

# Dark matter **subhalo** **boost** factor

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[based on 2007.10392, 2203.16440, and 2203.16491]

**Gaétan Facchinetti**  
**(Université Libre de Bruxelles)**

in collaboration with  
Thomas Lacroix, Martin Stref,  
Judit Pérez-Romero, Julien Lavalle,  
David Maurin, and Miguel A. Sánchez-Conde





Credit: Tom Gauld  
(for NEW SCIENTIST)

## INTRODUCTION

WHAT IS THE SUBHALO  
BOOST FACTOR?

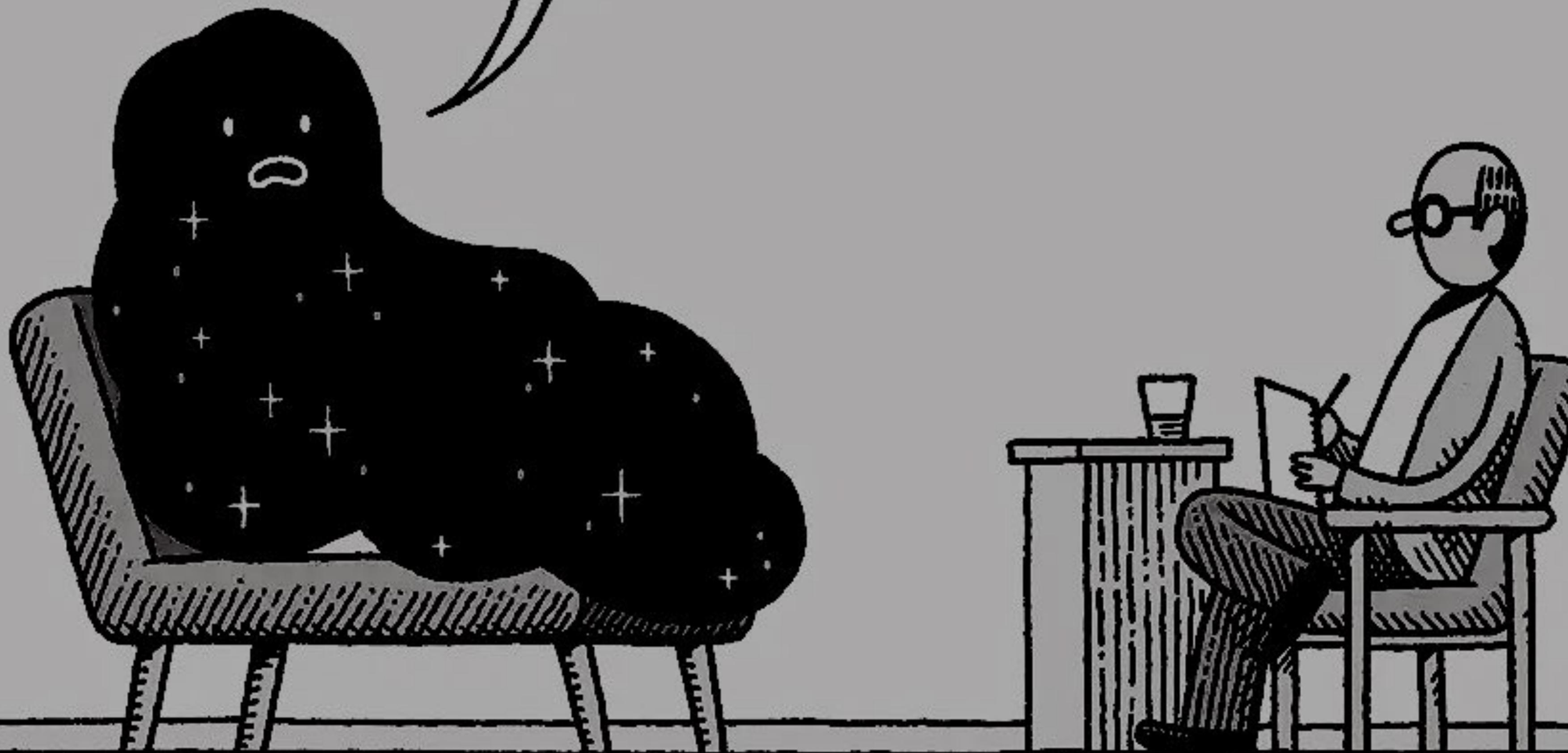
Subhalos,



Subhalos everywhere

**Let us first discuss  
DM structuring**

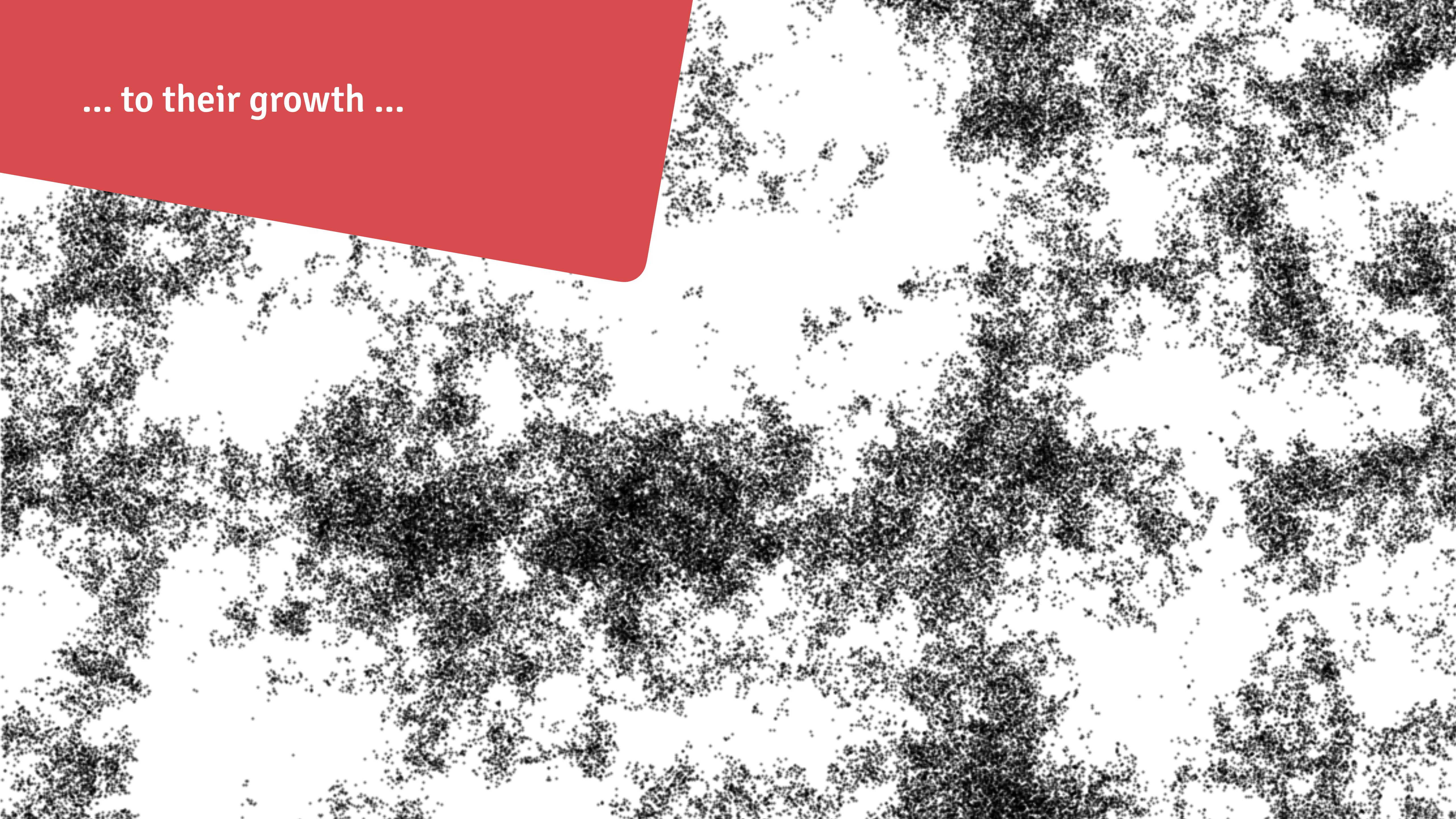
THEY ALL ASK "WHAT IS DARK MATTER?"  
AND "WHERE IS DARK MATTER?", BUT  
NOBODY ASKS "HOW IS DARK MATTER?"



Credit: Tom Gauld  
(for NEW SCIENTIST)



The story starts from  
the initial perturbations  
of the matter density field...



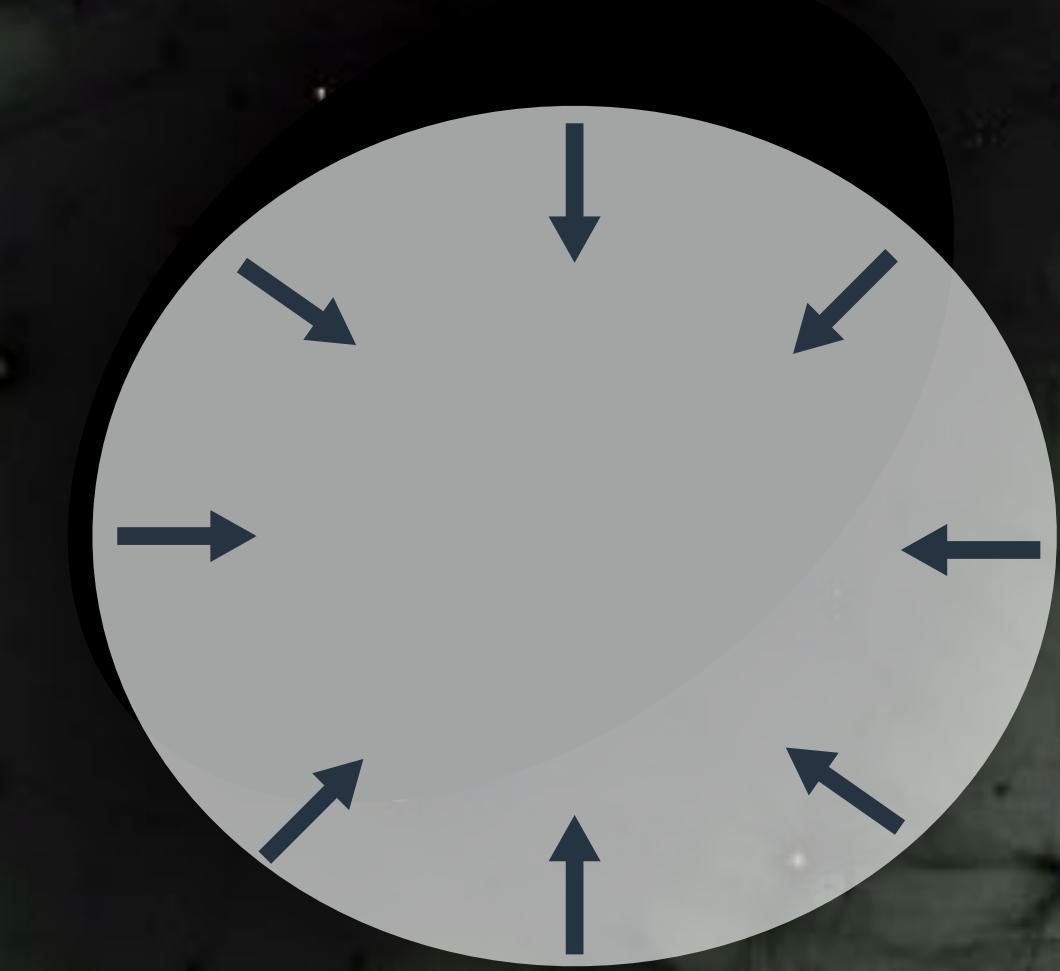
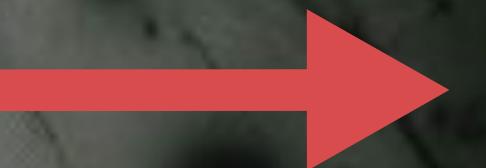
A black and white photograph of a forest. The trees are tall and slender, with dark, textured foliage. The sky above is bright and clear. In the upper left corner, there is a solid red rectangular overlay containing the text.

... to their growth ...

... and their collapse  
into halos (virialisation)



