

GAMBIT Dark matter

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Outline

- Global Fitting and GAMBIT
- Recent Code updates
- Recent Physics works (primarily focussing on DM):
 - Bounds on Neutrino Masses
 - Thermal WIMPS (Dirac DM EFT)
 - Cosmological decaying ALPs
 - Simplified DM models
 - Electroweakinos + gravitino
- What to look forward to from GAMBIT

BSM Physics is a balancing act

Theories may need to sit in some sweet spot to solve the problem they were created for (e.g. strong cp), whilst also satisfying all existing constraints.

This motivates comprehensive studies, including as many effects as possible.

Flavour constraints

Direct
Detection

LHC
limits

DM
overabundance

Vacuum
stability

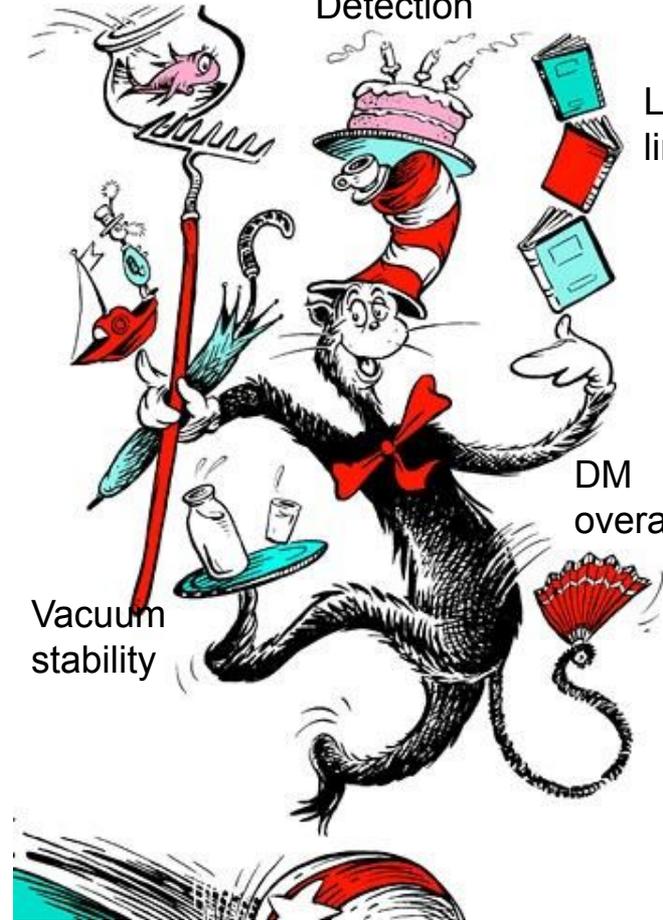


Image credit: Cat in the Hat

Global Fitting

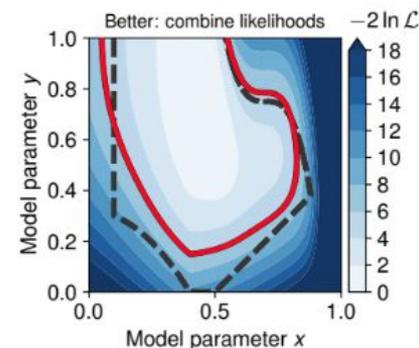
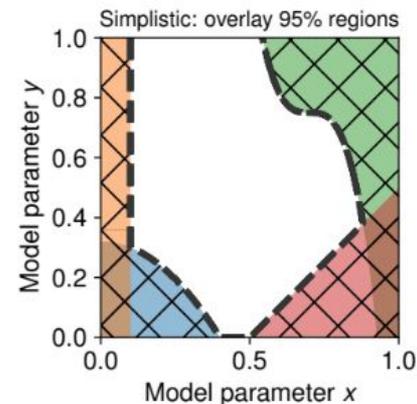
We want to test a theory against a range of experiments over the wide range of possible model parameters.

Need a joint likelihood function:

$$\mathcal{L}_{Total} = \mathcal{L}_{DD} \times \mathcal{L}_{ID} \times \mathcal{L}_{Collider} \times \dots$$

Simple overlays of limits may not be fully informative.

Can be complicated and slow -> needs an efficient way to do this.



GAMBIT: The Global And Modular BSM Inference Tool

gambit.hepforge.org

github.com/GambitBSM

EPJC 77 (2017) 784

arXiv:1705.07908

- Extensive model database, beyond SUSY
- Fast definition of new datasets, theories
- Extensive observable/data libraries
- Plug&play scanning/physics/likelihood packages
- Various statistical options (frequentist /Bayesian)
- Fast LHC likelihood calculator
- Massively parallel
- Fully open-source



Members of: ATLAS, Belle-II, CLIC, CMS, CTA, Fermi-LAT, DARWIN, IceCube, LHCb, SHiP, XENON

Authors of: BubbleProfiler, Capt'n General, Contur, DarkAges, DarkSUSY, DDCalc, DirectDM, Diver, EasyScanHEP, ExoCLASS, FlexibleSUSY, gamLike, GM2Calc, HEPLike, IsaTools, MARTY, nuLike, PhaseTracer, PolyChord, Rivet, SOFTSUSY, SuperIso, SUSY-AI, xsec, Vevacious, WIMPSim

Recent collaborators: P Athron, C Balázs, A Beniwal, S Bloor, T Bringmann, A Buckley, J-E Camargo-Molina, C Chang, M Chruszcz, J Conrad, J Cornell, M Danninger, J Edsjö, T Emken, A Fowlie, T Gonzalo, W Handley, J Harz, S Hoof, F Kahlhoefer, A Kvellestad, P Jackson, D Jacob, C Lin, N Mahmoudi, G Martinez, MT Prim, A Raklev, C Rogan, R Ruiz, P Scott, N Serra, P Stöcker, W. Su, A Vincent, C Weniger, M White, Y Zhang, ++

70+ participants in many experiments and numerous major theory codes

GAMBIT Physics Modules

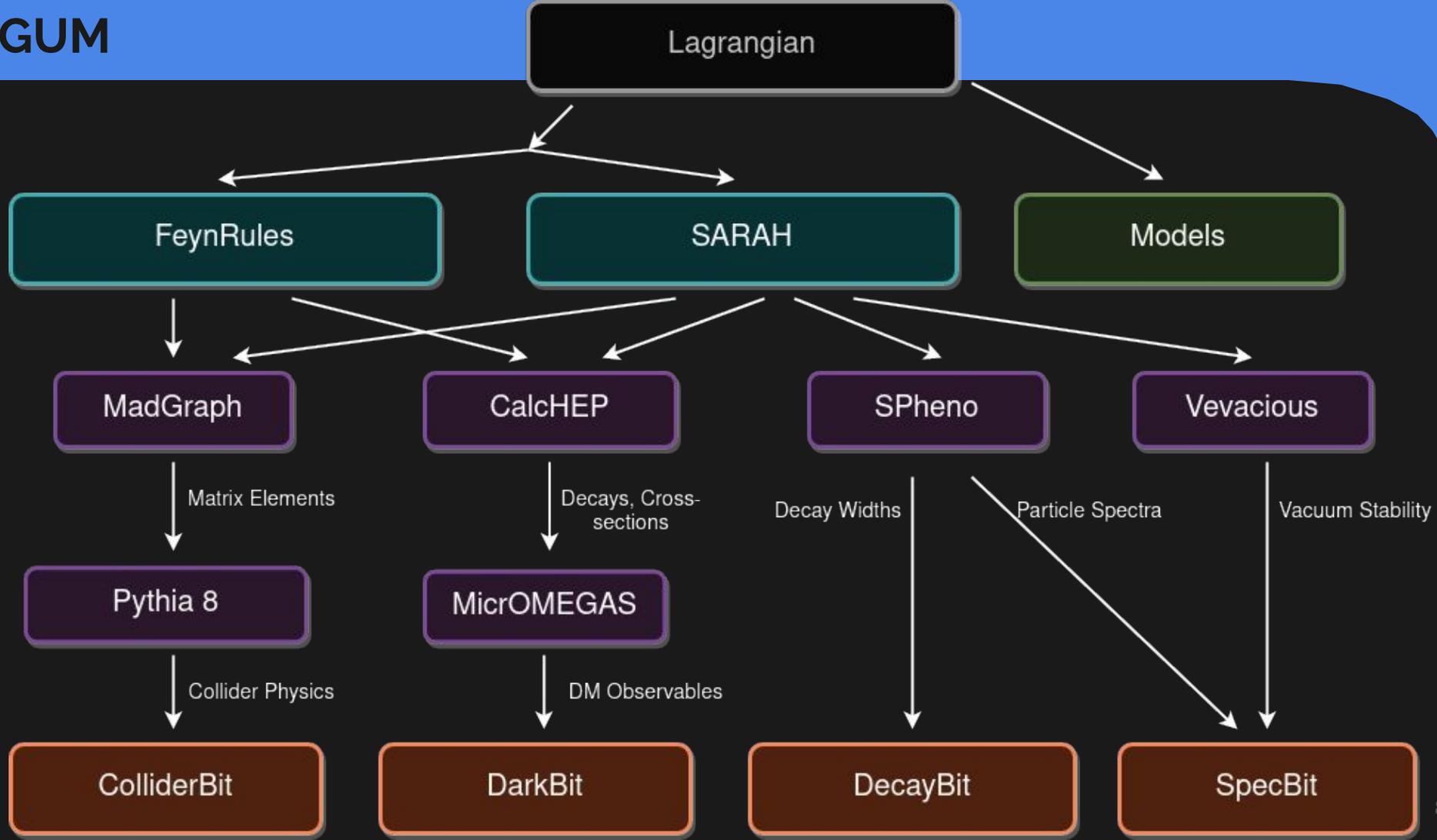
| | |
|-----------------------------------------------------|--------------------|
| ColliderBit: Collider observables | [arXiv:1705.07919] |
| DarkBit: Relic Density, Direct detection,... | [arXiv:1705.07920] |
| FlavBit: Flavour observables | [arXiv:1705.07933] |
| SpecBit: Computes particle masses | [arXiv:1705.07936] |
| DecayBit: Particle decay rates | [arXiv:1705.07936] |
| PrecisionBit: tests based on rare processes | [arXiv:1705.07936] |
| NeutrinoBit: neutrino physics | [arXiv:1908.02302] |
| CosmoBit: cosmology | [arXiv:2009.03286] |

