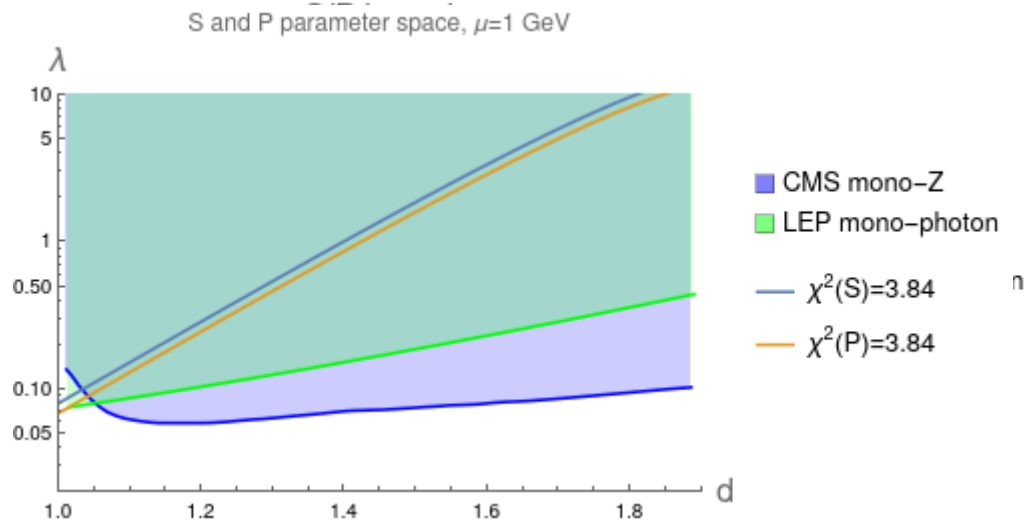
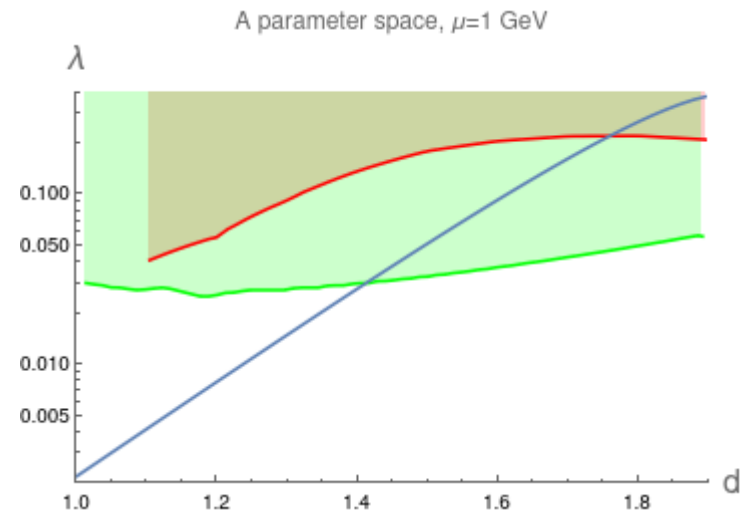
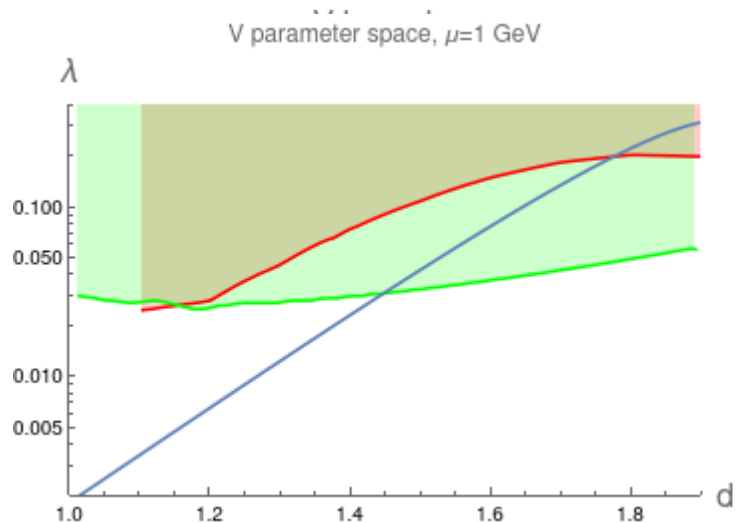
SM prediction of number of events