Initial overdensity

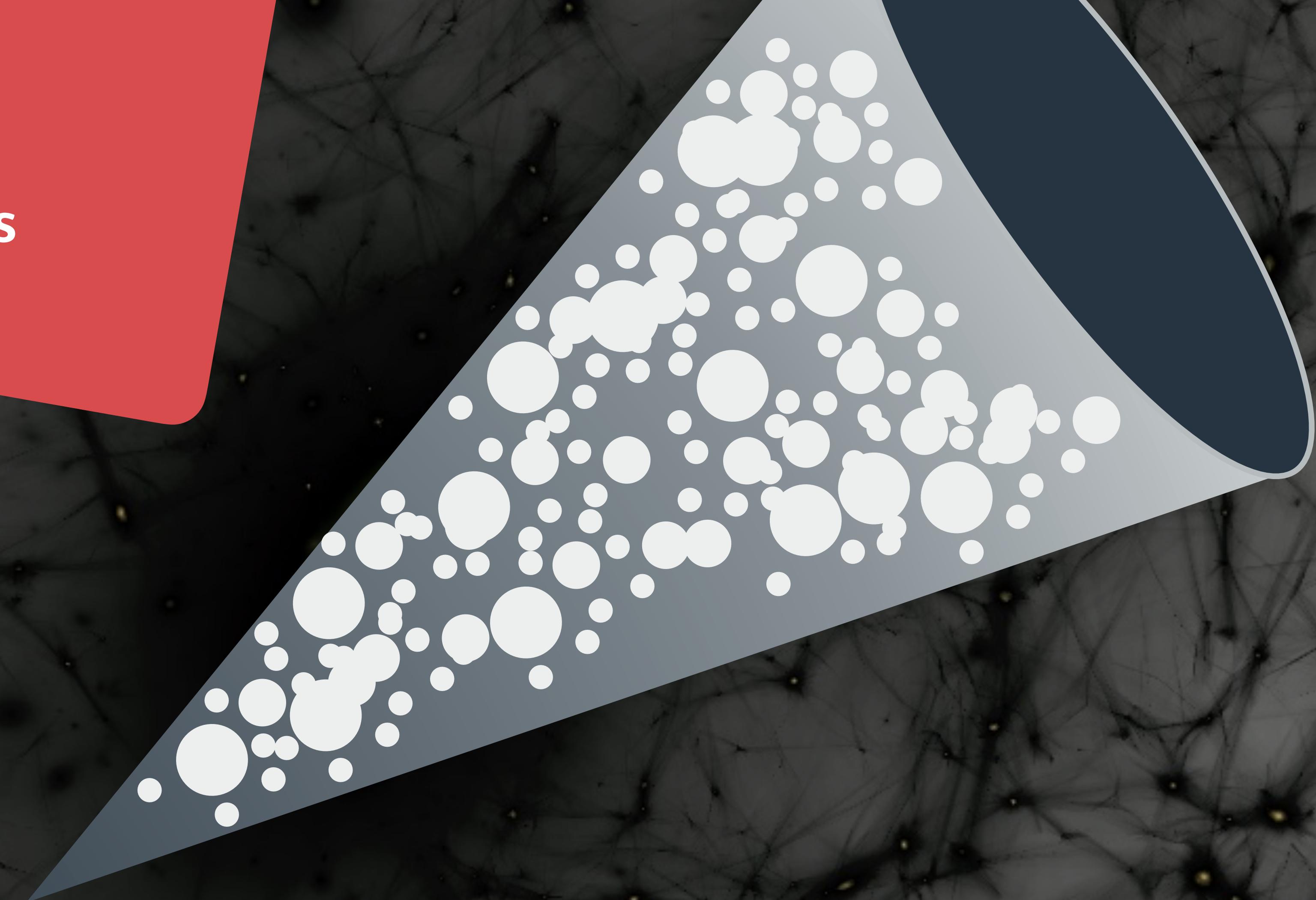


Collapse and virialisation



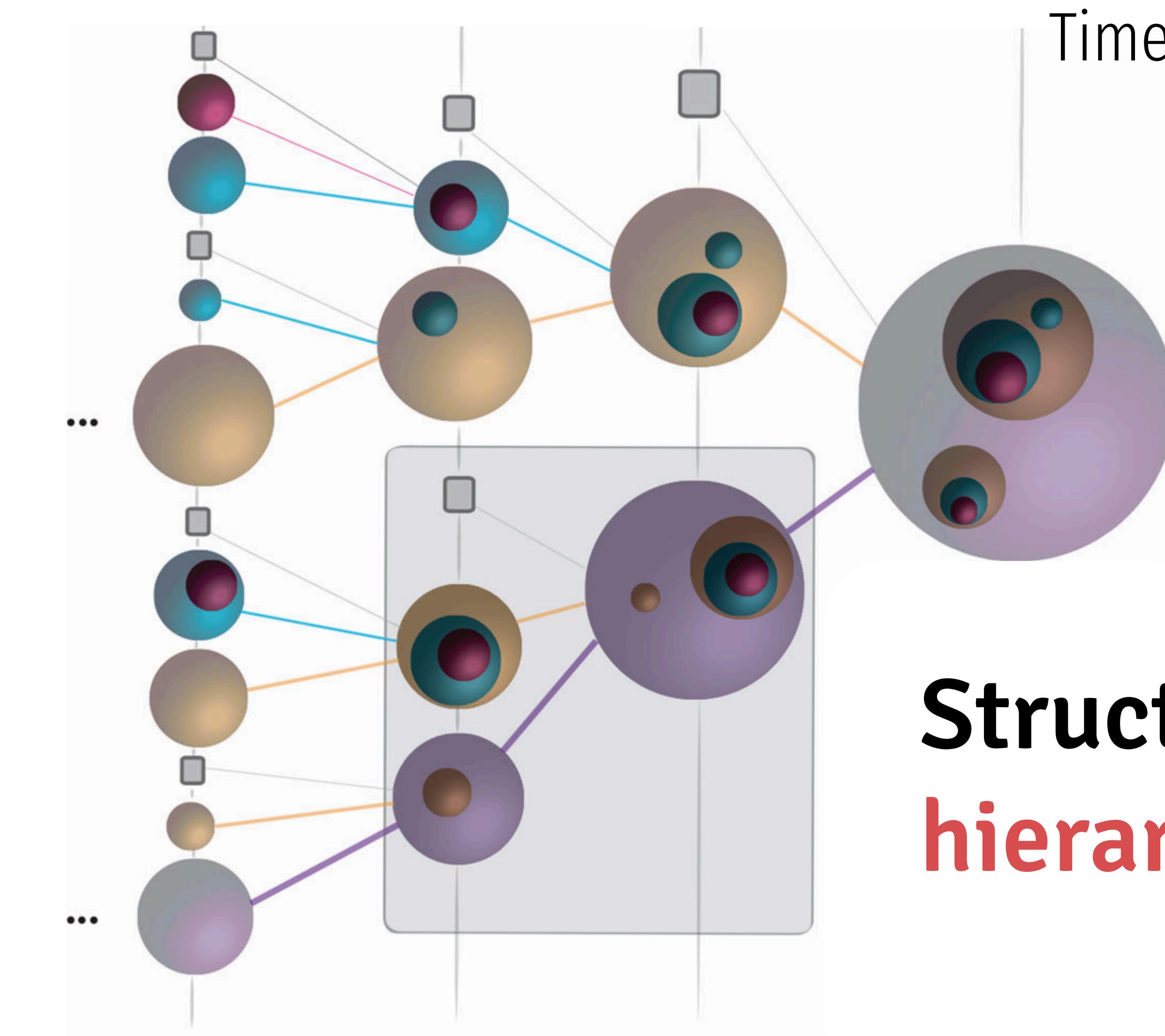
DM Halo

In the CDM paradigm,  
the Universe is populated  
by **aliens** dark matter halos



**What about  
DM subhalos?**

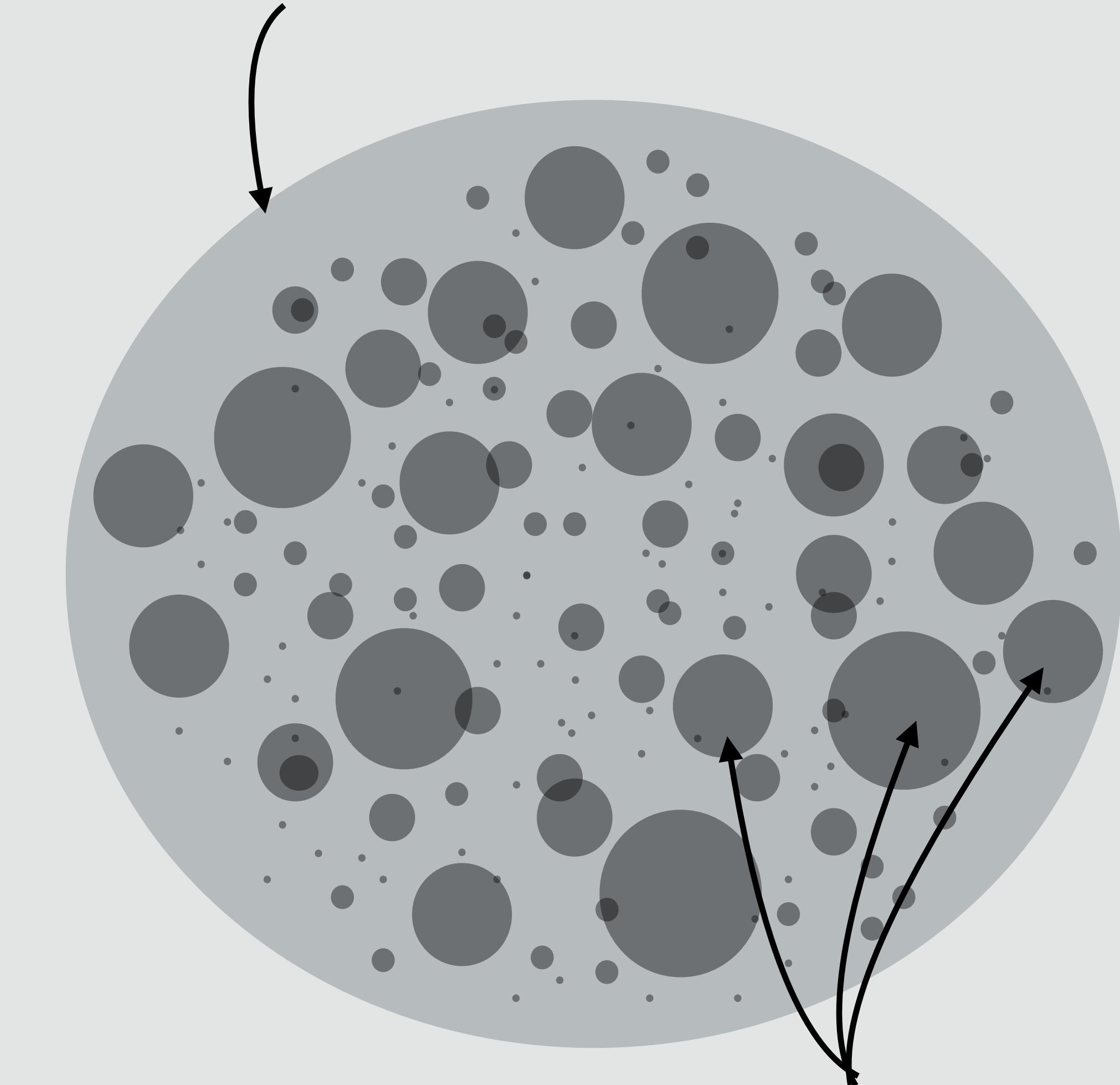
[figure from Jiang+14]



**Structures form  
hierarchically**

# Halos are clumpy

Dark matter host halo (smooth)

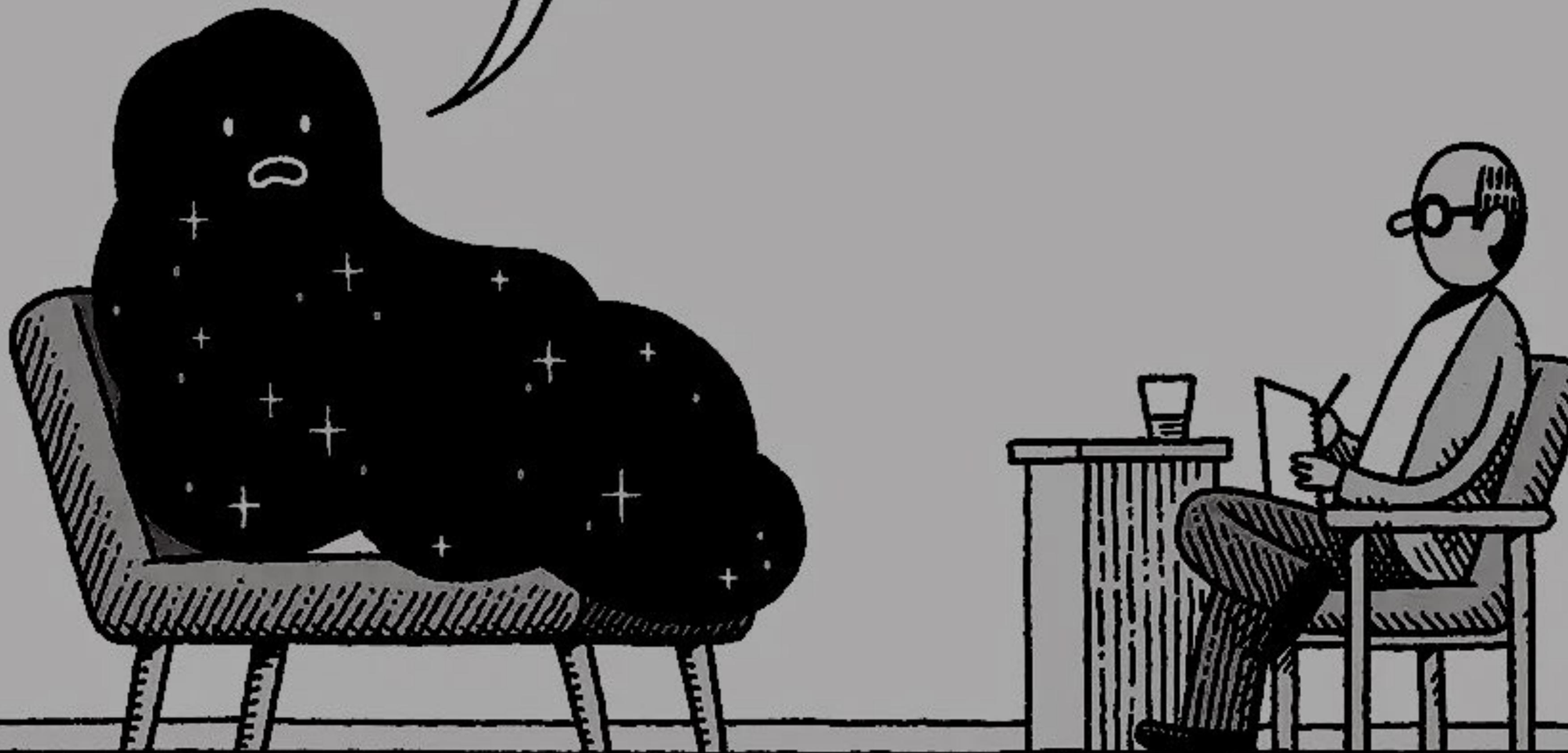


Dark matter **CLUMPS/Subhalos**  
(CDM paradigm)