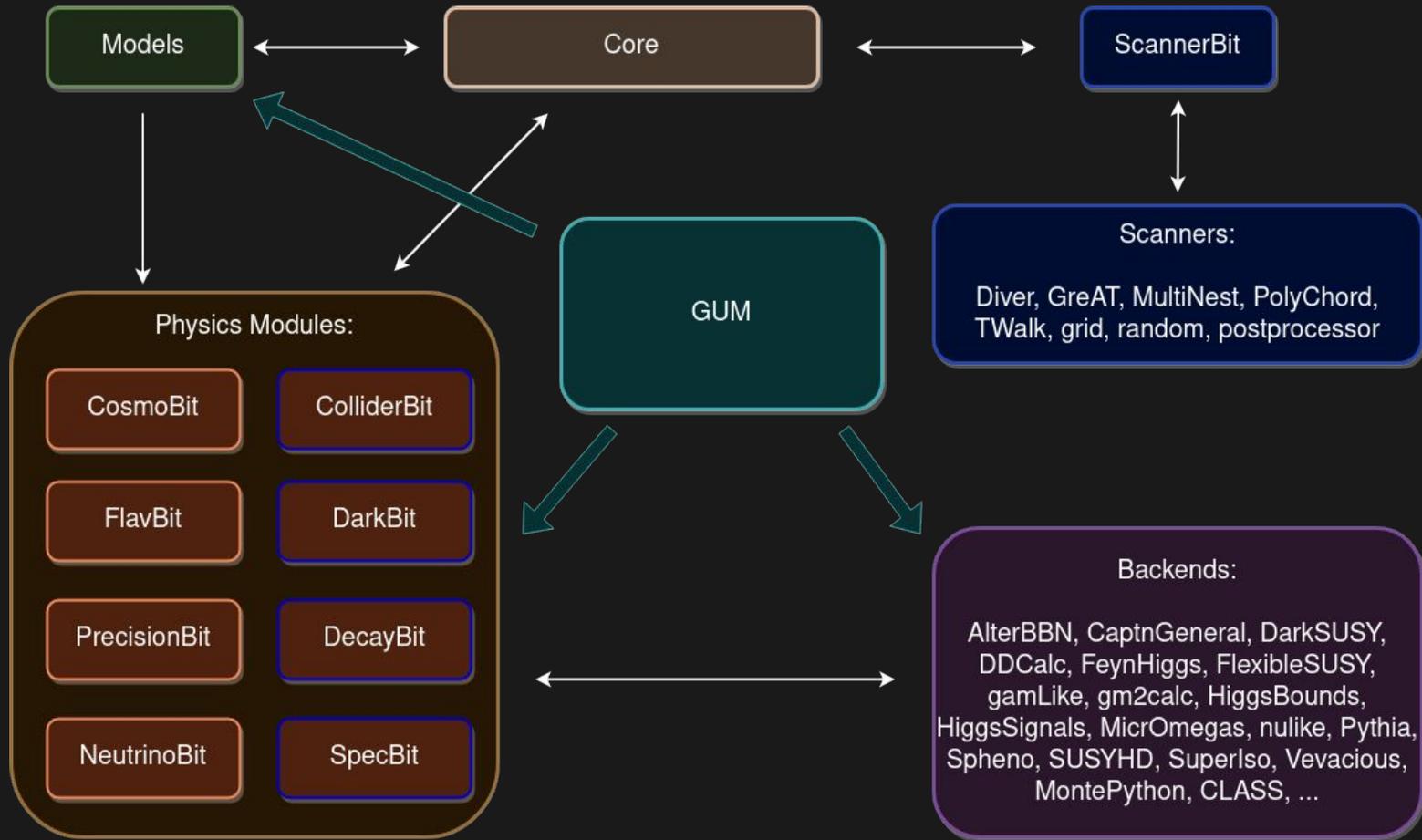
For details on most likelihood implementations, see the relevant GAMBIT module code papers.

Recent code updates

- GUM: The GAMBIT Model Machine: from Lagrangians to Likelihoods
- Many new backend interfaces, including:
 - > SM Measurements (Contur, Rivet)
 - > Cosmology (AlterBBN, ACROPOLIS, MontePython)
 - > EFT matching (DirectDM)
 - > Vacuum Stability (Vevacious)
- Support for ATLAS full likelihoods
- Full support for Clang and AppleClang
- Many quality of life core changes







GAMBIT results relevant to DM

EFTS (Scalar Singlet, simplified models, ...):

- arXiv:1705.07931
- arXiv:1806.11281
- arXiv:1808.10465
- arXiv:2106.02056
- arXiv:2209.13266

Axions/ALPS:

- arXiv:1810.07192
- arXiv:2007.05517
- arXiv:2205.13549

SUSY (constrained SUSY, low-dim SUSY):

- arXiv:1705.07935
- arXiv:1705.07917
- arXiv:1809.02097

Neutrinos (SM Neutrinos, Right-handed):

- arXiv:1908.02302
- arXiv:2009.03287

Code Papers:

- arXiv:1705.07920: DarkBit
- arXiv:2107.00030: CosmoBit
- arXiv:2009.03286: GUM

I will focus on...

- arxiv:2009.03287: Bounds on the lightest neutrino mass
- arxiv:2106.02056: Thermal WIMPs and the Scale of New Physics
- arxiv:2205.13549: Cosmological Constraints on decaying ALPs
- arxiv:2209.13266: s-channel Simplified DM Models
- EW-MSSM + light gravitino \implies Both the gravitino and NLSP would contribute to the DM relic abundance, but we focus only on collider constraints.