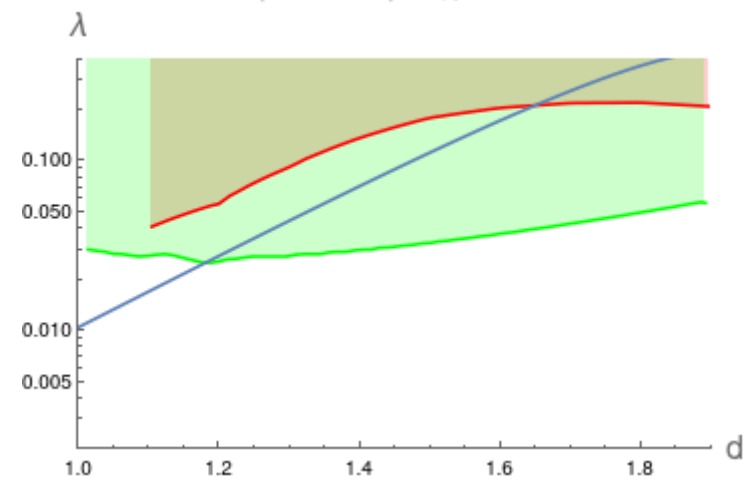
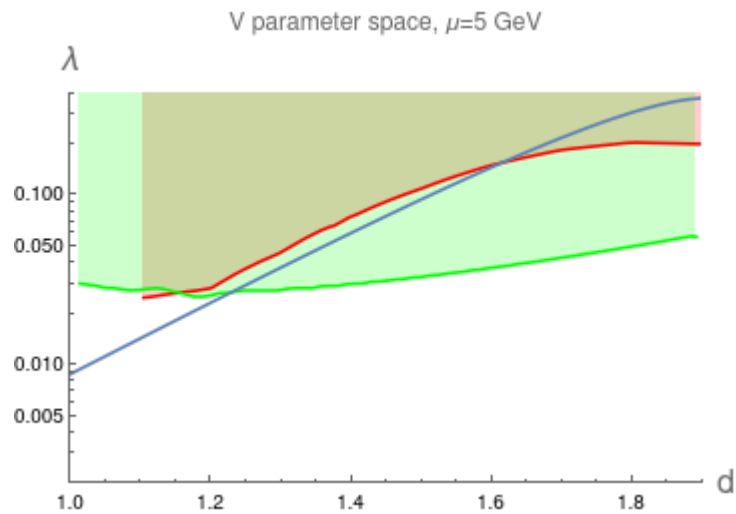
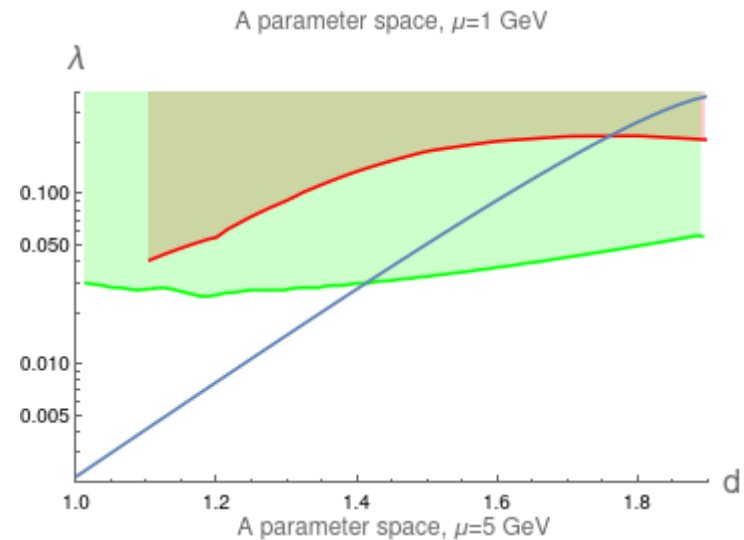
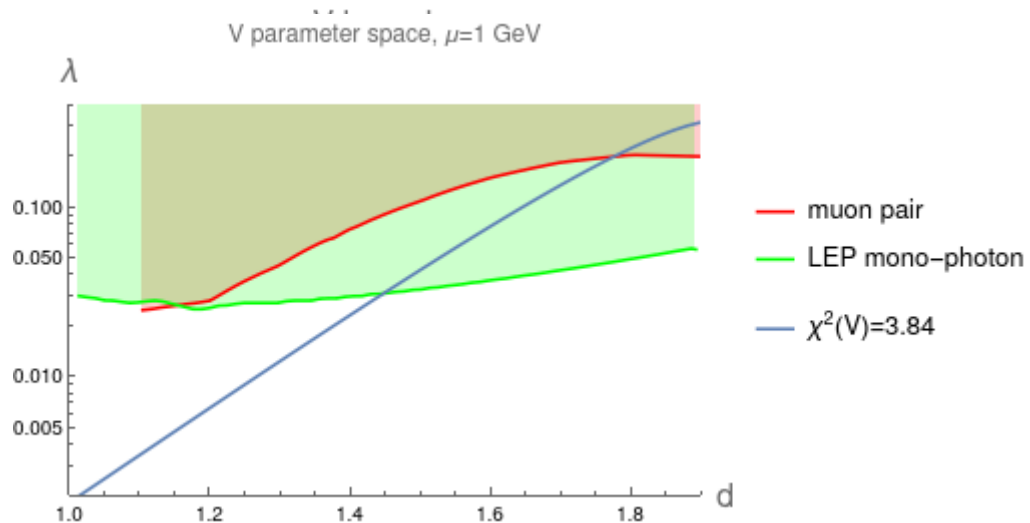
Fractional unparticle effects