**A boost?! Boost of what?**

**Let us now discuss  
DM annihilation**

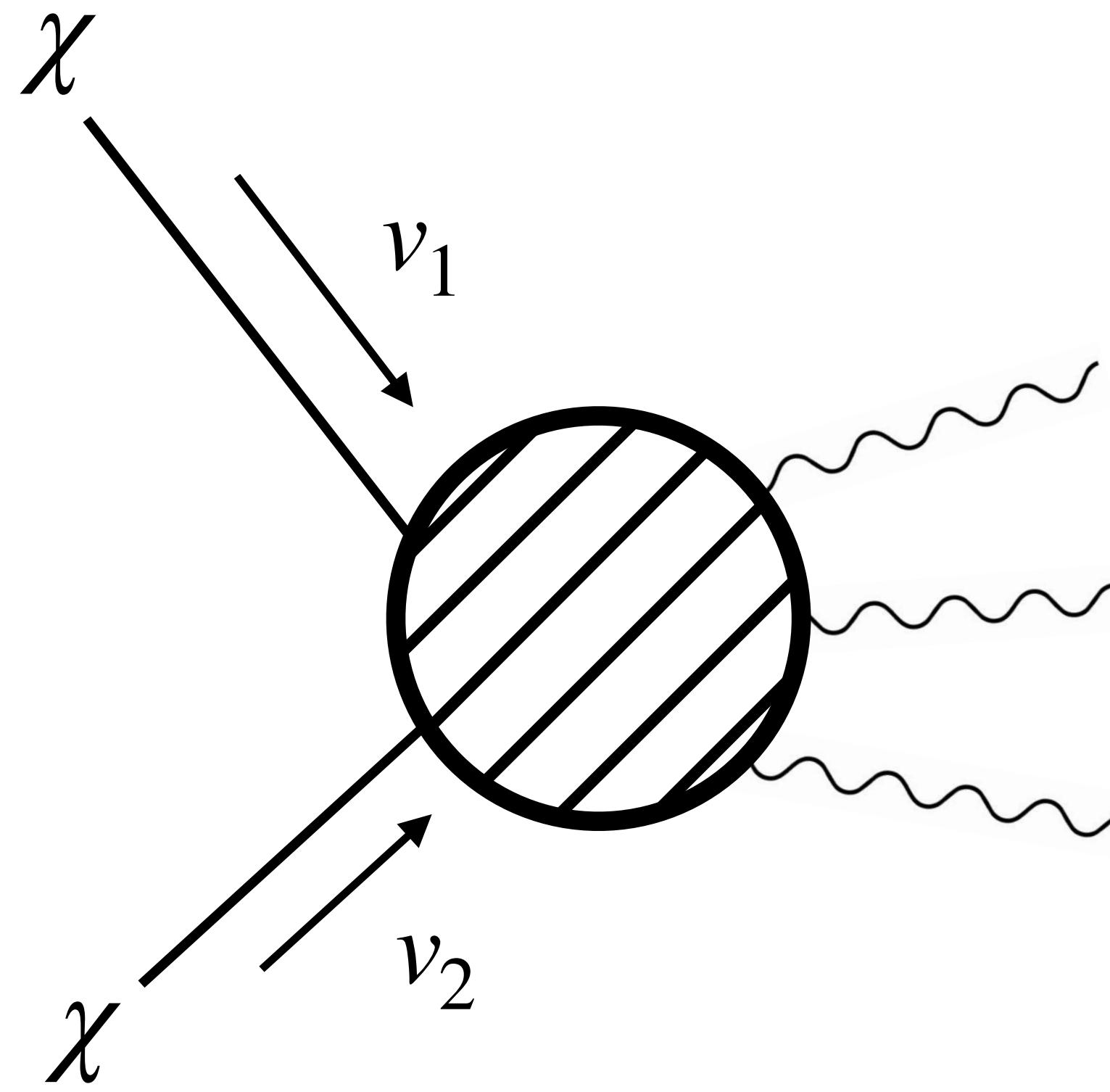
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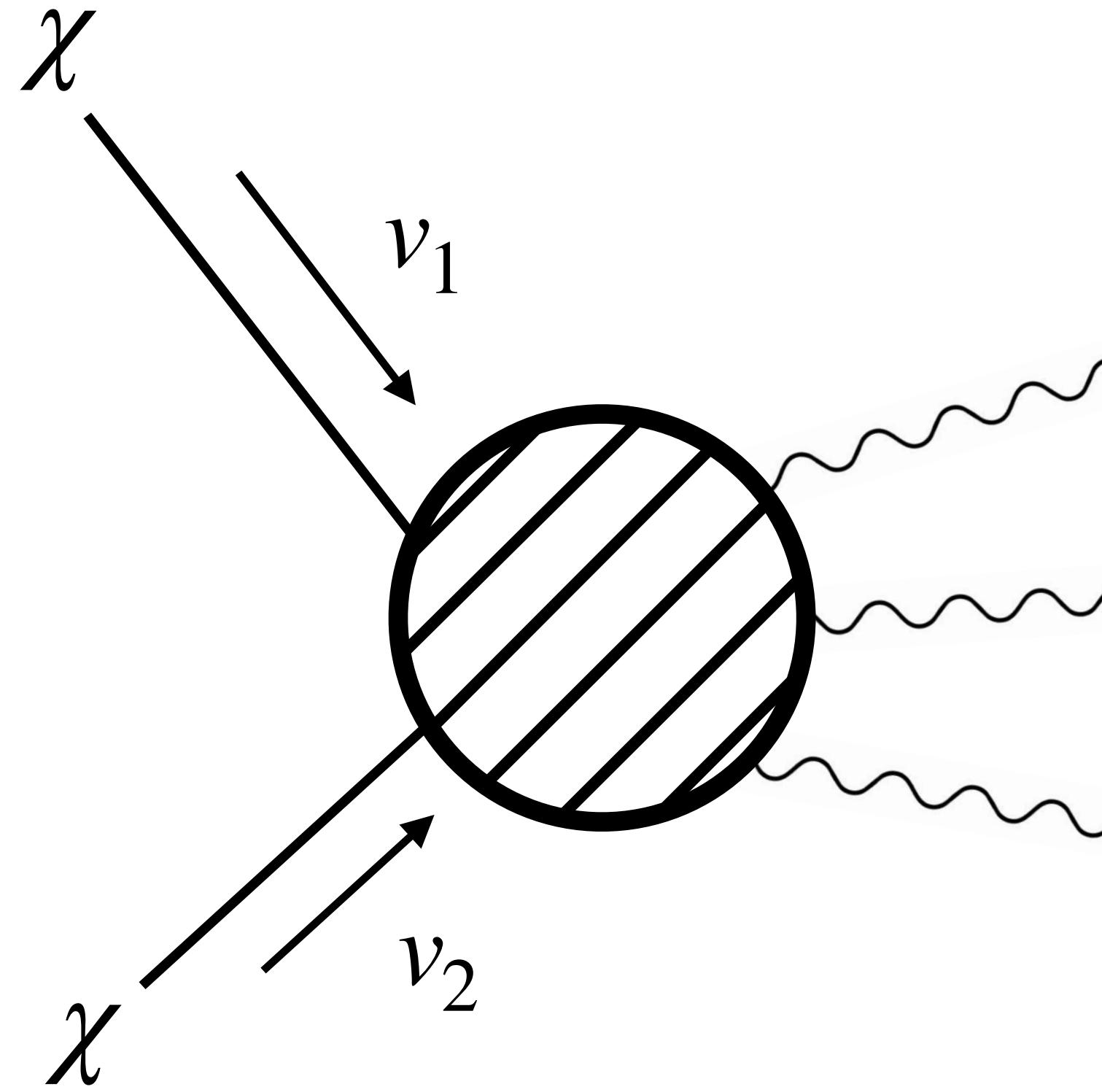
Credit: Tom Gauld  
(for NEW SCIENTIST)

**Write the annihilation  
cross-section in terms of  
the partial-wave expansion**

$$\sigma_{\text{ann}} v_{\text{rel}} = \sigma_0 + \sigma_1 v_{\text{rel}}^2 + \mathcal{O}(v_{\text{rel}}^4)$$



**Write the annihilation  
cross-section in terms of  
the partial-wave expansion**

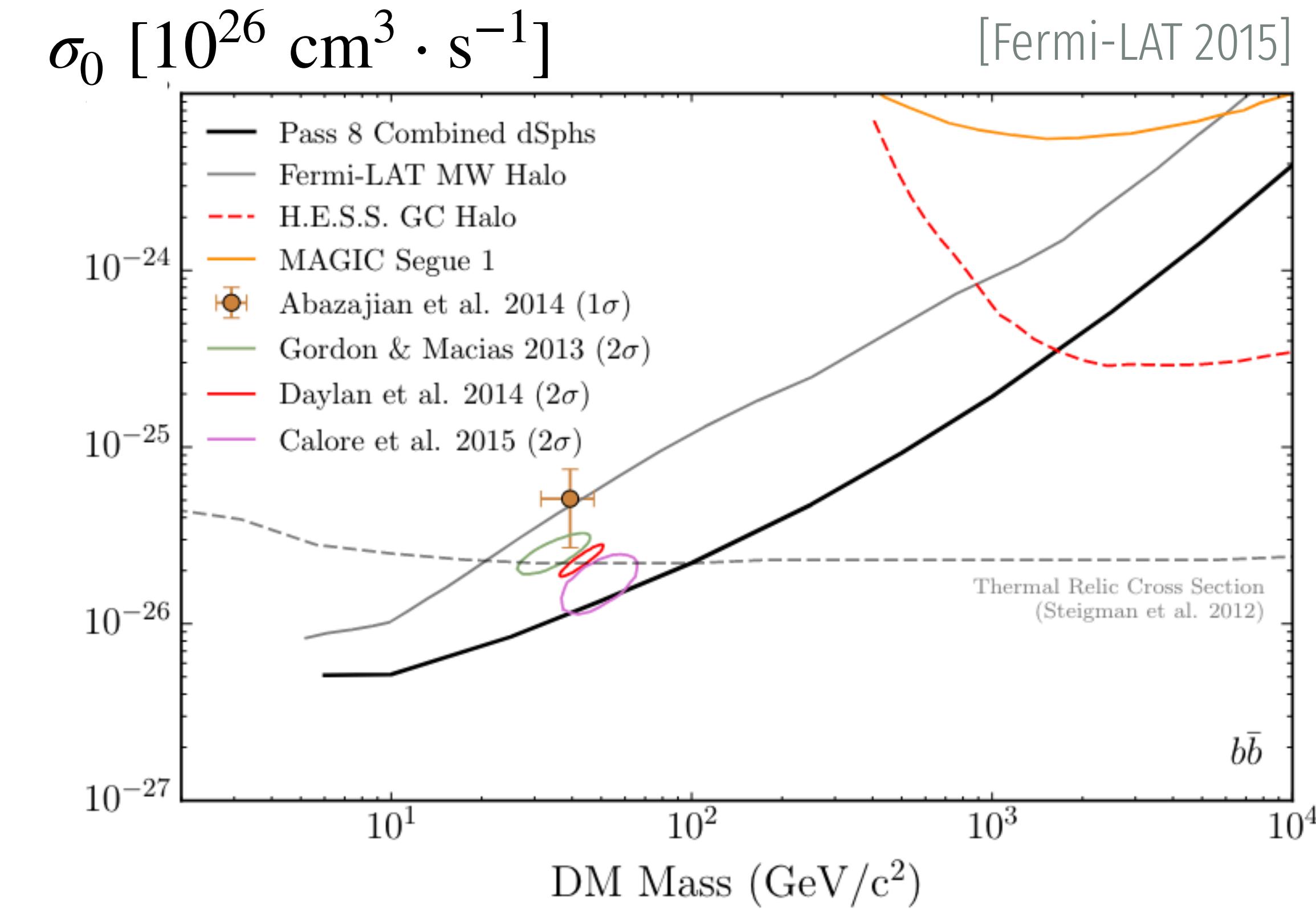
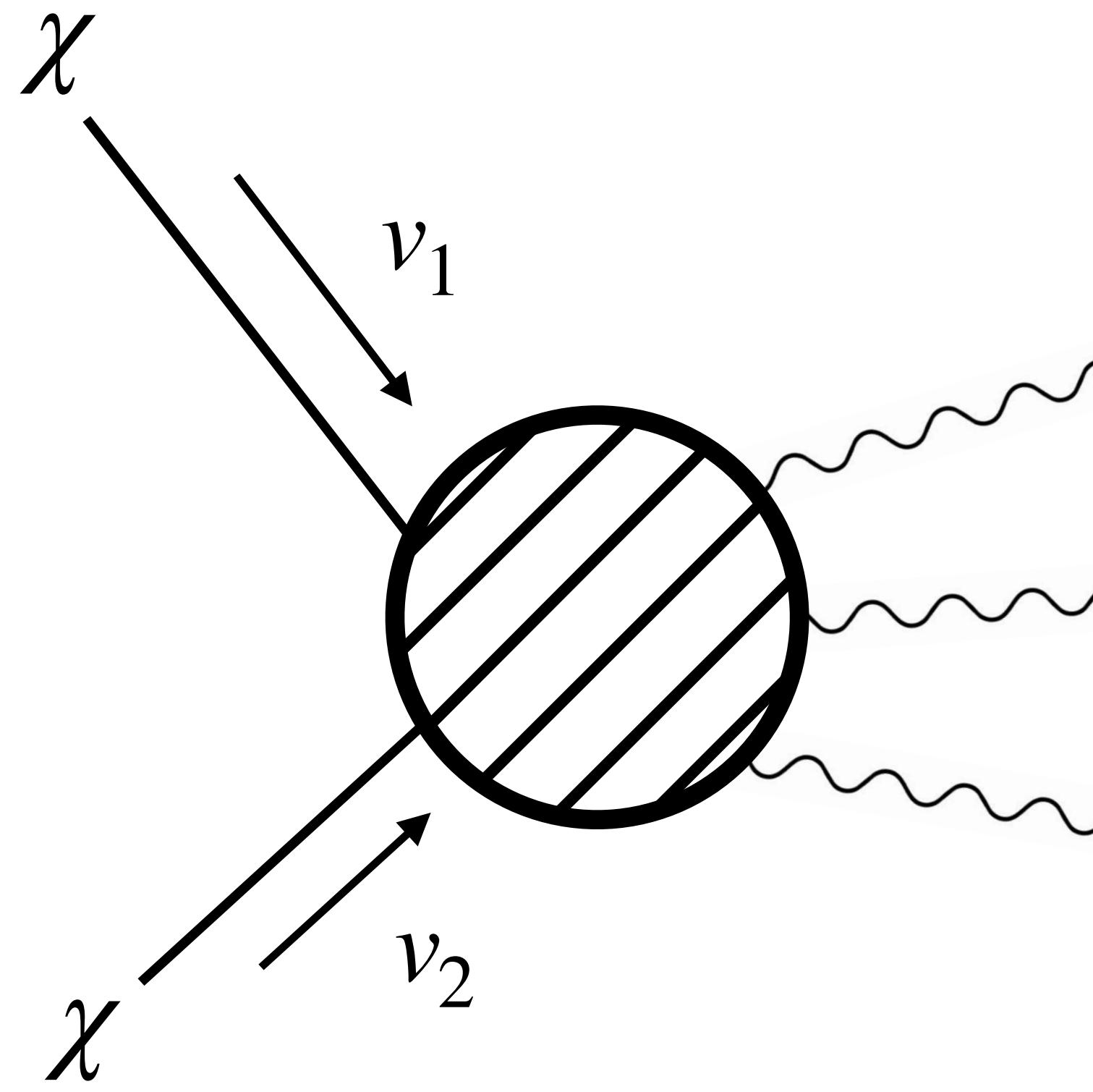


$$\sigma_{\text{ann}} v_{\text{rel}} = \sigma_0 + \sigma_1 v_{\text{rel}}^2 + \mathcal{O}(v_{\text{rel}}^4)$$

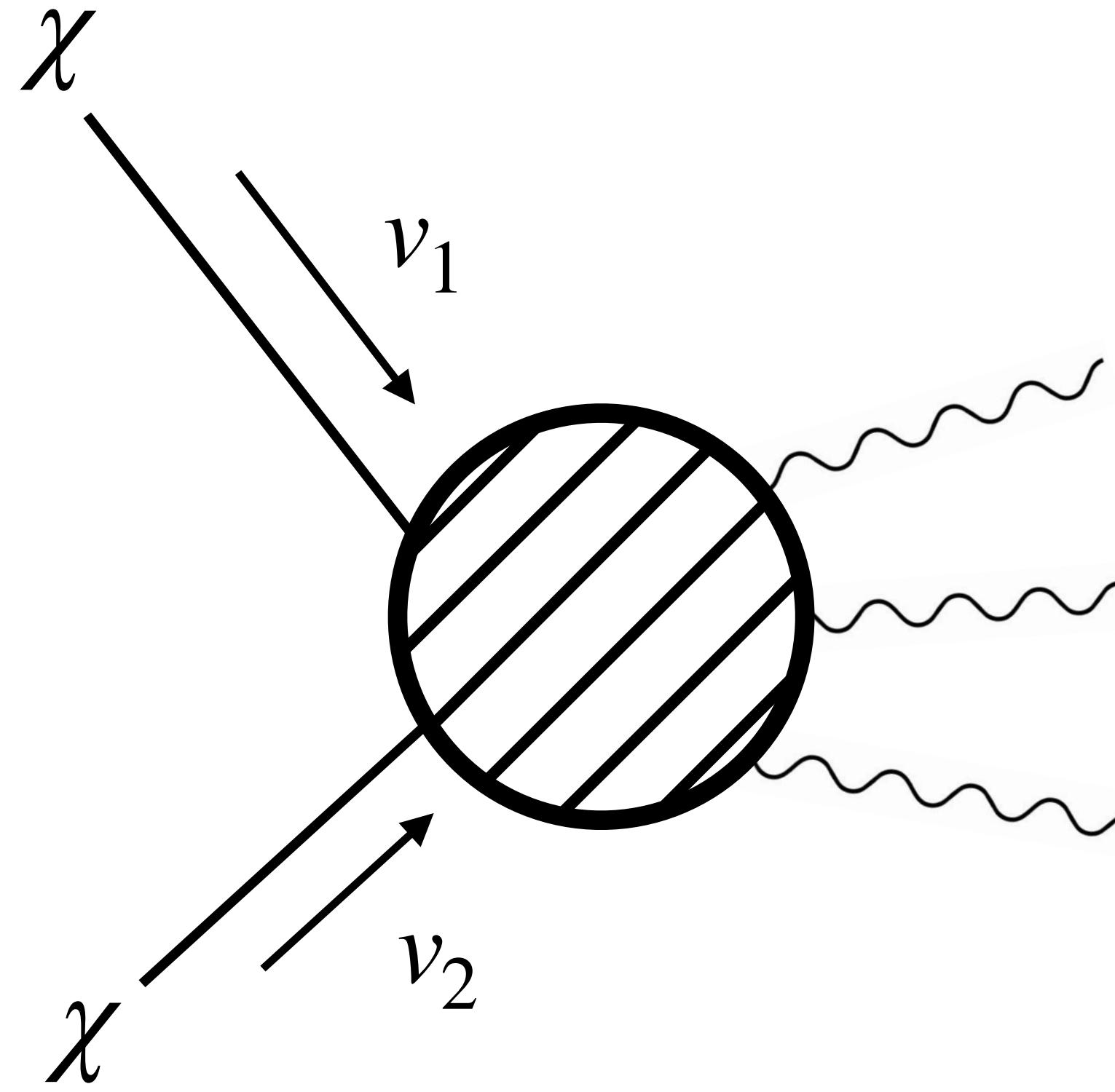
$\sigma_0$ : s-wave

Most studied case  
Strong constraints from indirect searches  
Thermal cross-section excluded  $m_\chi < 10\text{-}100 \text{ GeV}$