Thermal WIMPs and the scale of new physics: arXiv:2106.02056

Sum of all dimension 6 & 7 operators for a Dirac fermion EFT valid up to some new physics scale Λ .

$$\mathcal{L}_{\text{int}} = \sum_{a,d} \frac{\mathcal{C}_a^{(d)}}{\Lambda^{d-4}} \mathcal{Q}_a^{(d)}$$

Include couplings to SM quarks, gluons and photon.

$$\begin{aligned} \mathcal{Q}_{1,q}^{(6)} &= (\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu q), \\ \mathcal{Q}_{2,q}^{(6)} &= (\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu q), \\ \mathcal{Q}_{3,q}^{(6)} &= (\bar{\chi}\gamma_\mu\chi)(\bar{q}\gamma^\mu\gamma_5 q), \\ \mathcal{Q}_{4,q}^{(6)} &= (\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{q}\gamma^\mu\gamma_5 q). \end{aligned}$$

$$\begin{aligned} \mathcal{Q}_1^{(7)} &= \frac{\alpha_s}{12\pi} (\bar{\chi}\chi) G^{a\mu\nu} G_{\mu\nu}^a, \\ \mathcal{Q}_2^{(7)} &= \frac{\alpha_s}{12\pi} (\bar{\chi}i\gamma_5\chi) G^{a\mu\nu} G_{\mu\nu}^a, \\ \mathcal{Q}_3^{(7)} &= \frac{\alpha_s}{8\pi} (\bar{\chi}\chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a, \\ \mathcal{Q}_4^{(7)} &= \frac{\alpha_s}{8\pi} (\bar{\chi}i\gamma_5\chi) G^{a\mu\nu} \tilde{G}_{\mu\nu}^a, \\ \mathcal{Q}_{5,q}^{(7)} &= m_q (\bar{\chi}\chi)(\bar{q}q), \\ \mathcal{Q}_{6,q}^{(7)} &= m_q (\bar{\chi}i\gamma_5\chi)(\bar{q}q), \\ \mathcal{Q}_{7,q}^{(7)} &= m_q (\bar{\chi}\chi)(\bar{q}i\gamma_5 q), \\ \mathcal{Q}_{8,q}^{(7)} &= m_q (\bar{\chi}i\gamma_5\chi)(\bar{q}i\gamma_5 q), \\ \mathcal{Q}_{9,q}^{(7)} &= m_q (\bar{\chi}\sigma^{\mu\nu}\chi)(\bar{q}\sigma_{\mu\nu} q), \\ \mathcal{Q}_{10,q}^{(7)} &= m_q (\bar{\chi}i\sigma^{\mu\nu}\gamma_5\chi)(\bar{q}\sigma_{\mu\nu} q). \end{aligned}$$

Thermal WIMPs and the scale of new physics: arXiv:2106.02056

EFT is valid if $\Lambda >$ largest pT bin in analysis.

We use a falling missing energy spectra above Λ .

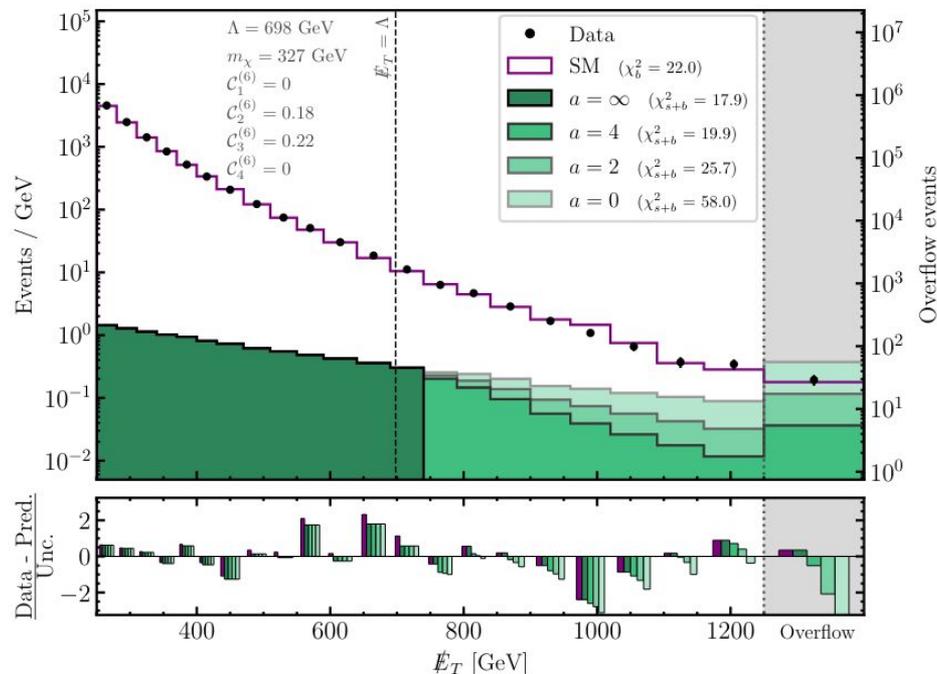
The nature of this signal drop off would depend on the UV completion.

$$\frac{d\sigma}{d\cancel{E}_T} \rightarrow \frac{d\sigma}{d\cancel{E}_T} \left(\frac{\cancel{E}_T}{\Lambda} \right)^{-a}$$

Two cases:

1. Hard cutoff ($a = \infty$)
2. Scanning over a range $[0,4]$ to find the best fit.

I will only show the results for the **hard cutoff**

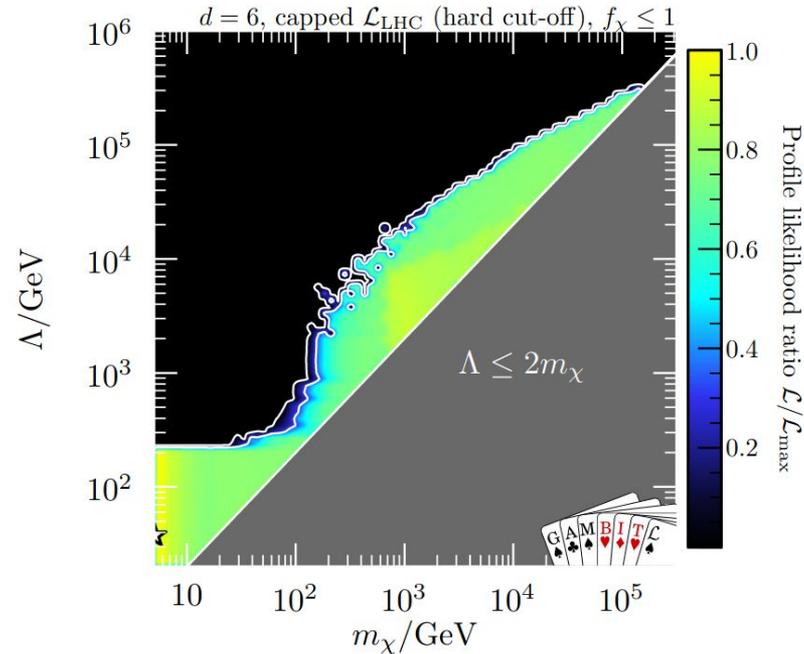


Thermal WIMPs and the scale of new physics: arXiv:2106.02056

High Λ high m_χ : Relic density exceeds planck measurement unless couplings become non-perturbative

High Λ low m_χ : Strong LHC constraints prevent any surviving regions that satisfy relic abundance.