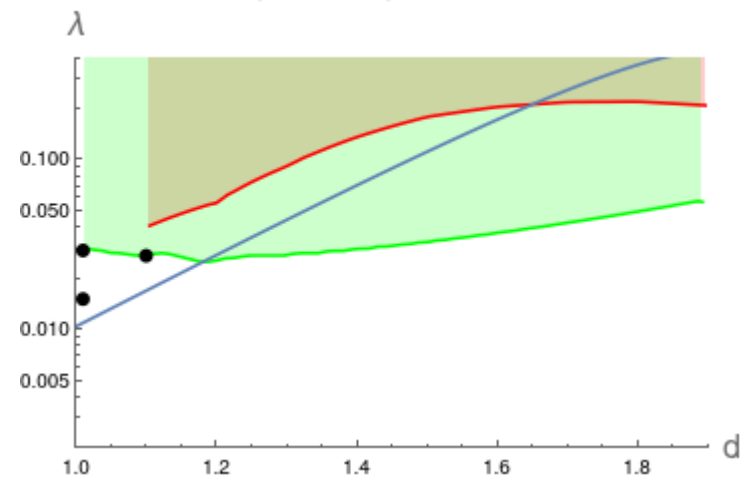
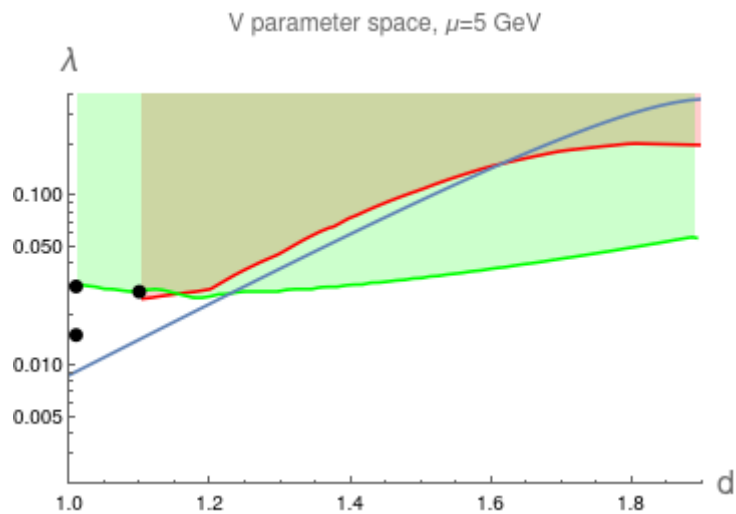
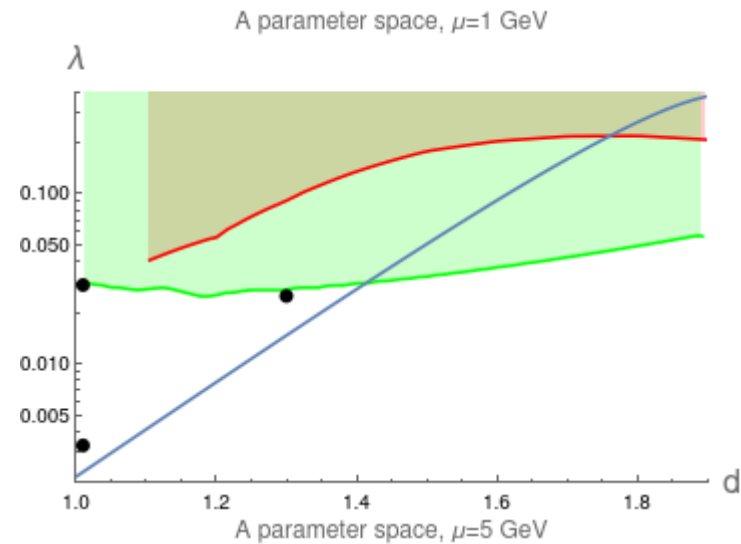
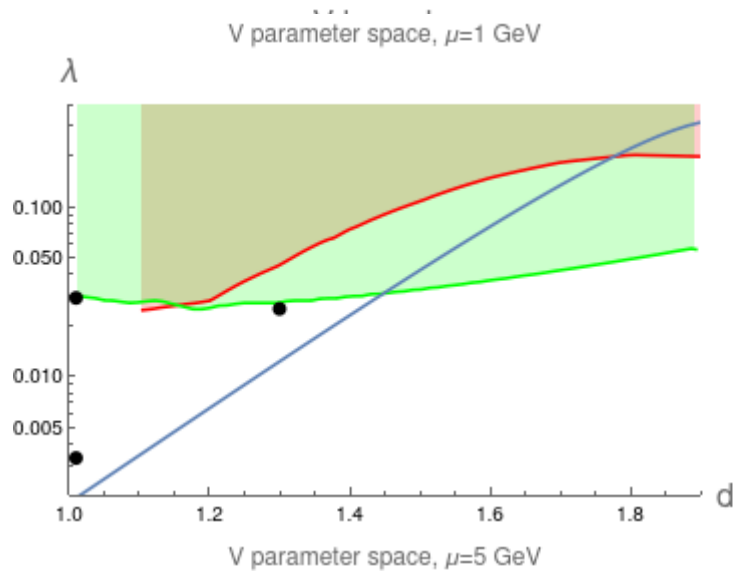
Sensitivity curves