Write the annihilation  
cross-section in terms of  
the partial-wave expansion



Write the annihilation  
cross-section in terms of  
the partial-wave expansion

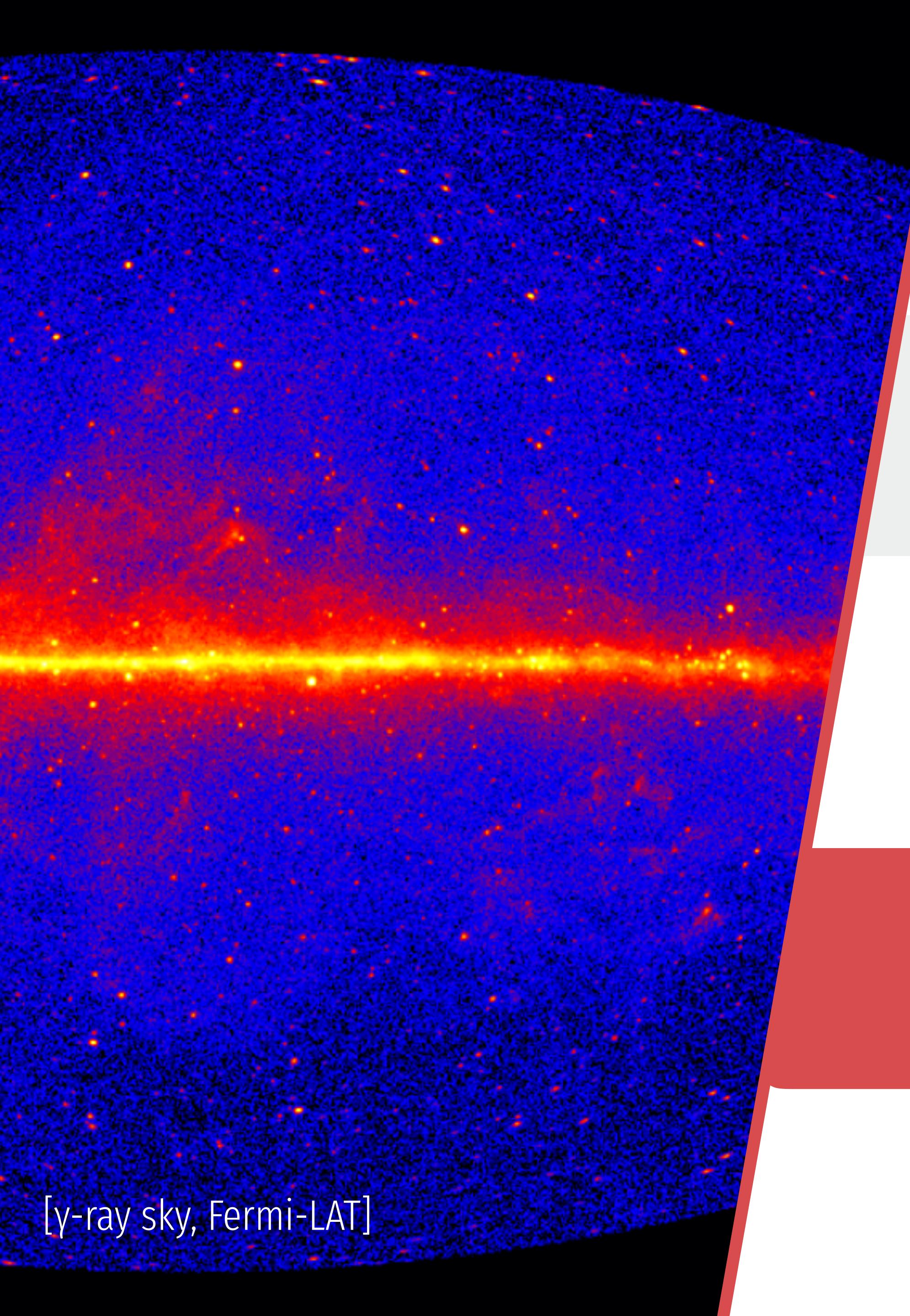


$$\sigma_{\text{ann}} v_{\text{rel}} = \sigma_0 + \sigma_1 v_{\text{rel}}^2 + \mathcal{O}(v_{\text{rel}}^4)$$

■  $\sigma_1 v^2$ : p-wave

Suppressed indirect signal  
Thermal cross-section unconstrained  
Not exotic (fermions with scalar couplings)

and in particular  
(for the example)  
**gamma rays**



[ $\gamma$ -ray sky, Fermi-LAT]

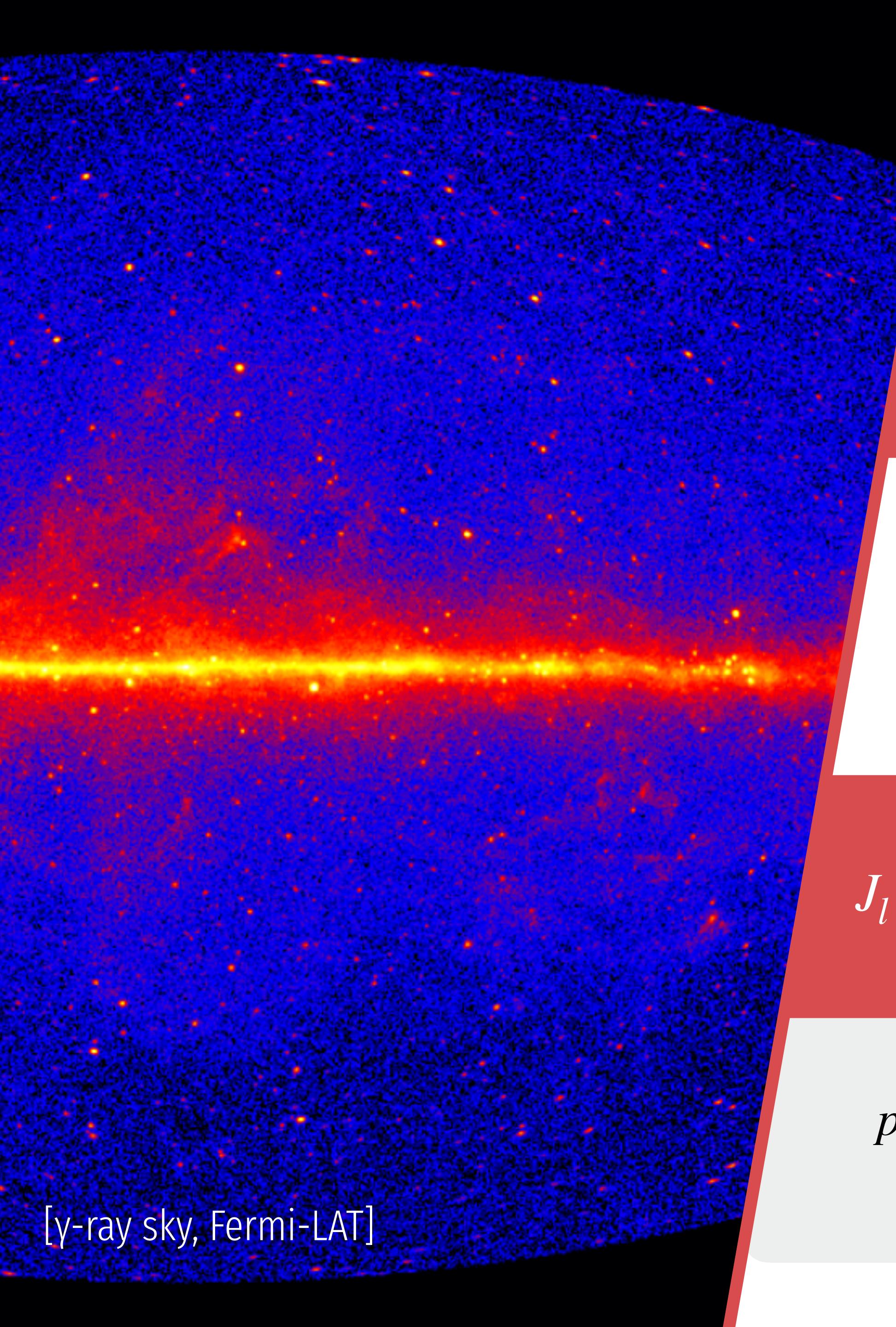
## Let us generalize

$$\sigma_{\text{ann}} v_{\text{rel}} = \sum_l \sigma_l \mathcal{S}_l(v_{\text{rel}}) \quad \mathcal{S}_l(v) = v^{2l}$$

(for partial waves)

## the associated $\gamma$ -ray flux

$$\frac{d\phi}{dE} = \frac{1}{8\pi m_\chi^2} \frac{dN_\gamma}{dE} \sum_l \sigma_l J_l$$



## Here we have the J-factor:

$$\frac{d\phi}{dE} = \frac{1}{8\pi m_\chi^2} \frac{dN_\gamma}{dE} \sum_l \sigma \mathcal{J}_l$$

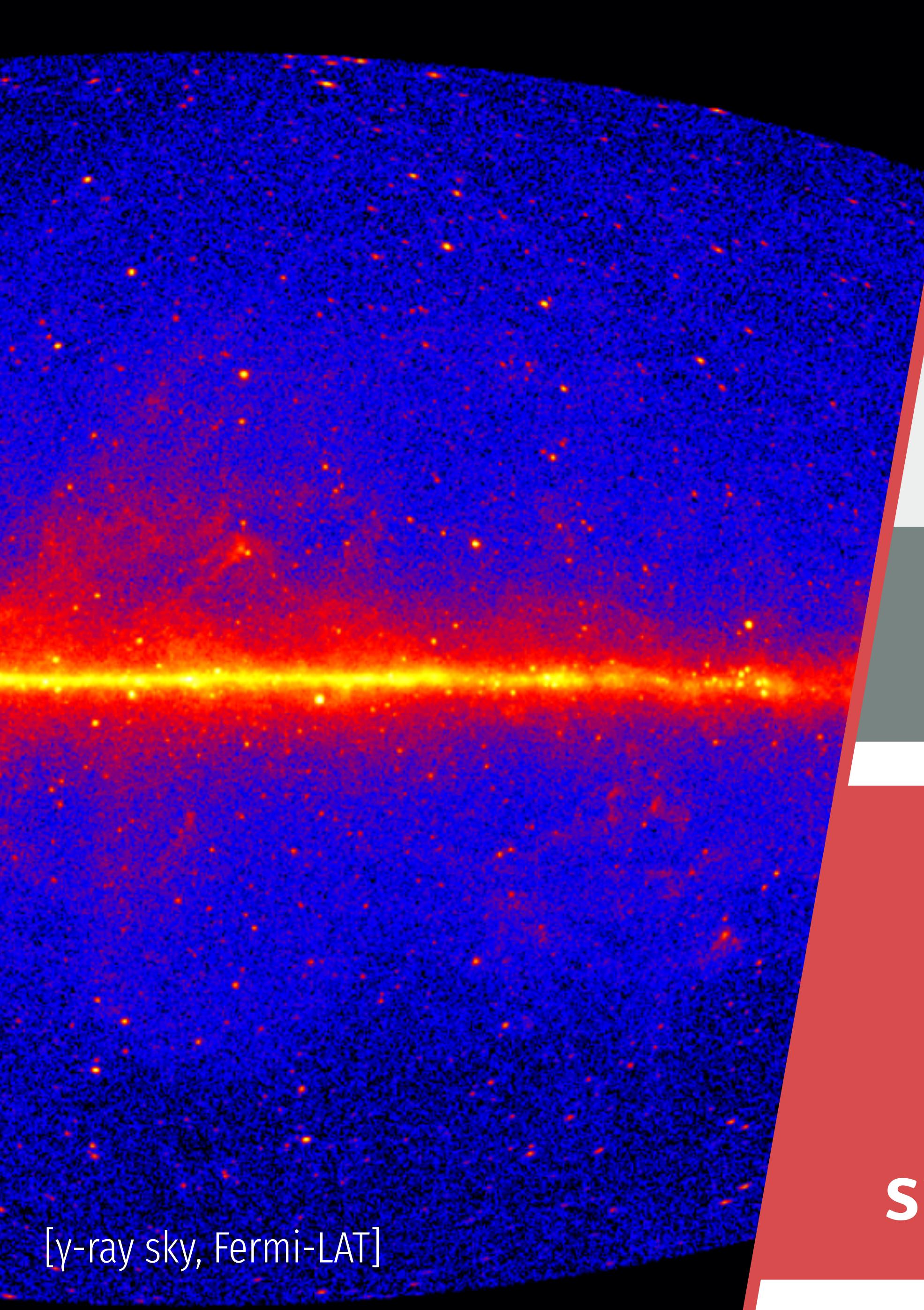
as a measure of the « luminosity »  
*(along a line of sight in a given solid angle,  
with weights from velocity factors)*

$$J_l = \int ds d\Omega \rho_\chi^2(\vec{r}(s, \vec{\theta})) \left[ \int_{\mathbb{R}^3} p_{\vec{v}_{\text{rel}}}(\vec{r}(s, \vec{\theta}), \vec{v}_{\text{rel}}) \mathcal{S}_l(v_{\text{rel}}) d^3 \vec{v}_{\text{rel}} \right]$$

$$p_{\vec{v}_{\text{rel}}}(\vec{r}, \vec{v}_{\text{rel}}) \equiv \frac{1}{\rho_\chi^2(r)} \int f\left(\vec{r}, \vec{v}_c - \frac{\vec{v}_{\text{rel}}}{2}\right) f\left(\vec{r}, \vec{v}_c + \frac{\vec{v}_{\text{rel}}}{2}\right) d^3 \vec{v}_c$$

(distribution of relative velocity)

[ $\gamma$ -ray sky, Fermi-LAT]



## (example) s-wave case:

$$\frac{d\phi}{dE} = \frac{\sigma_0}{8\pi m_\chi^2} \frac{dN_\gamma}{dE} J_0$$

$$J_0 = \int d\mathbf{s} d\Omega \rho_\chi^2(\vec{r}(s, \vec{\theta})) \left[ \int_{\mathbb{R}^3} p_{\vec{v}_{\text{rel}}}(\vec{r}(s, \vec{\theta}), \vec{v}_{\text{rel}}) d^3 \vec{v}_{\text{rel}} \right]$$

$$J_0 = \int d\mathbf{s} d\Omega \rho_\chi^2(\vec{r}(s, \vec{\theta})) \propto \frac{m}{4\pi D^2}$$

**Squared density: the presence of subhalos enhances/modifies the signal**

$$\text{boost factor} = \frac{\text{signal with clumpiness (true signal)}}{\text{signal if DM was smoothly distributed}}$$

The boost factor is encoded  
in the J-factors/luminosity

$$\text{boost factor} = \frac{J_l^{\text{smooth part}} + J_l^{\text{clumps}}}{J_l^{\text{all smooth}}}$$

I will talk more about this  
**astrophysical** examples ...