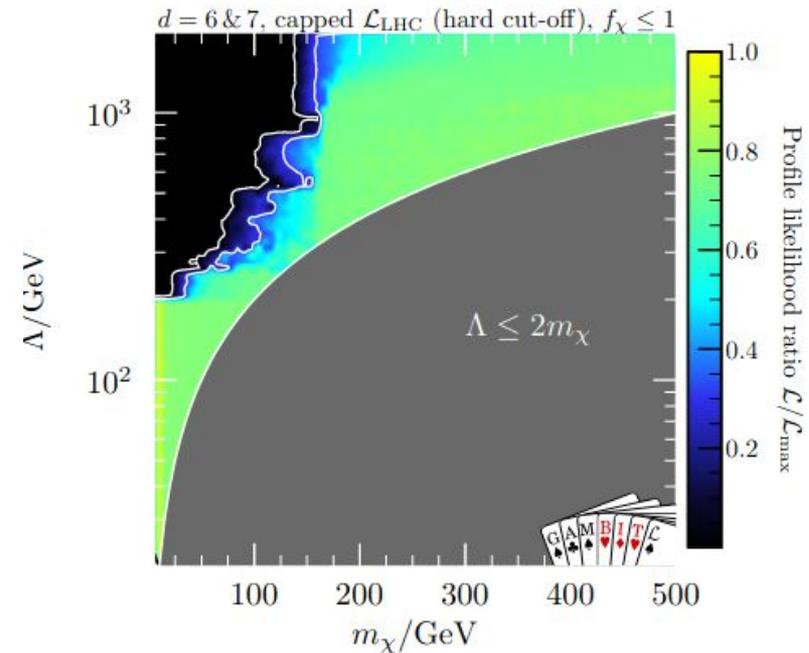
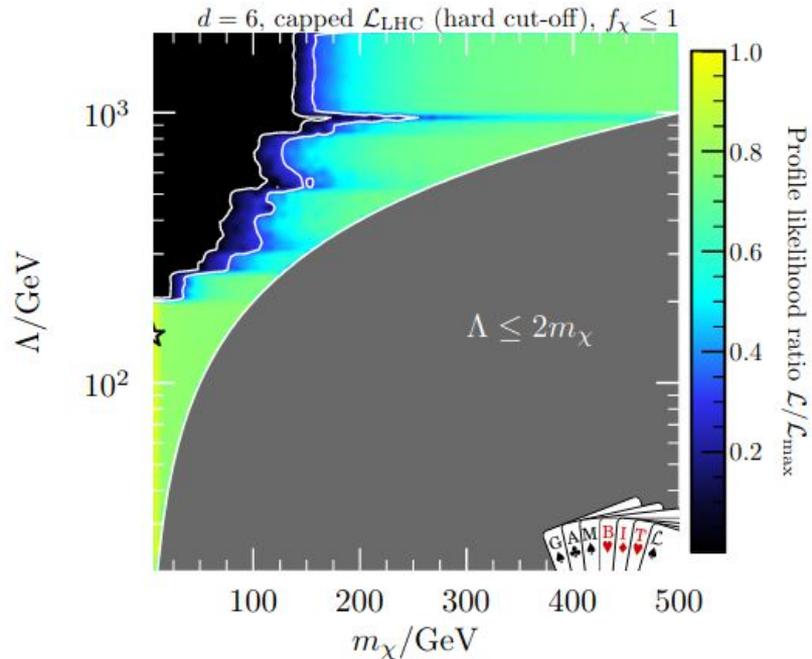
Below $\Lambda = 200$ GeV, LHC has no impact given EFT validity cutoff.



Thermal WIMPs and the scale of new physics: arXiv:2106.02056

The addition of the dimension 7 operators has very little impact in the m_χ/Λ plane.

Surviving parameter regions change abruptly with Λ , whenever additional signal regions enter the calculation.



Cosmological constraints on decaying ALPs: arXiv:2205.13549

Focus on Axion-Like-Particles (ALPs) in the keV - MeV mass range, and lifetimes: $10^4 \text{ s} < \tau_a < 10^{13} \text{ s}$.

Corresponds to ALPs that decay sometime between the end of BBN and CMB.

ALP-photon effective coupling:

$$\mathcal{L} = \frac{g_{a\gamma}}{4} a F^{\mu\nu} \tilde{F}_{\mu\nu}$$

4 ALP free parameters: ALP mass, ALP lifetime, ALP-photon coupling and ALP abundance (ξ)

$$\xi = \frac{m_a n_{a,0}^{\tau_a \rightarrow \infty}}{\rho_{\text{DM},0}}$$

ξ cannot be arbitrarily small, due to the irreducible contribution from freeze-in.

Cosmological constraints on decaying ALPs: arXiv:2205.13549

Constraints from energy injection:

- CMB power spectra
- Modifications to N_{eff} & photon-baryon ratio (η_b)
- Primordial Nucleosynthesis (light element photo-disintegration)
- ALP induced spectral distortions
- ALP decays after SN1987A
- Stellar Evolution (ratio of HB/RGB stars)

| Model parameter | | Scan range | |
|-----------------------------|----------------------|--------------------|------------------------------------|
| ALP mass | m_a | [0.001, 200] | MeV |
| ALP lifetime | τ_a | $[10^4, 10^{13}]$ | s |
| ALP abundance | ξ | $[10^{-12}, 10^2]$ | |
| Baryon abundance | ω_b | [0.020, 0.024] | |
| Dark matter abundance | ω_{DM} | [0.10, 0.13] | |
| Hubble constant | H_0 | [62, 74] | $\text{km s}^{-1} \text{Mpc}^{-1}$ |
| Redshift of reionisation | z_{reio} | [4.5, 9.5] | |
| Primordial curvature | $\ln(10^{10} A_s)$ | [2.9, 3.2] | |
| Scalar spectral index | n_s | [0.9, 1.1] | |
| Neutron lifetime | τ_n | [875, 895] | s |
| Planck nuisance parameter | A_{Planck} | [0.9, 1.1] | |
| Pantheon nuisance parameter | M | [-20, -18] | |

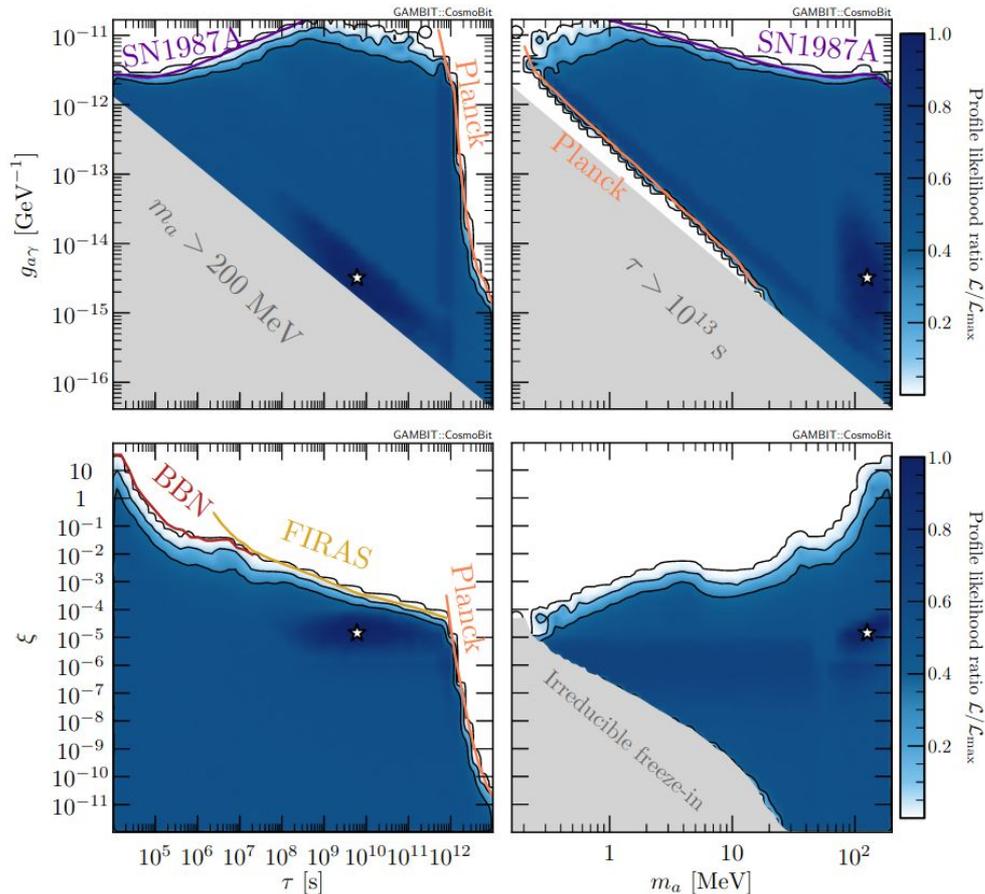
Cosmological constraints on decaying ALPs: arXiv:2205.13549

We find $m_a > 300$ keV, which can only be evaded if ALPS are stable on cosmological timescales.

Viable regions exist that could constitute a high ALP abundance in the very early universe, but decay away before recombination.

No strong effect from N_{eff} , η_b or R parameter.

Slight preference at best-fit point is due to additional deuterium abundance from photodisintegration.