Sensitivity curves



Sensitivity curves



Unparticle effects on the measurement of a_μ^{had}

| | SM | Axial-vector | Vector |
|------|----------------------------|-----------------------------|-----------------------------|
| P1 | $6903(29) \times 10^{-11}$ | $6957(29) \times 10^{-11}$ | $6986(29) \times 10^{-11}$ |
| Pull | 0 | 1.9 | 2.8 |
| P2 | — | $6980(29) \times 10^{-11}$ | $7019(29) \times 10^{-11}$ |
| Pull | — | 2.6 | 4.0 |
| P3 | — | $11073(29) \times 10^{-11}$ | $13250(29) \times 10^{-11}$ |
| Pull | — | 143 | 218 |
| P4 | — | $6954(29) \times 10^{-11}$ | $6979(29) \times 10^{-11}$ |
| Pull | — | 1.7 | 2.6 |
| P5 | — | $6971(29) \times 10^{-11}$ | $7006(29) \times 10^{-11}$ |
| Pull | — | 2.3 | 3.5 |
| P6 | — | $7091(29) \times 10^{-11}$ | $7186(29) \times 10^{-11}$ |
| Pull | — | 6.4 | 9.7 |

$$Pull = (a^{had,new} - a^{had,SM}) / \Delta_{SM}$$

Summary and outlook

- MUonE promise a novel approach of evaluating the hadronic contribution to the muon's $g-2$.
- Such a precise experiment can help us to detect small new physics effect such as unparticles.
- Unparticles with broken scale invariance are still possible, but the parameter space is shrinking.
- MUonE is sensitive to (axial-)vector unparticles with $1 < d \leq 1.4$ and $1 \leq \mu \leq 12$ GeV.
- Further works: re-do the constraints (lepton magnetic moments, CMS mono-Z) carefully with $\mu > 0$.

Acknowledgement :

This research is funded by the Vietnam National Foundation for Science and Technology Development (NAFOSTED) under grant number 103.01-2020.17.