... but the boost may also be  
important in **cosmology**  
(e.g. for the 21cm)

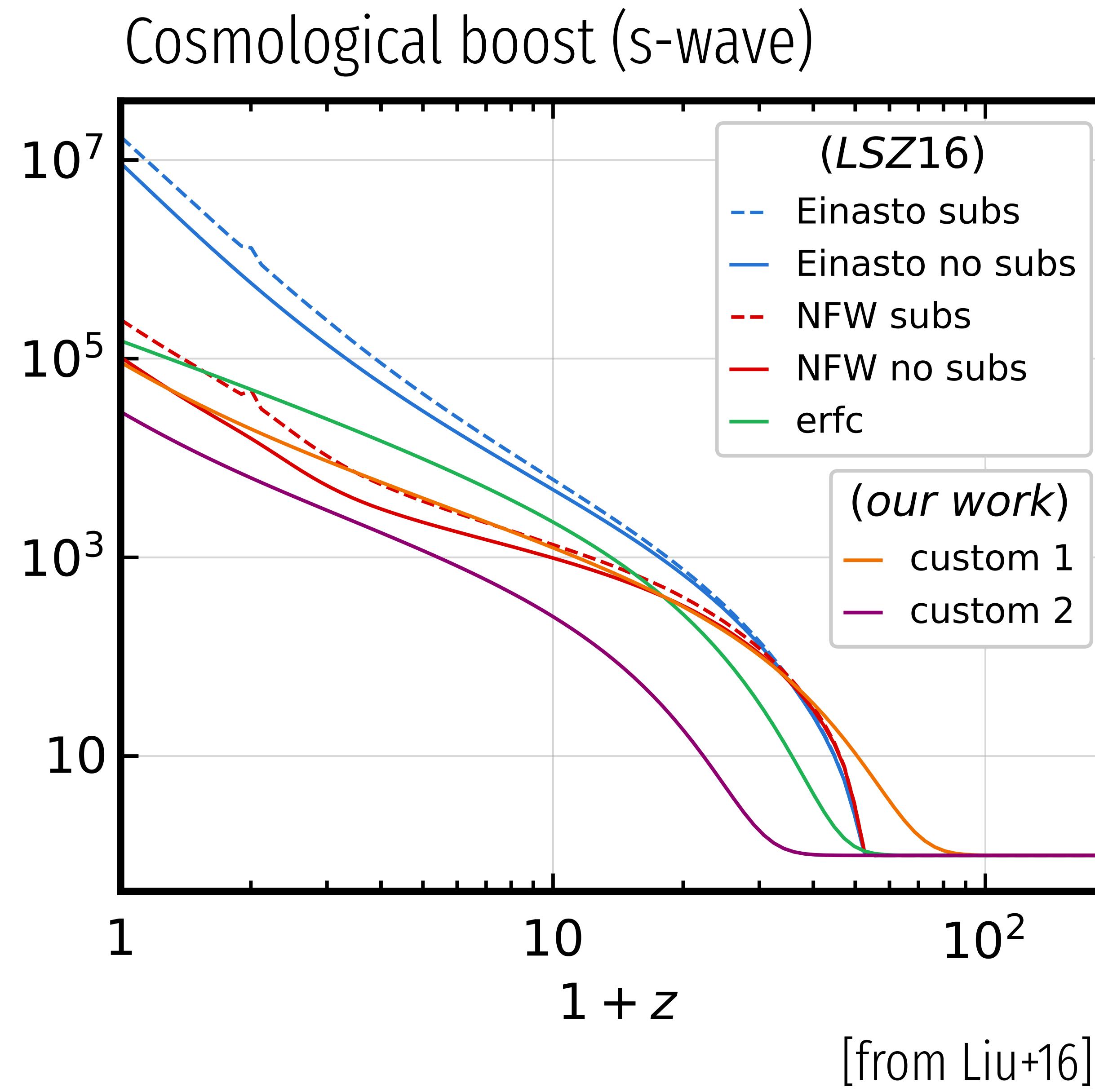
**Enhancement due to  
subhalos**



Boost from halos  
but  
without subhalos



Boost from halos  
and  
without subhalos



II.

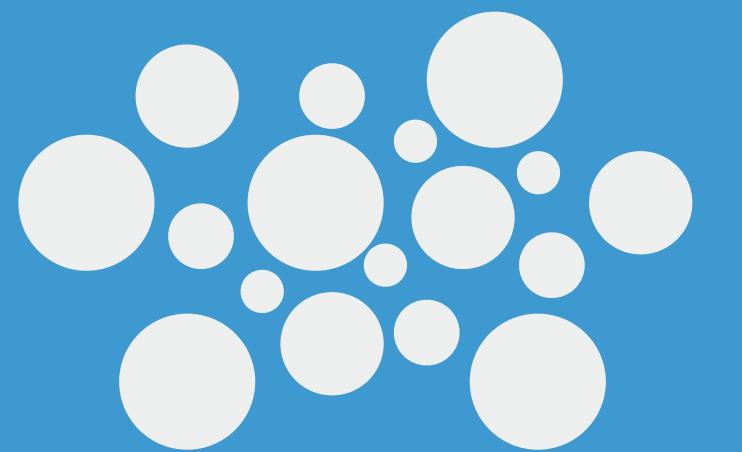
Sommerfeld enhanced  
and subhalo boosted  
J-factors

I.

A generic analytical  
subhalo model

III.

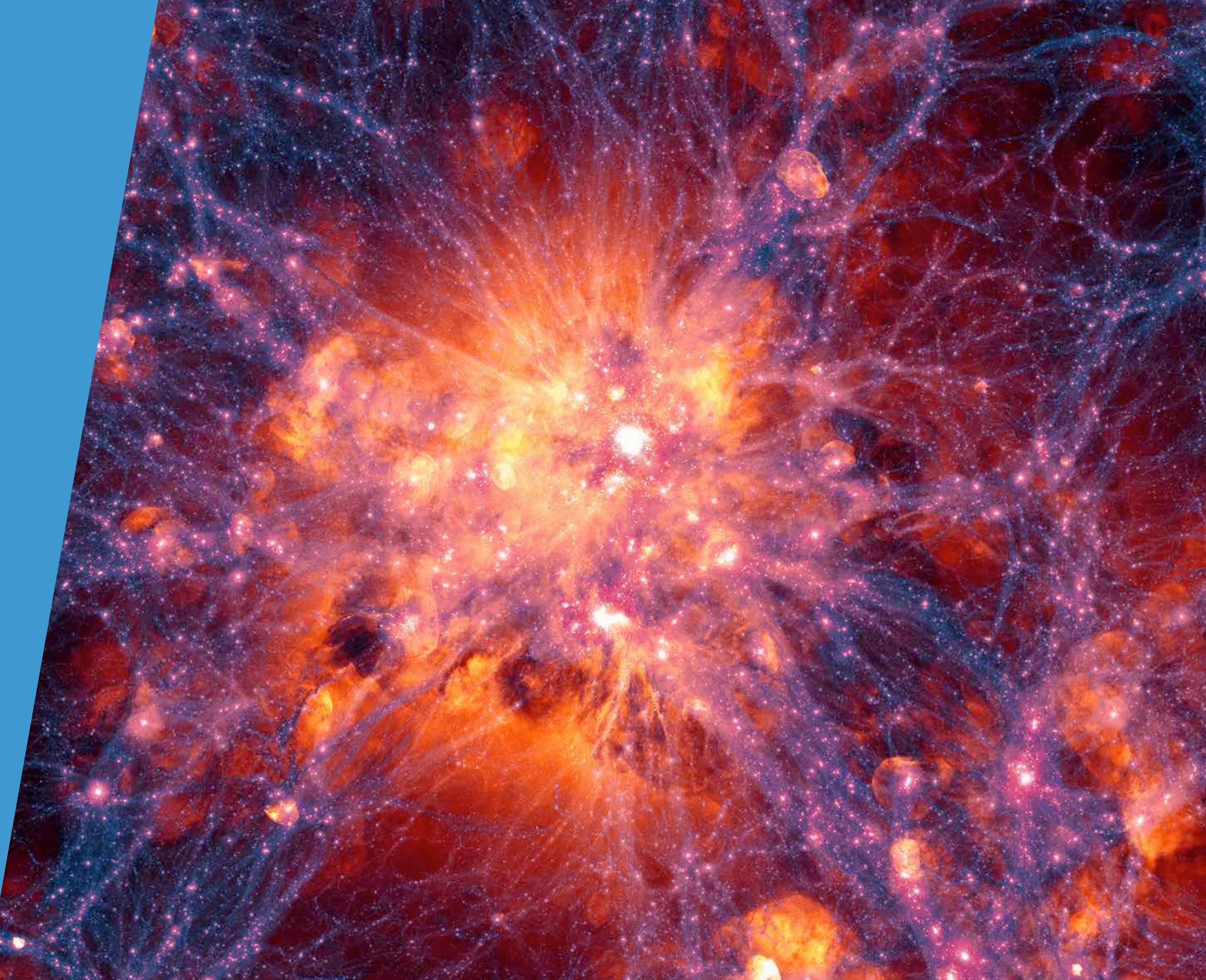
Subhalos as point sources  
in the Milky Way



## Part 1:

### A generic analytical subhalo model

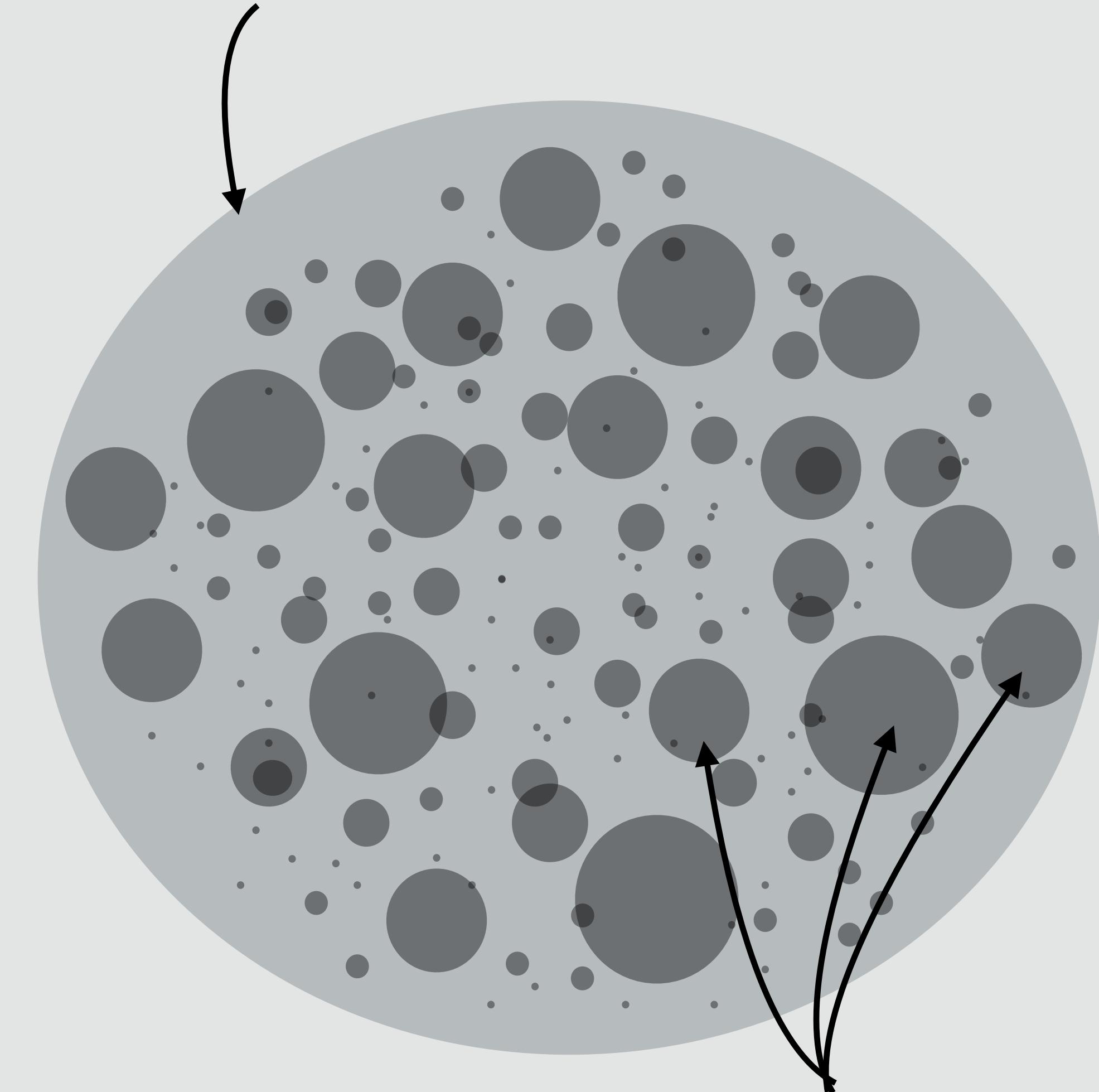
[1610.02233, 2201.09788]



# **Consider a host halo**

$$\rho_\chi = \rho_{\text{smooth}} + \sum_{i=1}^{N_{\text{sub}}} \rho_i$$

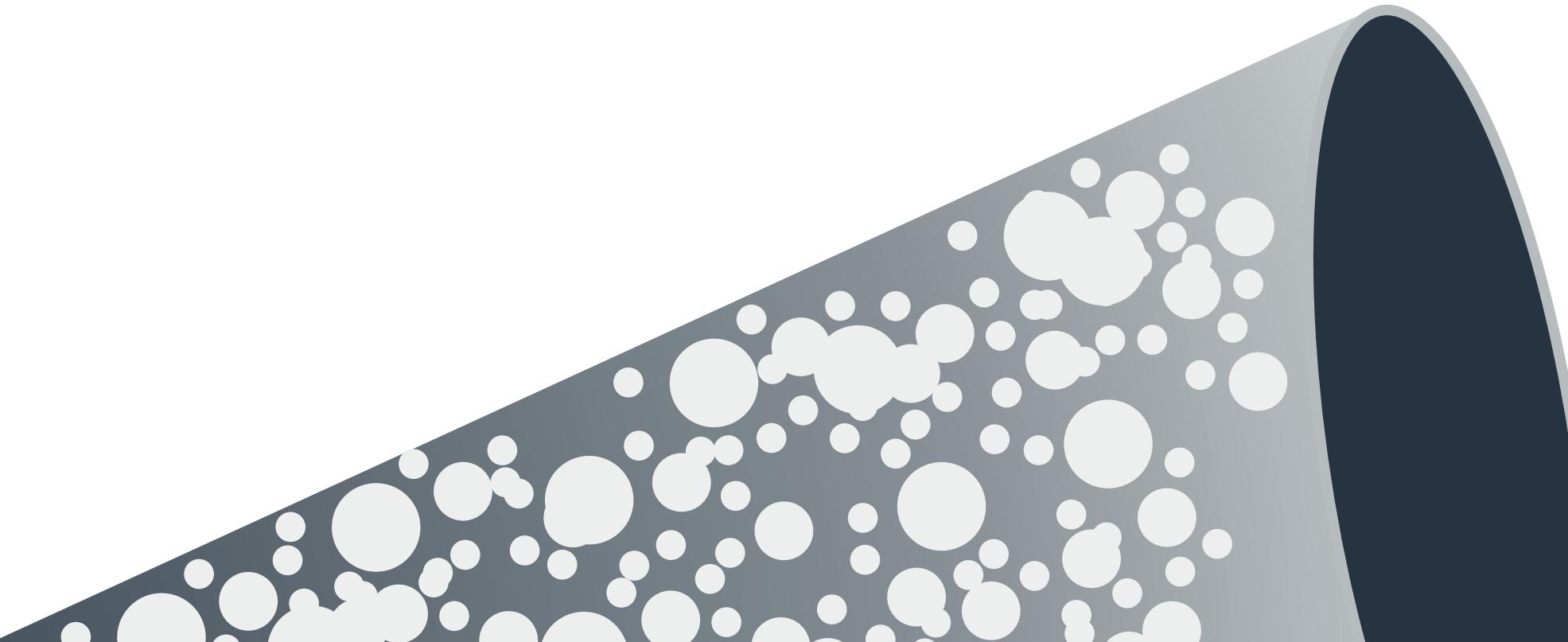
Dark matter host halo (smooth)