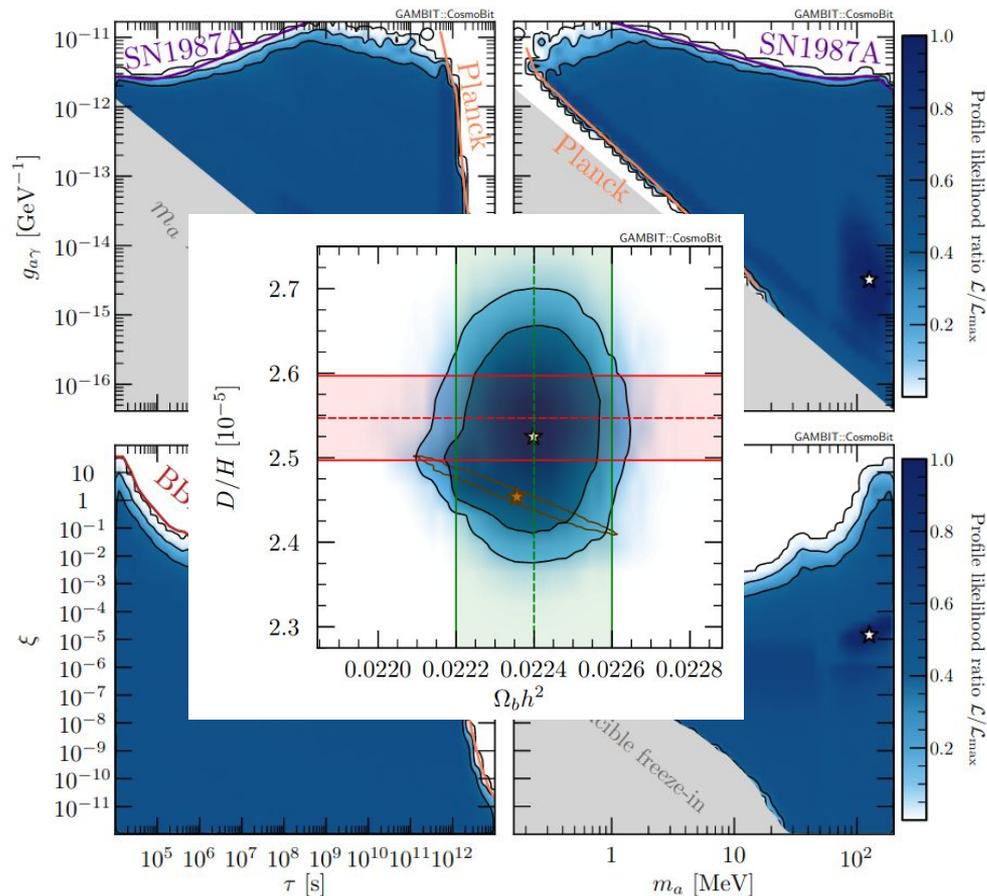
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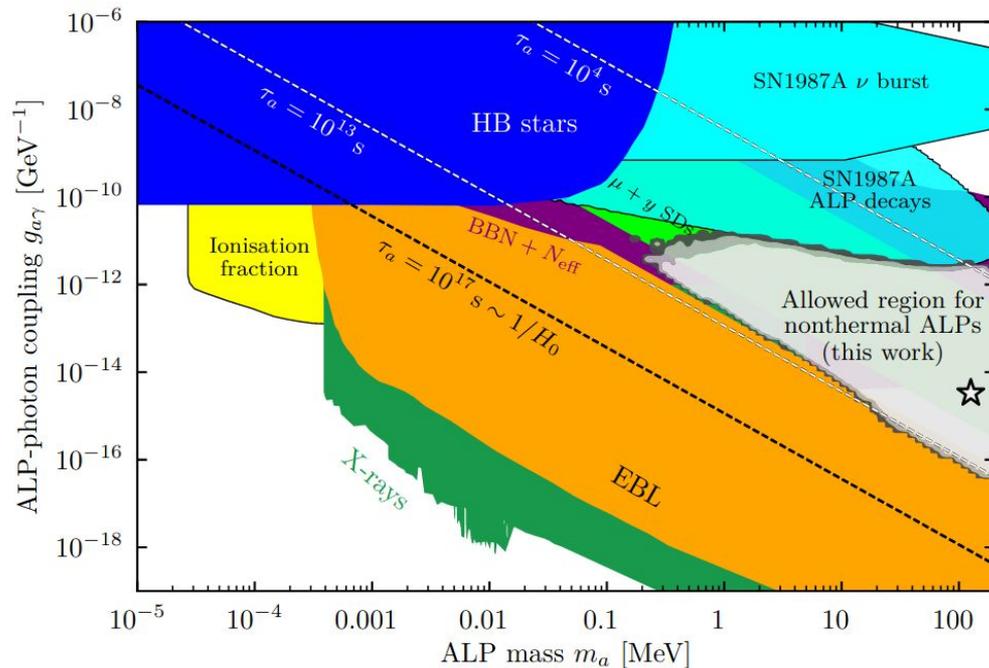
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S-channel Simplified Models: arXiv:2209.13266

Scalar DM:

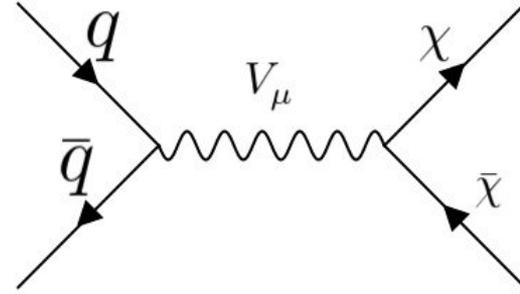
$$\mathcal{L}_{BSM} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_{DM}^2 \phi^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_M^2 V_\mu V^\mu + g_q V_\mu \bar{q} \gamma^\mu q + i g_{DM}^V V_\mu (\phi^\dagger (\partial^\mu \phi) - (\partial^\mu \phi^\dagger) \phi)$$

Dirac fermion DM:

$$\mathcal{L}_{BSM} = i \bar{\chi} \gamma^\mu \partial_\mu \chi - m_{DM} \bar{\chi} \chi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_M^2 V_\mu V^\mu + g_q V_\mu \bar{q} \gamma^\mu q + V_\mu \bar{\chi} (g_{DM}^V + g_{DM}^A \gamma^5) \gamma^\mu \chi$$

Majorana fermion DM:

$$\mathcal{L}_{BSM} = \frac{1}{2} i \bar{\psi} \gamma^\mu \partial_\mu \psi - \frac{1}{2} m_{DM} \bar{\psi} \psi - \frac{1}{4} V_{\mu\nu} V^{\mu\nu} - \frac{1}{2} m_M^2 V_\mu V^\mu + g_q V_\mu \bar{q} \gamma^\mu q + \frac{1}{2} g_{DM}^A V_\mu \bar{\psi} \gamma^5 \gamma^\mu \psi$$



Assumptions:

No lepton couplings

No axial-vector quark couplings

Flavour universal couplings

In each model, there are 4 or 5 model parameters: DM mass (m_{DM}), Mediator mass (m_M), mediator-quark coupling (g_q), mediator-DM coupling (g_{DM}) (either vector or axial-vector)

S-channel Simplified Models: arXiv:2209.13266

The presence of an axial-vector couplings for the Dirac and Majorana models implies a bound from unitarity: [3]

$$m_{DM} \leq \sqrt{\frac{\pi}{2}} \frac{m_M}{g_{DM}^A}$$

Vector DM models will face strong unitarity violation -> An upcoming paper will go into detail about this violation.

S-channel Simplified Models: arXiv:2209.13266

Scan Parameters:

| Parameters | Range |
|-------------------------------------------------|----------------------------------|
| DM mass, m_{DM} | [50, 10000] GeV |
| Mediator mass, m_M | [50, 10000] GeV |
| quark-mediator coupling, g_q | [0.01, 1.0] |
| mediator-DM coupling (vector), g_{DM}^V | [0.01, 3.0] |
| mediator-DM coupling (axial vector), g_{DM}^A | [0.01, 3.0] |
| Nuisance Parameters | |
| Pion-nucleon sigma term, $\sigma_{\pi N}$ | [5, 95] MeV |
| strange quark cont. to nucleon spin, Δ_s | [-0.062, -0.008] |
| strange quark nuclear tensor charge, g_T^s | [-0.075, 0.021] |
| strange quark proton charge radius, r_s^2 | [-0.22, -0.01] GeV ⁻² |
| Local DM density, ρ_0 | [0.2, 0.8] GeV cm ⁻³ |
| Most probably speed, v_{esc} | [216, 264] km s ⁻¹ |
| Galactic escape speed, v_{peak} | [453, 603] km s ⁻¹ |

Constraints:

Experiment

CDMSlite [4]
CRESST-II [5]
CRESST-III [6]
DarkSide 50 [7]
LUX 2016 [8]
PICO-60 [9, 10]
PandaX [11, 12]
XENON1T [13]
LZ 2022 [28]

LHC Dijets [14–22]
ATLAS monojet [23]
CMS monojet [24]

Fermi-LAT [25]
Planck 2018: Ωh^2 [26]

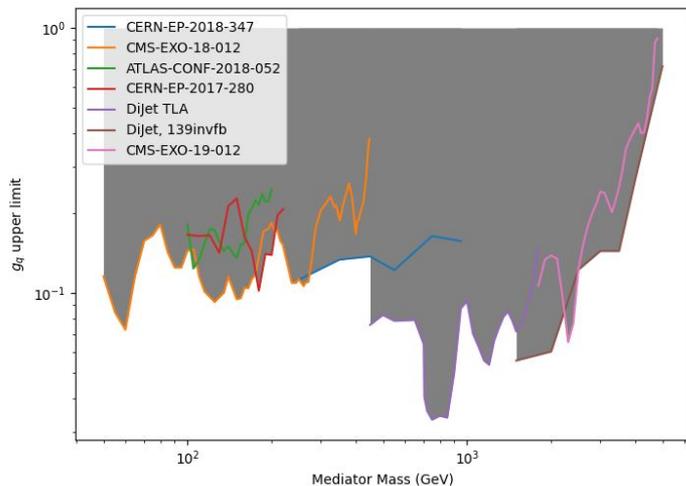
Nuisances

S-channel Simplified Models: arXiv:2209.13266

Direct Detection: Majorana model only has NREOs that lead to suppressed signals.

| Effective Operator | Relevant models |
|---------------------------------------------------------------------------------------------------------------------------------------|-----------------|
| $1_{DM}1_N$ | Scalar, Dirac |
| $i\hat{\mathbf{S}} \cdot (\hat{\mathbf{S}}_N \times \frac{\hat{\mathbf{q}}}{m_N}), \hat{\mathbf{S}} \cdot \hat{\mathbf{v}}^\perp 1_N$ | Dirac, Majorana |

Di-jets:

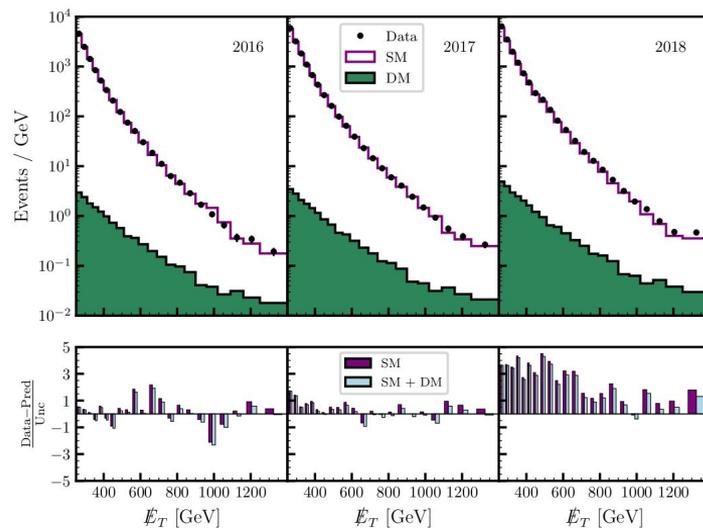


DM Abundance:

2 Annihilation channels:

- DM DM \rightarrow quark pair
- DM DM \rightarrow mediator pair

Collider Missing Energy searches:

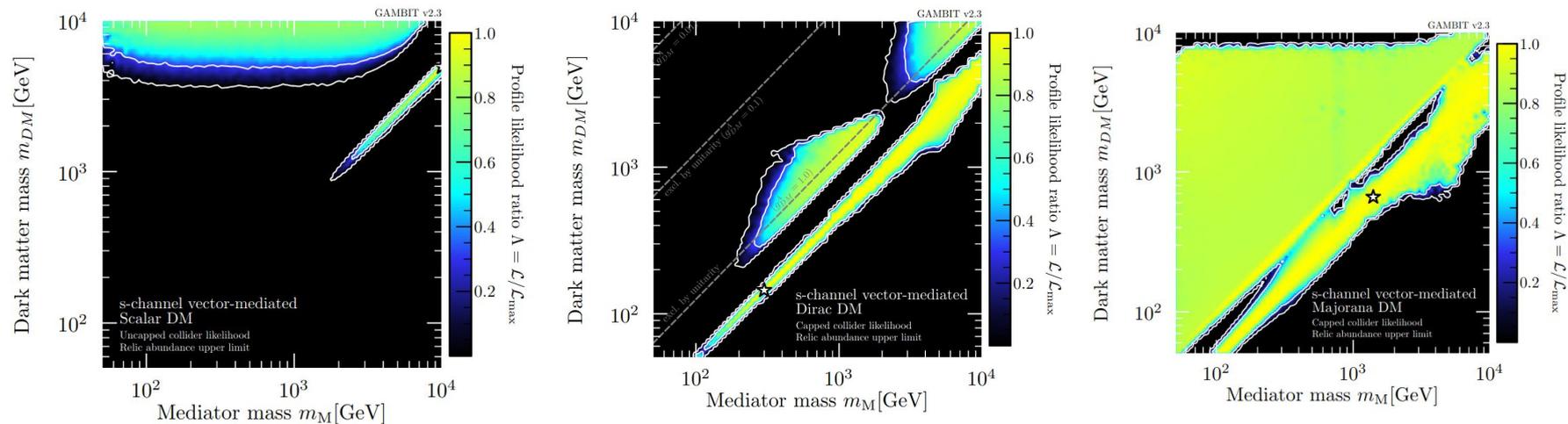


S-channel Simplified Models: arXiv:2209.13266

On/off-resonance regions are generally split by the annihilation channels (qq/MM).