Dark matter **CLUMPS/Subhalos**  
(CDM paradigm)

# Subhalo contribution to the J-factors (in a given volume)

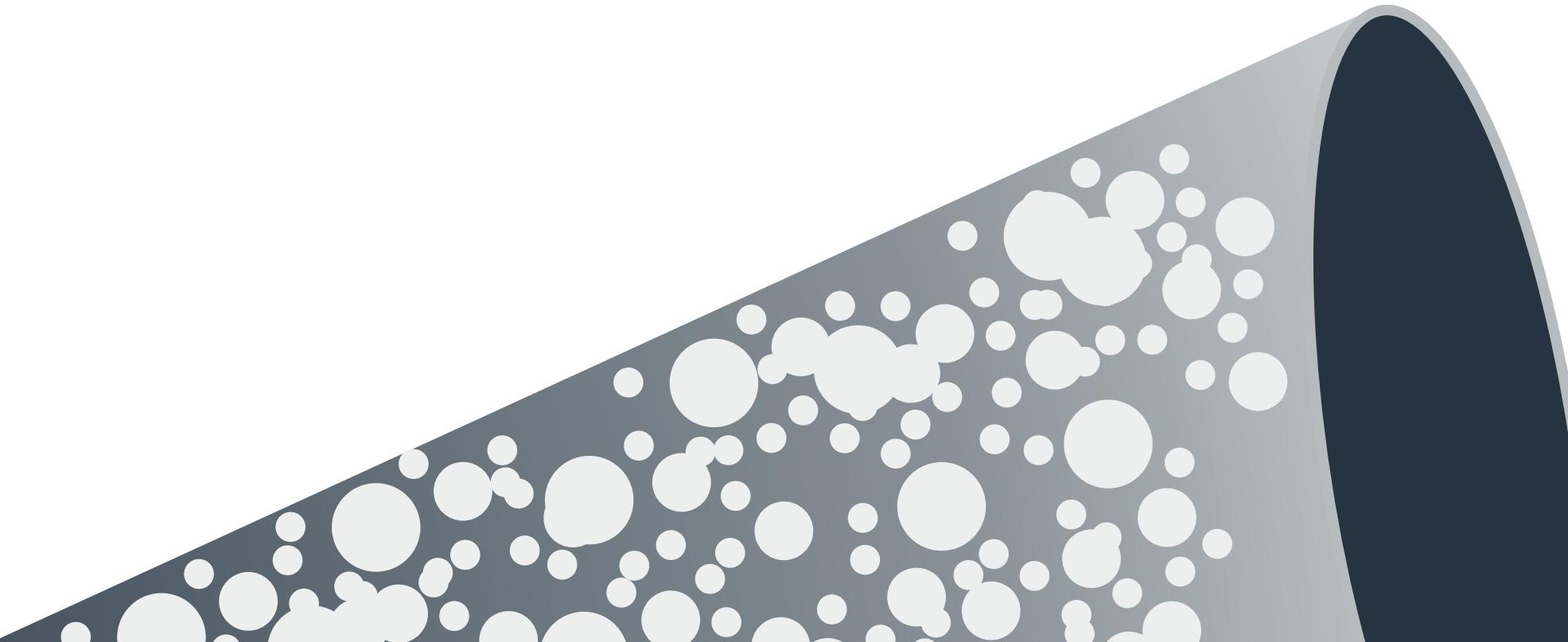
$$J_l^{\text{sub}} \equiv \frac{1}{D^2} \int_{\mathcal{V}} \frac{\partial^2 n}{\partial m_t \partial c} \Big|_{\text{sub}} (m_t, c, \vec{R}) \mathcal{L}_l(m_t, c) dm_t dc d^3 \vec{R}$$



# Subhalo contribution to the J-factors (in a given volume)

$$J_l^{\text{sub}} \equiv \frac{1}{D^2} \int_{\mathcal{V}} \frac{\partial^2 n}{\partial m_t \partial c} \Big|_{\text{sub}} (m_t, c, \vec{R}) \mathcal{L}_l(m_t, c) dm_t dc d^3 \vec{R}$$

Luminosity of a  
single subhalo



# Subhalo contribution to the J-factors (in a given volume)

$$J_l^{\text{sub}} \equiv \frac{1}{D^2} \int_{\mathcal{V}} \frac{\partial^2 n}{\partial m_t \partial c} \Big|_{\text{sub}} (m_t, c, \vec{R}) \mathcal{L}_l(m_t, c) dm_t dc d^3 \vec{R}$$

Luminosity of a  
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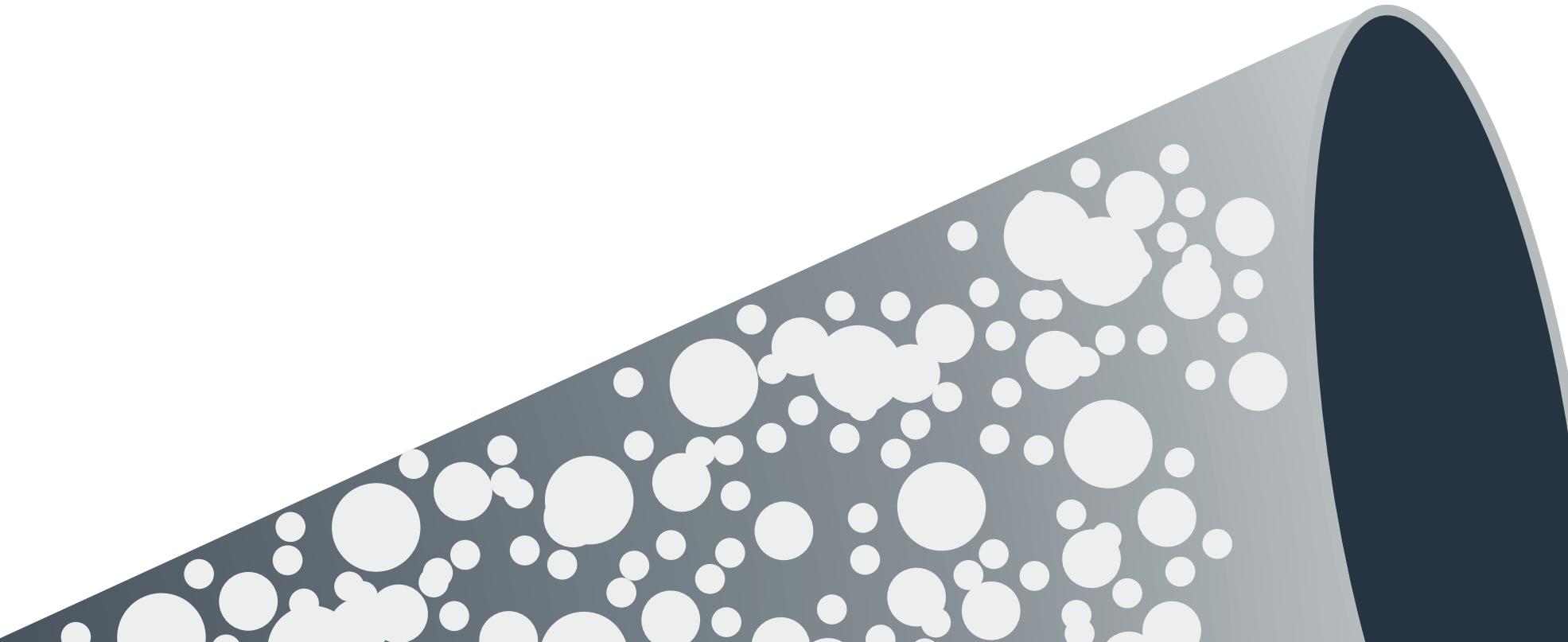
« Luminosity » of a subhalo (of mass  $m_t$  and concentration  $c$ )

$$\begin{aligned} \mathcal{L}_l(m_t, c) &= \int \rho_{\text{sub}}^2(\vec{r}) \int f_{\vec{v}_{\text{rel}}}^{\text{sub}}(\vec{r}, \vec{v}_{\text{rel}}) \mathcal{S}_l(v_{\text{rel}}) d^3 \vec{v}_{\text{rel}} d^3 \vec{r} \\ &\simeq \int \rho_{\text{sub}}^2(\vec{r}) \mathcal{S}_l[\langle v_{\text{rel}} \rangle_n(\vec{r})] d^3 \vec{r} \quad \text{with} \\ \langle v_{\text{rel}} \rangle_n(\vec{r}) &= \left[ \int p_{\vec{v}_{\text{rel}}}^{\text{sub}}(\vec{r}, \vec{v}_{\text{rel}}) v_{\text{rel}}^n \right]^{1/n} \end{aligned}$$

# Subhalo contribution to the J-factors (in a given volume)

$$J_l^{\text{sub}} \equiv \frac{1}{D^2} \int_{\mathcal{V}} \left. \frac{\partial^2 n}{\partial m_t \partial c} \right|_{\text{sub}} (m_t, c, \vec{R}) \mathcal{L}_l(m_t, c) dm_t dc d^3 \vec{R}$$

subhalo distribution



# How to describe the subhalo population?

[GF, Stref and Lavalle 2022, Stref+17,  
Benson+12, Bartels+15,  
Hiroshima+18, Hiroshima+22,  
Zavala+14,  
Van den Bosch+05,  
Peñarrubia+05, ...]

with analytical models

Number of CDM subhalos in the targets  $\gg 10^5$

Use a statistical description of the subhalos

A recipe from  
[Stref and Lavalle 2017]  
[GF, Stref and Lavalle 2022]



# How to describe the subhalo population?

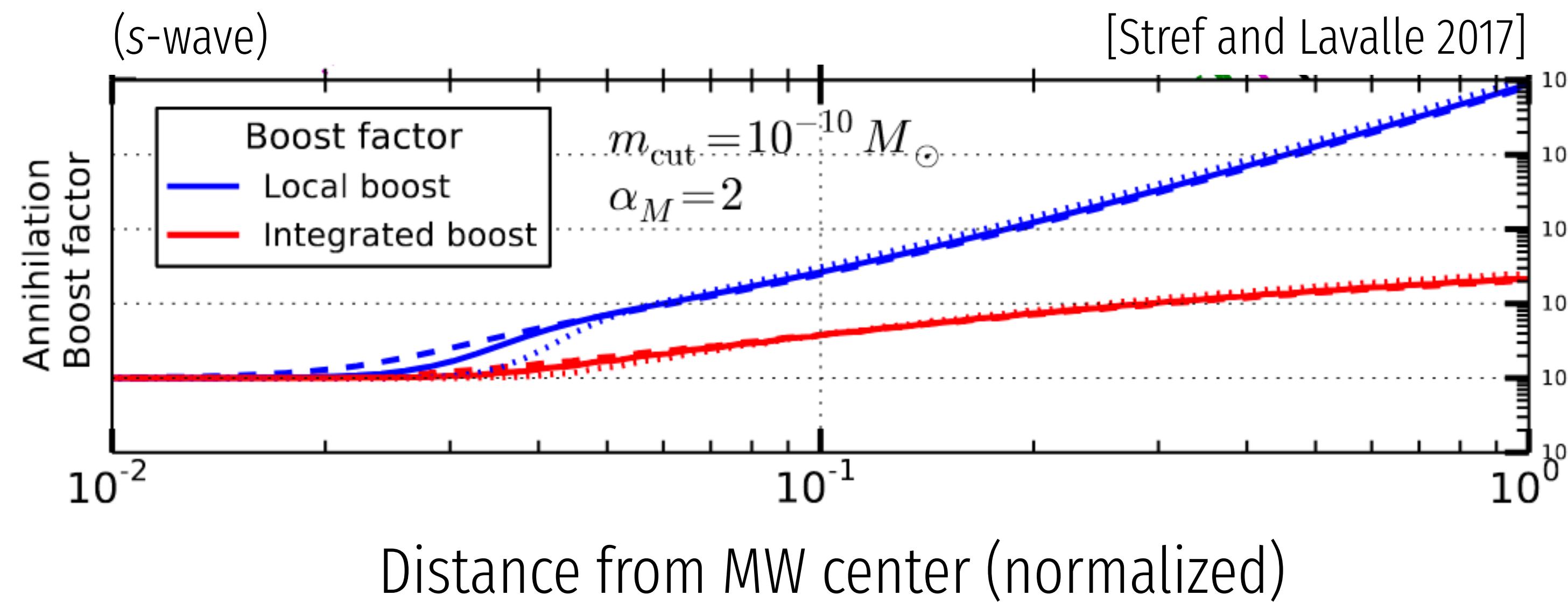
[GF, Stref and Lavalle 2022, Stref+17,  
Benson+12, Bartels+15,  
Hiroshima+18, Hiroshima+22,  
Zavala+14,  
Van den Bosch+05,  
Peñarrubia+05, ...]

in two lines:

- 1. Cosmological subhalo distribution
- 2. Dynamical effects in the host  
(Make subhalo shrink over time)



# First application of the model:



## Boost factor in the Milky Way

# Building an analytical model for a subhalo population: the recipe

which, in equation, gives:

$$\frac{\partial^2 n}{\partial m_t \partial c} \Big|_f = \int \frac{\partial^2 n}{\partial m \partial c} \Big|_i \Theta \left( \frac{r_t(m, c, \vec{R}, z)}{r_s(m, c, z)} - \epsilon_t \right) \delta[m_t - m_t^\star(m, c, \vec{R}, z)] dm$$