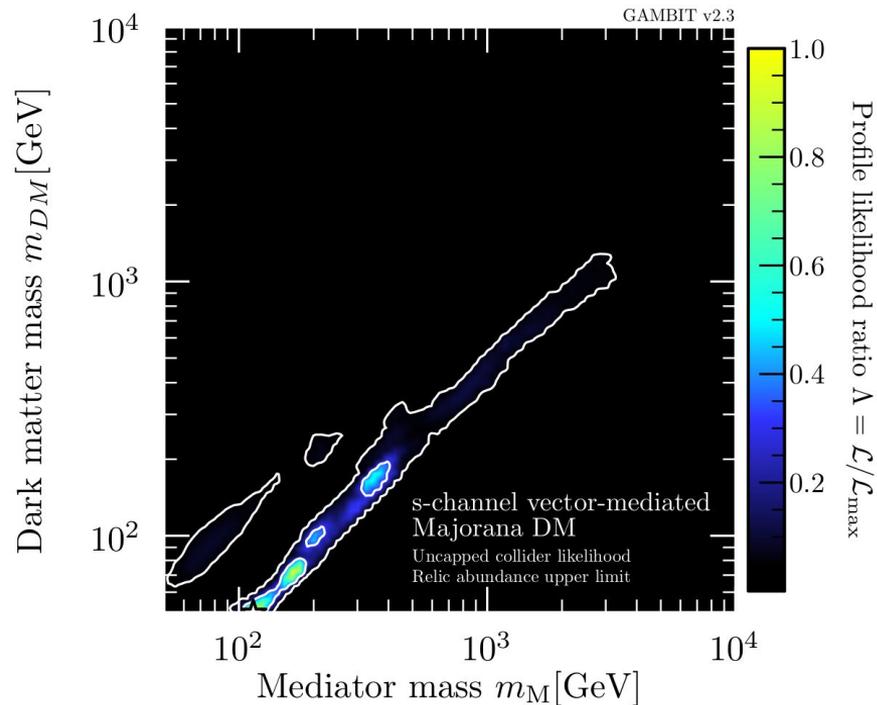
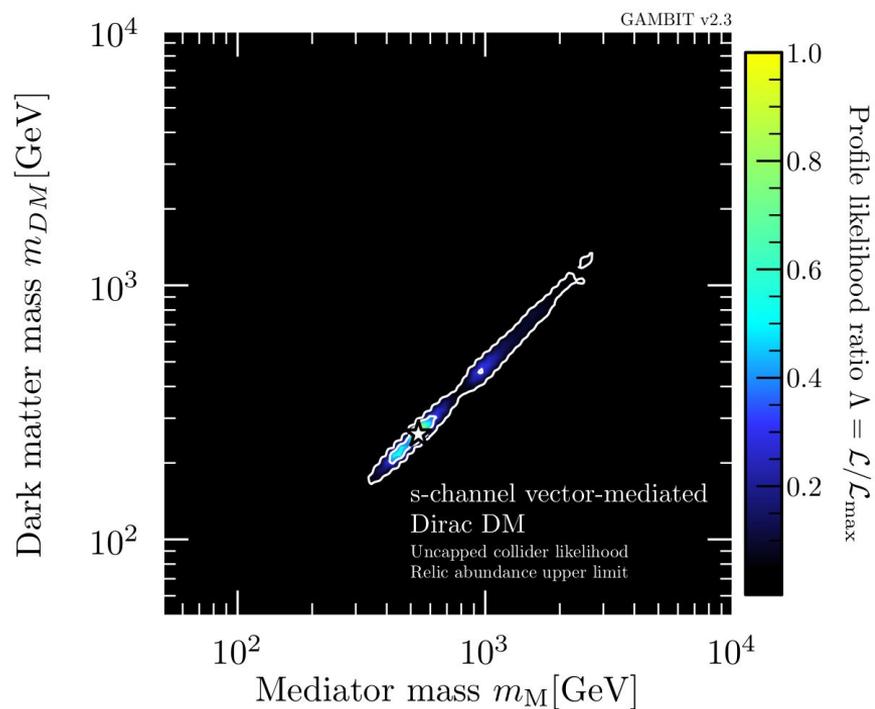
Dirac DM: Combination of unitarity violation and direct detection limits give two distinct regions off resonance.

Majorana DM: Weak exclusion when DM mass > Mediator mass due to suppressed NREOs.



S-channel Simplified Models: arXiv:2209.13266

By fitting the fluctuations in the monojet data, the Dirac and Majorana DM models have an overall preference for a small regions primarily along the resonance.



Bounds on the lightest neutrino mass: arXiv:2009.03287

Use both terrestrial and cosmological experiments to set a CL on the lightest neutrino mass and the sum of neutrino masses.

Constraints:

Neutrino oscillations (chi-squared tables from NuFit)

Primordial BBN abundances

CMB

Supernova Type 1a

BAO

$$\text{NH} : (m_1^2, m_2^2, m_3^2) = m_{\nu_0}^2 + (0, \Delta m_{21}^2, \Delta m_{3l}^2)$$

$$\text{IH} : (m_3^2, m_1^2, m_2^2) = m_{\nu_0}^2 + (0, |\Delta m_{3l}^2| - \Delta m_{21}^2, |\Delta m_{3l}^2|).$$

| Sector | Parameter | Range |
|------------------------|--------------------------------------------|-----------------------------------------------|
| ν masses | m_{ν_0} | [0, 1.1] eV |
| | Δm_{21}^2 | $[6, 9] \times 10^{-5} \text{ eV}^2$ |
| | (NH) Δm_{3l}^2 | $[2.2, 2.8] \times 10^{-3} \text{ eV}^2$ |
| (IH) Δm_{3l}^2 | $[-2.8, -2.2] \times 10^{-3} \text{ eV}^2$ | |
| Λ CDM | H_0 | $[50, 80] \text{ km s}^{-1} \text{ Mpc}^{-1}$ |
| | $\Omega_b h^2$ | [0.020, 0.024] |
| | $\Omega_{\text{cdm}} h^2$ | [0.10, 0.15] |
| | $\tau_{\text{reionization}}$ | [0.004, 0.20] |
| | $\ln(10^{10} A_s)$ | [2.5, 3.5] |
| | n_s | [0.90, 1.10] |
| N_{eff} | r_ν | [0.75, 1.15] |
| Nuisance | SN Ia abs. magnitude M | [-20, -18] |
| | Neutron lifetime τ_n | [876, 883] s |
| | <i>Planck</i> likelihood | 21 parameters varied |

Bounds on the lightest neutrino mass: arXiv:2009.03287

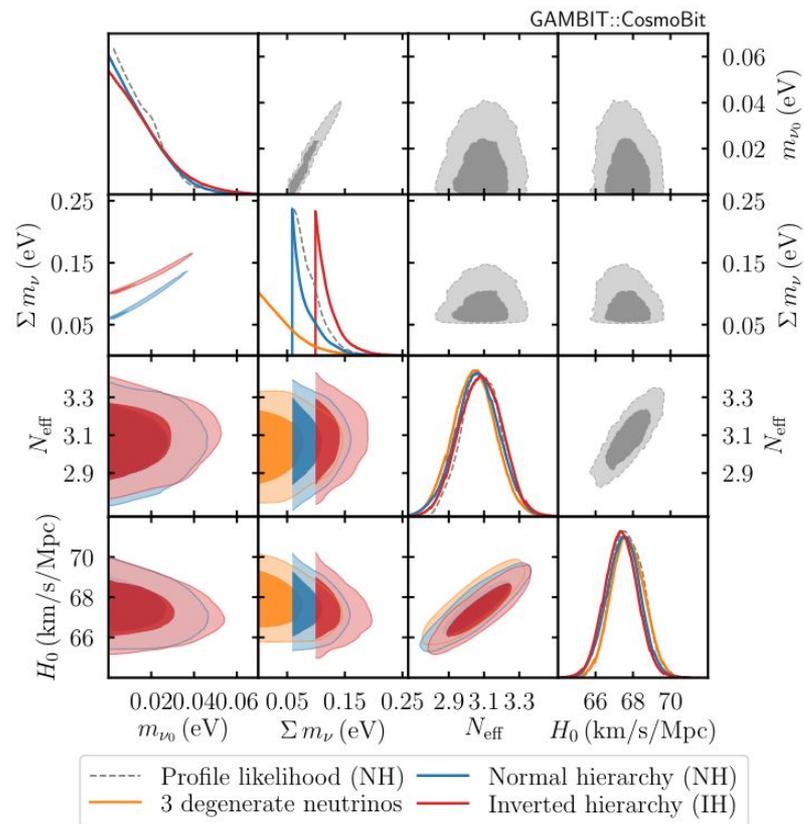
$$m_{\nu 0} < 0.037 \text{ eV} \quad (\text{NH})$$

$$m_{\nu 0} < 0.042 \text{ eV} \quad (\text{IH})$$

$$0.058 < \sum m_\nu < 0.139 \text{ eV} \quad (\text{NH})$$

$$0.098 < \sum m_\nu < 0.174 \text{ eV} \quad (\text{IH})$$

Highest likelihood corresponds to a massless lightest neutrino



Electroweakinos + light gravitino

Electroweakino effective field theory + light (1 eV) gravitino LSP (rest of SUSY decoupled).

The gravitino (and NLSP) would contribute to the DM relic abundance, but we don't include these constraints (collider focussed study).