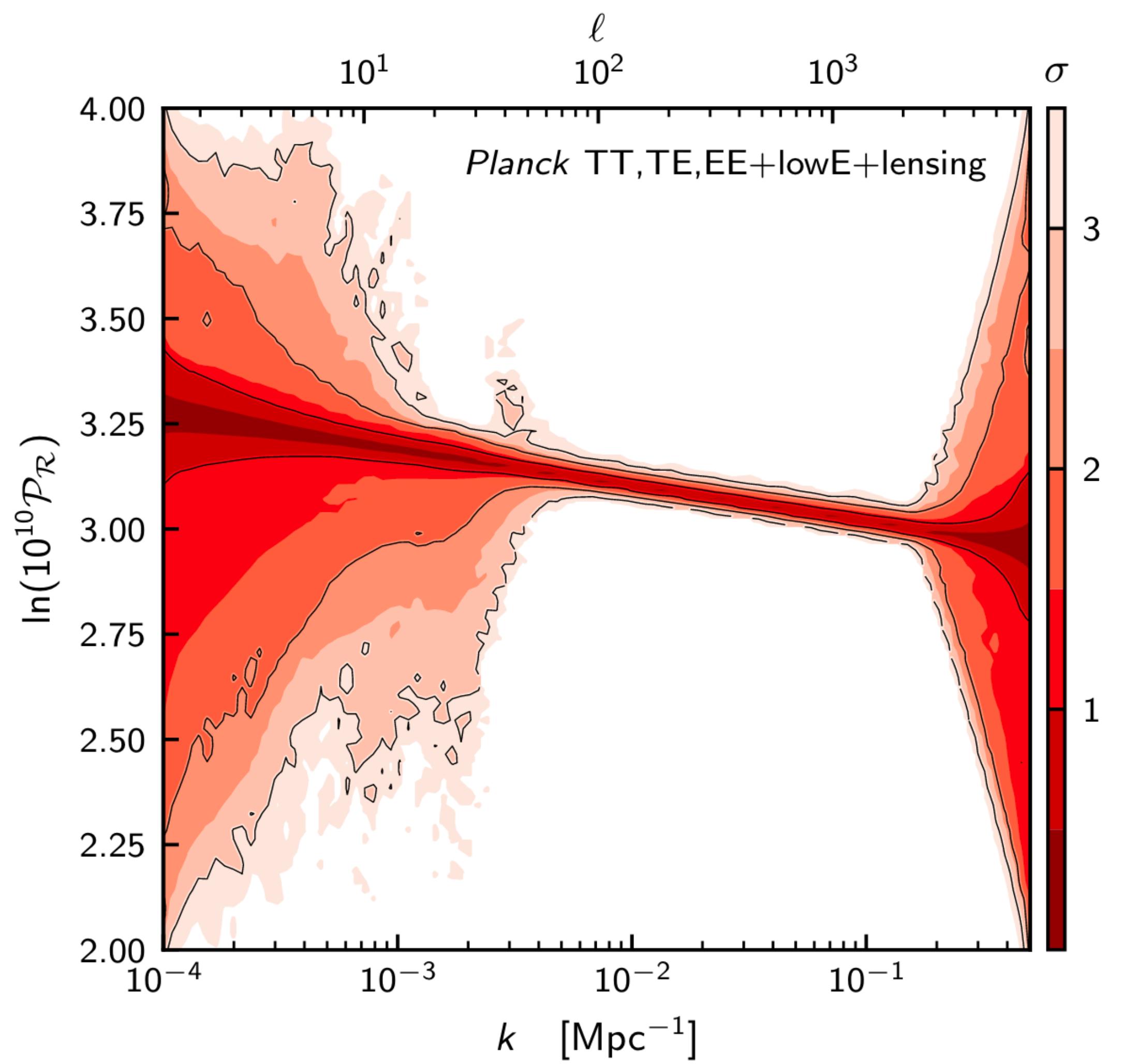
Initial (cosmological) distribution

$$\frac{\partial^2 n}{\partial m \partial c} \Big|_i = \frac{dN_{\text{sub}}}{dm}(m | M_{\text{host}}, z) p_{\vec{R}}(\vec{R}) p_c(c | m)$$

Tidal disruption/stripping terms



# The cosmological subhalo mass function



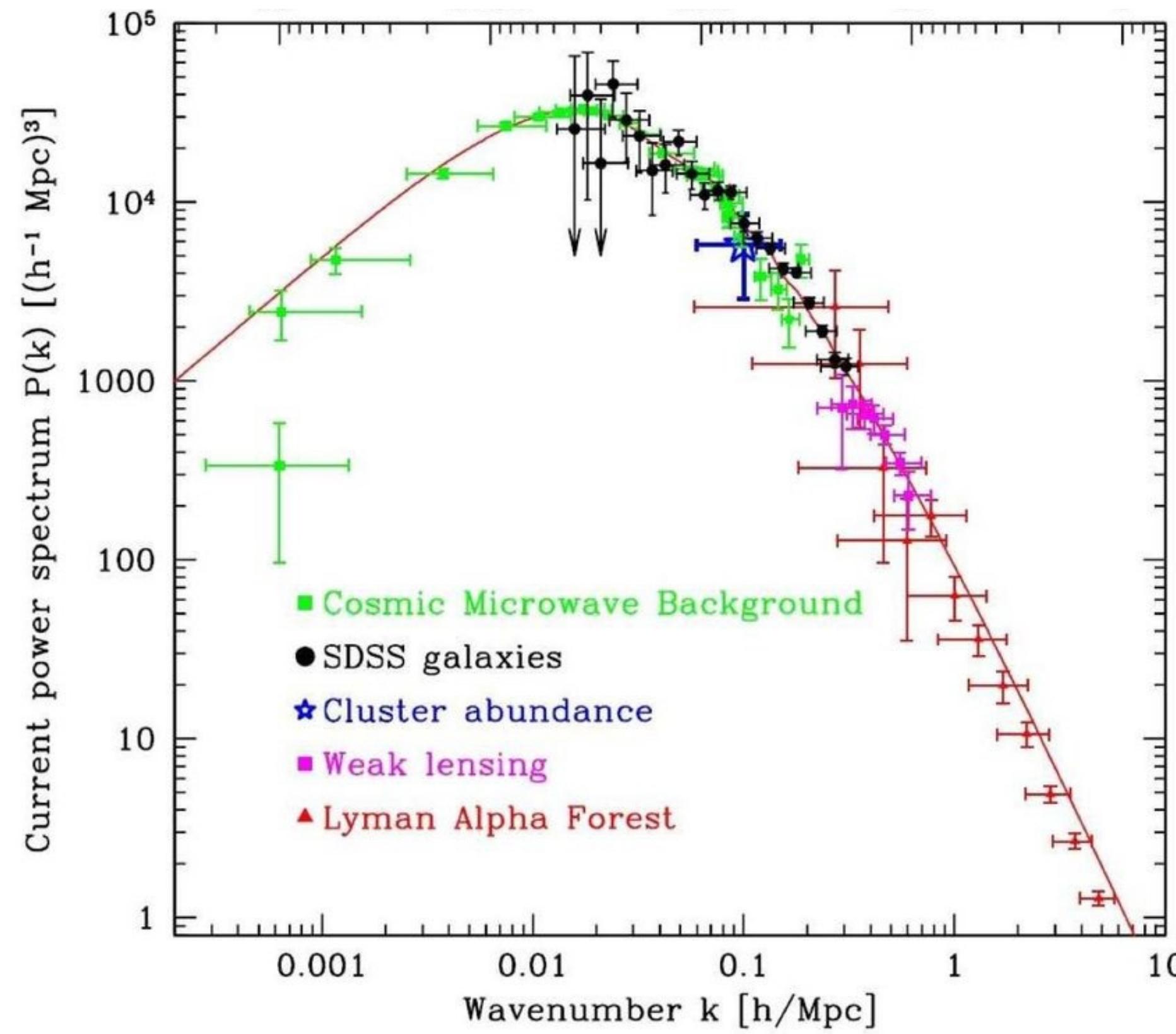
[Planck18]

## Primordial power spectrum

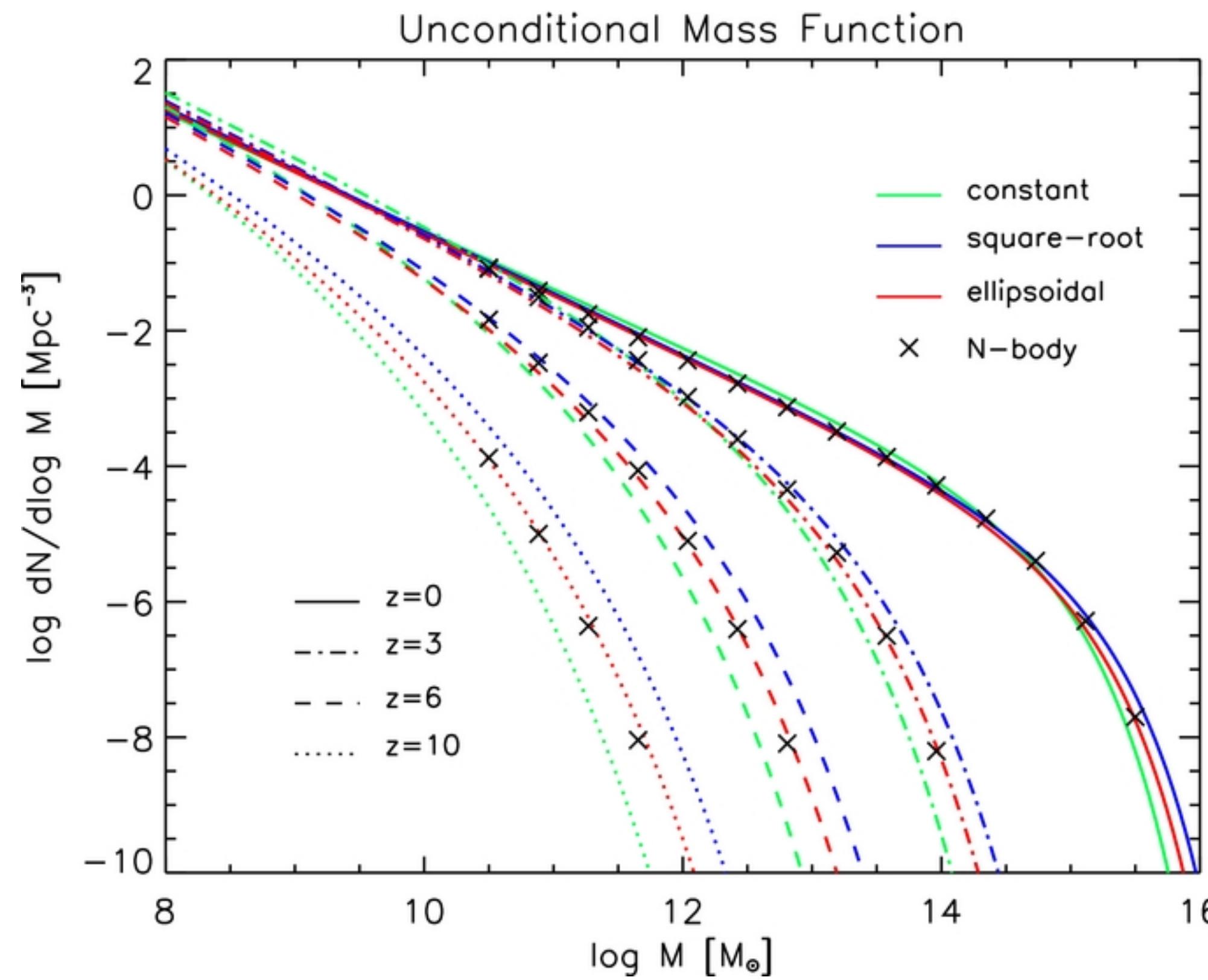
$$\mathcal{P}_{\mathcal{R}}(k) = \mathcal{A}_s \left( \frac{k}{k_0} \right)^{n_s - 1}$$

# Linear matter power spectrum

$$\mathcal{P}_m(k, z) = \frac{4}{25} \left[ \frac{D_1(z) k^2}{\Omega_{m,0} H_0^2} T(k) \right]^2 \mathcal{P}_{\mathcal{R}}(k)$$

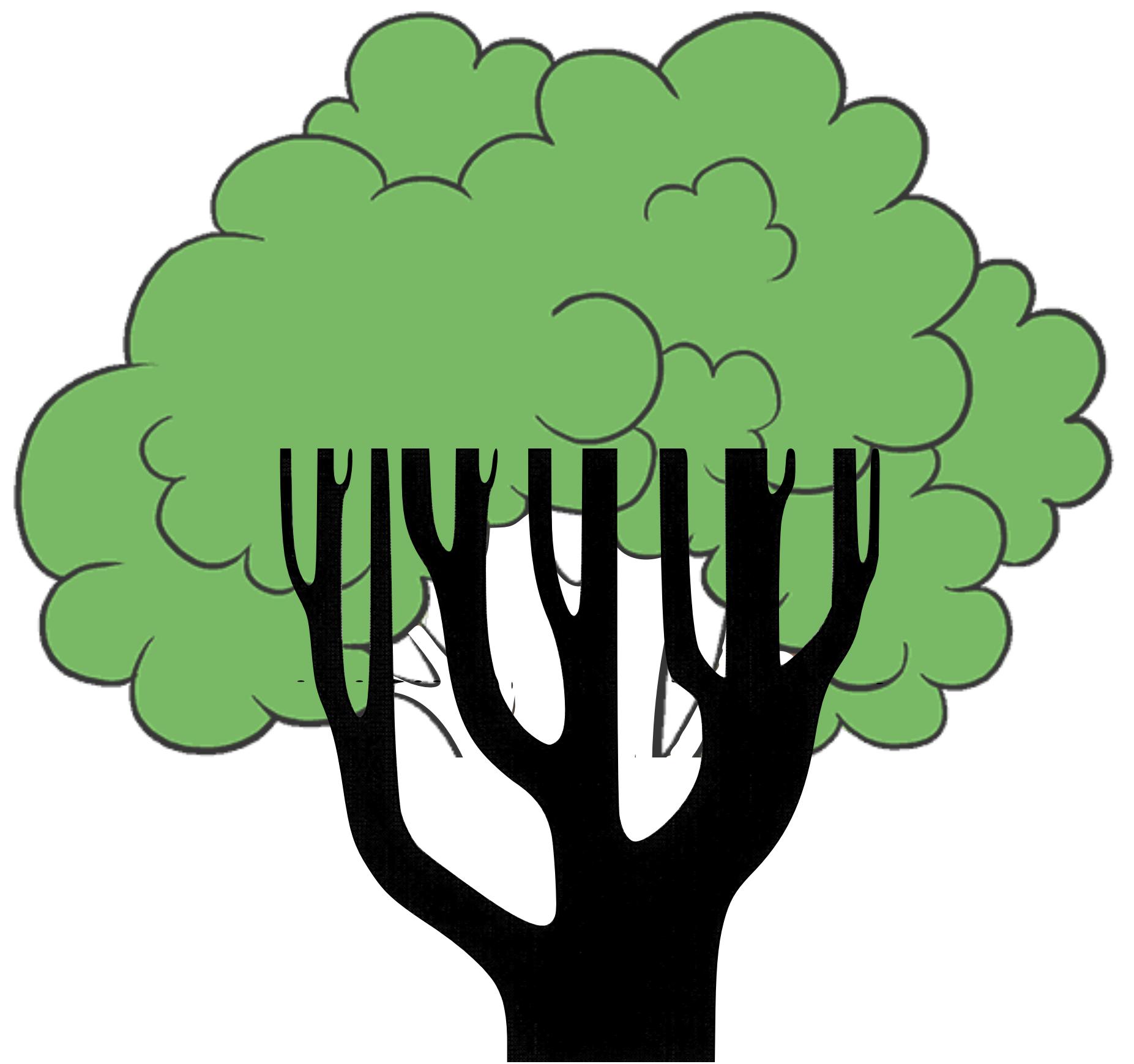


[Tegmark+04]