Main collider signal: production + decay of light EWinos

15 ATLAS + 12 CMS searches for new physics, and collection of SM measurements.

Paper to be released soon.

| Parameter | Range/value | Sampling priors |
|----------------------------------------|---------------|-----------------|
| $M_1(Q)$ | $[-1, 1]$ TeV | hybrid, flat |
| $M_2(Q)$ | $[0, 1]$ TeV | hybrid, flat |
| $\mu(Q)$ | $[-1, 1]$ TeV | hybrid, flat |
| $\tan\beta(m_Z)$ | $[1, 70]$ | log, flat |
| $m_{3/2}$ | 1 eV | fixed |
| Q | 3 TeV | fixed |
| $\alpha_s^{\overline{\text{MS}}}(m_Z)$ | 0.1181 | fixed |
| Top quark pole mass | 171.06 GeV | fixed |
| Higgs mass | 125.09 GeV | fixed |

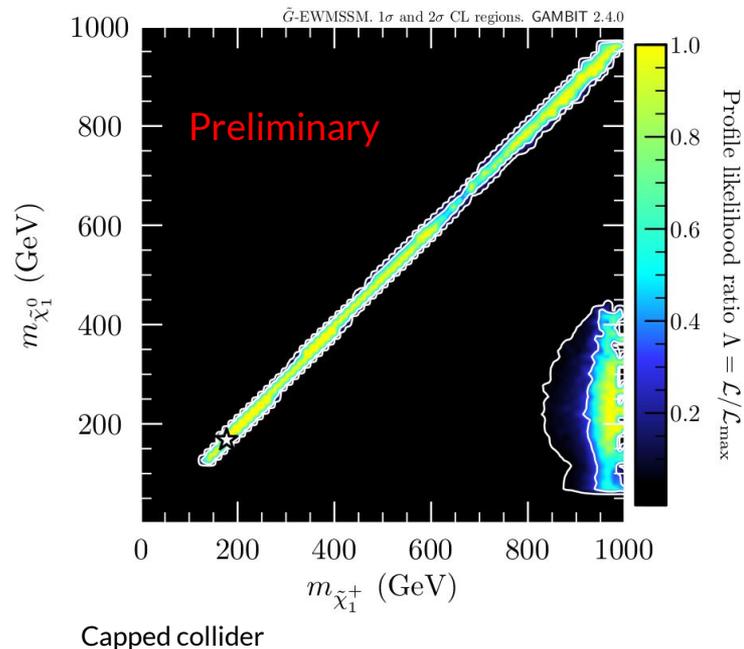
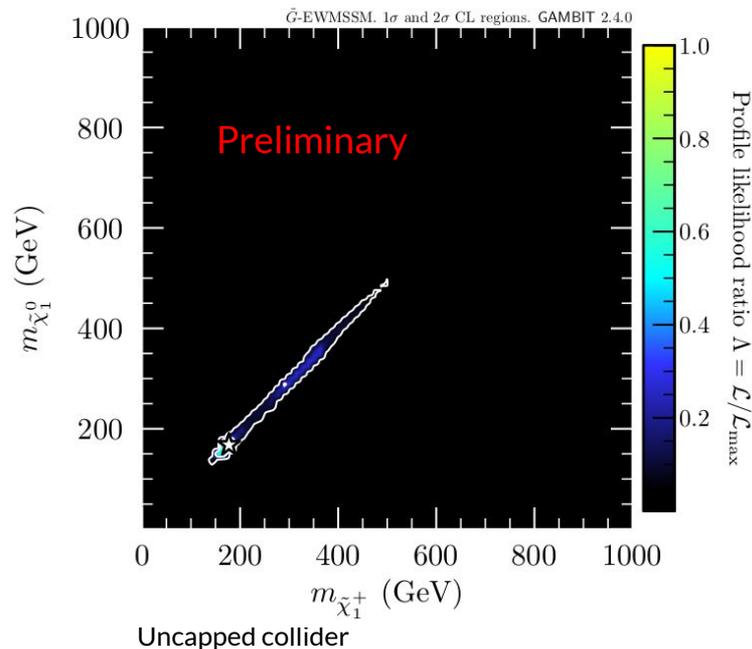
| Search label | Luminosity | Source |
|--------------------------|----------------------|----------------------------------------------------------------|
| ATLAS_2BoostedBosons | 139 fb ⁻¹ | ATLAS hadronic chargino/neutralino search [102] |
| ATLAS_0lep | 139 fb ⁻¹ | ATLAS 0-lepton search [94] |
| ATLAS_0lep_stop | 36 fb ⁻¹ | ATLAS 0-lepton stop search [103] |
| ATLAS_1lep_stop | 36 fb ⁻¹ | ATLAS 1-lepton stop search [104] |
| ATLAS_2lep_stop | 139 fb ⁻¹ | ATLAS 2-lepton stop search [105] |
| ATLAS_2OSlep_Z | 139 fb ⁻¹ | ATLAS stop search with Z/H final states [107] |
| ATLAS_2OSlep_chargino | 139 fb ⁻¹ | ATLAS 2-lepton chargino search [95] |
| ATLAS_2b | 36 fb ⁻¹ | ATLAS 2- <i>b</i> -jet stop/sbottom search [108] |
| ATLAS_3b | 24 fb ⁻¹ | ATLAS 3- <i>b</i> -jet Higgsino search [109] |
| ATLAS_3lep | 139 fb ⁻¹ | ATLAS 3-lepton chargino/neutralino search [96] |
| ATLAS_4lep | 139 fb ⁻¹ | ATLAS 4-lepton search [97] |
| ATLAS_MultiLep_strong | 139 fb ⁻¹ | ATLAS leptons + jets search [98] |
| ATLAS_PhotonGGM_1photon | 139 fb ⁻¹ | ATLAS 1-photon GGM search [110] |
| ATLAS_PhotonGGM_2photon | 36 fb ⁻¹ | ATLAS 2-photon GGM search [111] |
| ATLAS_Z_photon | 80 fb ⁻¹ | ATLAS Z + photon search [112] |
| CMS_0lep | 137 fb ⁻¹ | CMS 0-lepton search [113] |
| CMS_1lep_bb | 36 fb ⁻¹ | CMS 1-lepton + <i>b</i> -jets chargino/neutralino search [115] |
| CMS_1lep_stop | 36 fb ⁻¹ | CMS 1-lepton stop search [116] |
| CMS_2lep_stop | 36 fb ⁻¹ | CMS 2-lepton stop search [117] |
| CMS_2lep_soft | 36 fb ⁻¹ | CMS 2 soft lepton search [118] |
| CMS_2OSlep | 137 fb ⁻¹ | CMS 2-lepton search [119] |
| CMS_2OSlep_chargino_stop | 36 fb ⁻¹ | CMS 2-lepton chargino/stop search [120] |
| CMS_2Slep_stop | 137 fb ⁻¹ | CMS 2 same-sign lepton stop search [121] |
| CMS_MultiLep | 137 fb ⁻¹ | CMS multilepton chargino/neutralino search [100] |
| CMS_photon | 36 fb ⁻¹ | CMS 1-photon GMSB search [122] |
| CMS_2photon | 36 fb ⁻¹ | CMS 2-photon GMSB search [123] |
| CMS_1photon_1lepton | 36 fb ⁻¹ | CMS 1-photon + 1-lepton GMSB search [124] |

Electroweakinos + light gravitino

Most preferred: 3 nearly degenerate higgsinos which can fit small excesses in several searches.

Escapes photon + MET searches with low BR to photons.

Allowed region in capped collider scan in the bottom right is due to diminishing production xsec.



What's coming up from GAMBIT

Code:

Computational speed up for performing global fits:

- > Machine learning of cross sections, cosmic ray fluxes
- > GPU parallelisation of computational bottlenecks

ColliderBit Solo

- > Allows processing collider events generated elsewhere and passed to ColliderBit

Physics:

Light dark matter

- > Sub-GeV DM fits, including Beam dump experiments, direct & indirect detection, cosmology, ...

Vector DM Simplified Models

- > Studying the effect of unitarity violation and its impact on fits

2HDMs

- > Fits of Type 1, Type 2 and inert doublet models

Neutrino Oscillation

- > Fit of left-handed neutrinos reproducing oscillation data with detailed treatment of systematic uncertainties.

Summary

We have GAMBIT studies of many different theories, ranging from SUSY, ALPs, EFTs,...

I've discussed results from a few DM related recent studies.

Stay tuned for more physics results to come.

Website (tutorials, documentation): <https://gambitbsm.github.io>

GitHub: <https://github.com/GambitBSM/gambit>

Code Paper: GAMBIT: The Global and Modular Beyond-the-Standard-Model Inference Tool, Eur. Phys. J. C 77 (2017) 784, Addendum: Eur. Phys. J. C 78 (2018) 98, [arXiv:1705.07908](https://arxiv.org/abs/1705.07908)

GUM Paper: The GAMBIT Universal Model Machine: from Lagrangians to Likelihoods, Eur. Phys. J. C 81 (2021) 12, 1103, [arXiv:2107.00030](https://arxiv.org/abs/2107.00030).