Excursion set theory gives the  
halo mass function

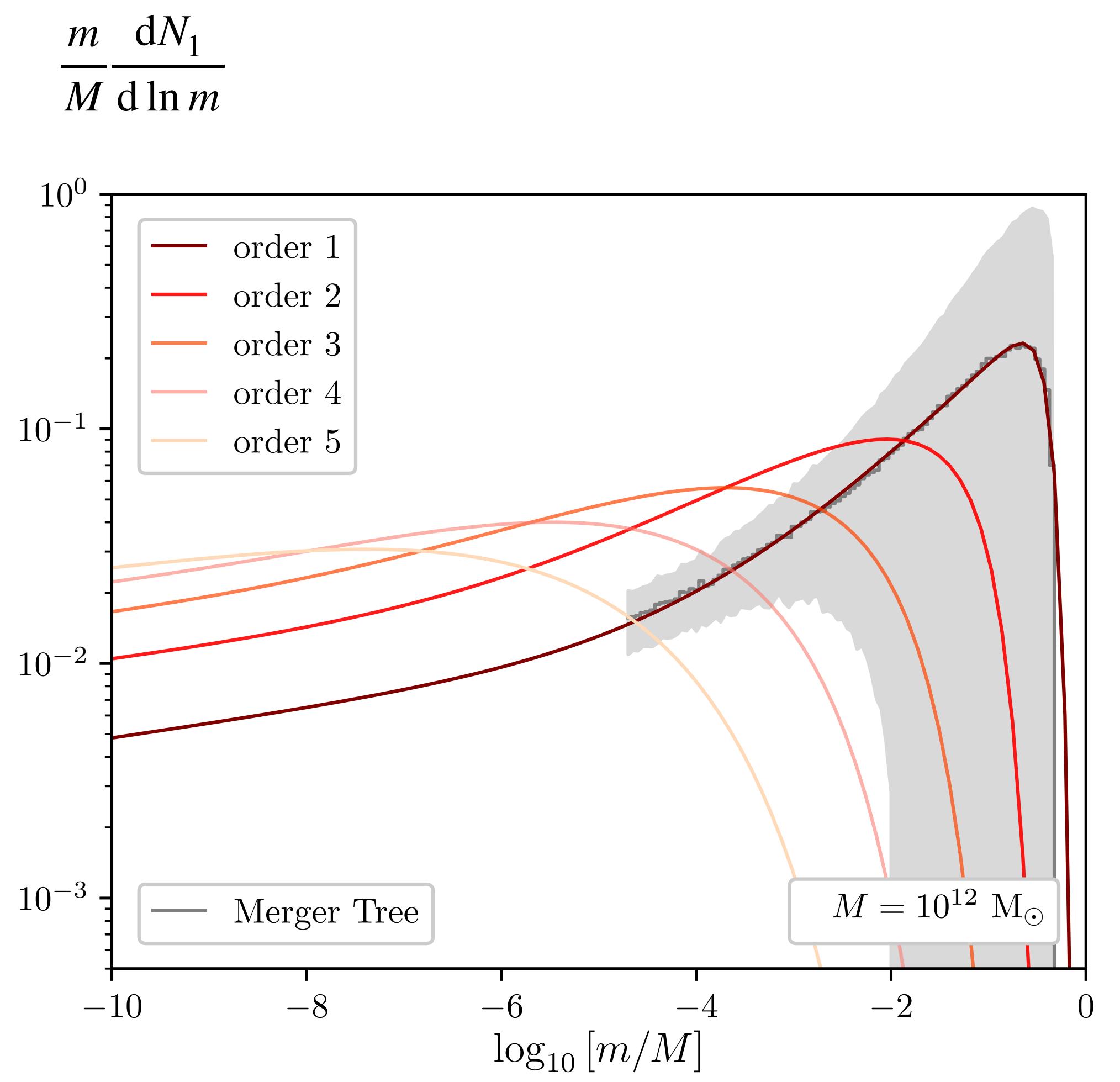
$$\frac{dn(M)}{dM} = \frac{\bar{\rho}_{m,0}}{M} \frac{\nu(M)}{2S(M)} \left| \frac{dS}{dM} \right| f[\nu(M)]$$



[image from Lacey+93]

Excursion set formalism +  
merger tree algorithms give  
the « cosmological »  
subhalo mass function

$$\frac{dn(m \mid M)}{dm}$$



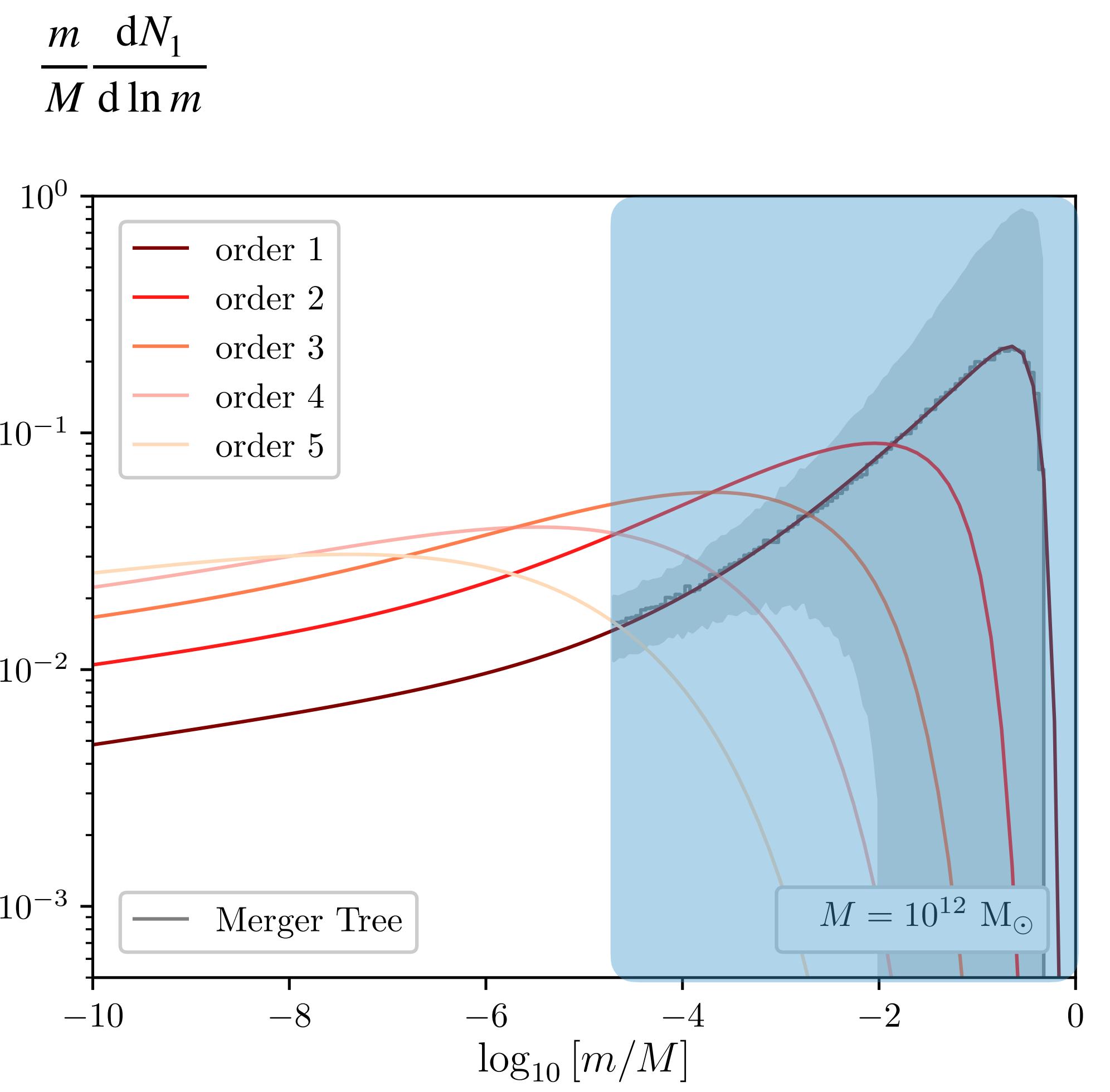
# We fit the subhalo mass function at z=0

Run the Cole+00 algorithm  
gives the mass function at large mass

Fit with the function

$$f(m, M) = \frac{1}{m} \left[ \sum_{i=1,2} \gamma_i \left( \frac{m}{M} \right)^{-\alpha_i} \right] \exp \left\{ -\beta \left( \frac{m}{M} \right)^{\zeta} \right\}$$

[GF and Lavalle (in prep.),  
see also e.g., Jiang+14, Giocolo+08]

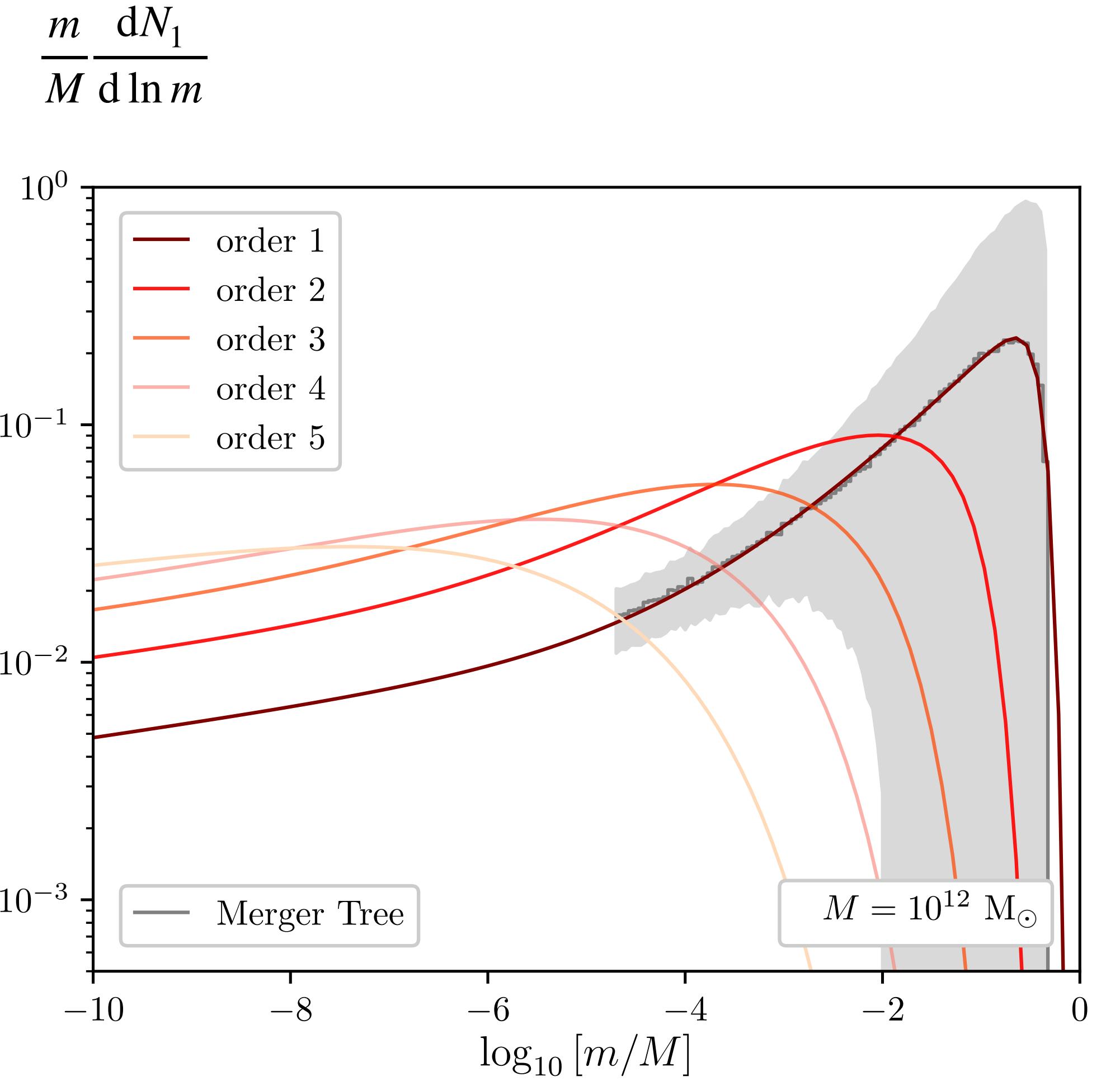


We fit the subhalo mass function at z=0

But ....

mass function at small mass inferred only from the behaviour at large mass

[GF and Lavalle (in prep.),  
see also e.g., Jiang+14, Giocolo+08]



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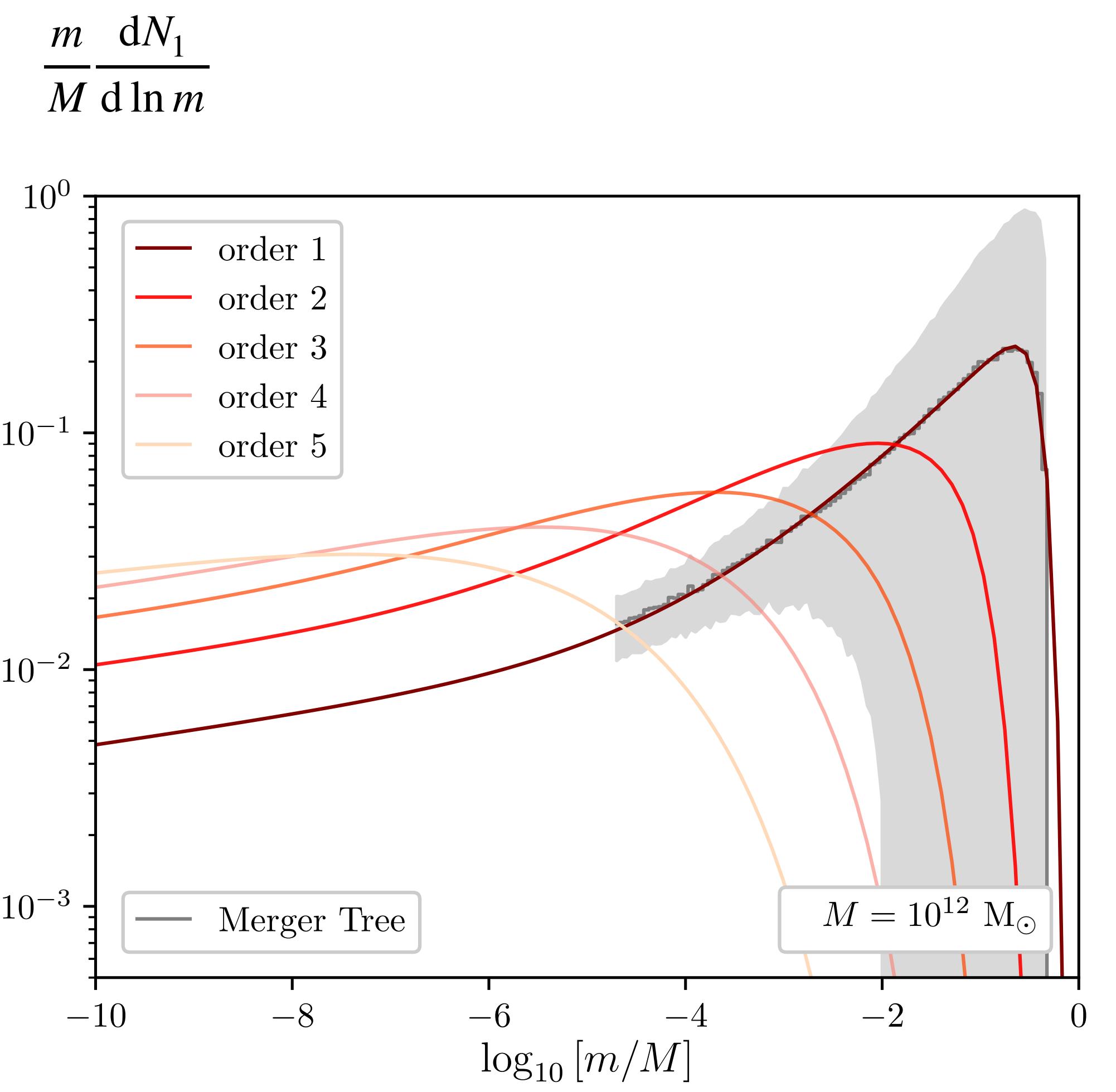
# We fit the subhalo mass function at z=0

Introduce a specific fitting procedure

Constrain the fit with the condition:

$$\frac{1}{M} \int_0^M m \frac{dN_1}{dm} dm = 1$$

The host halo is entirely made of subhalos (fractal picture)



# We fit the subhalo mass function at z=0

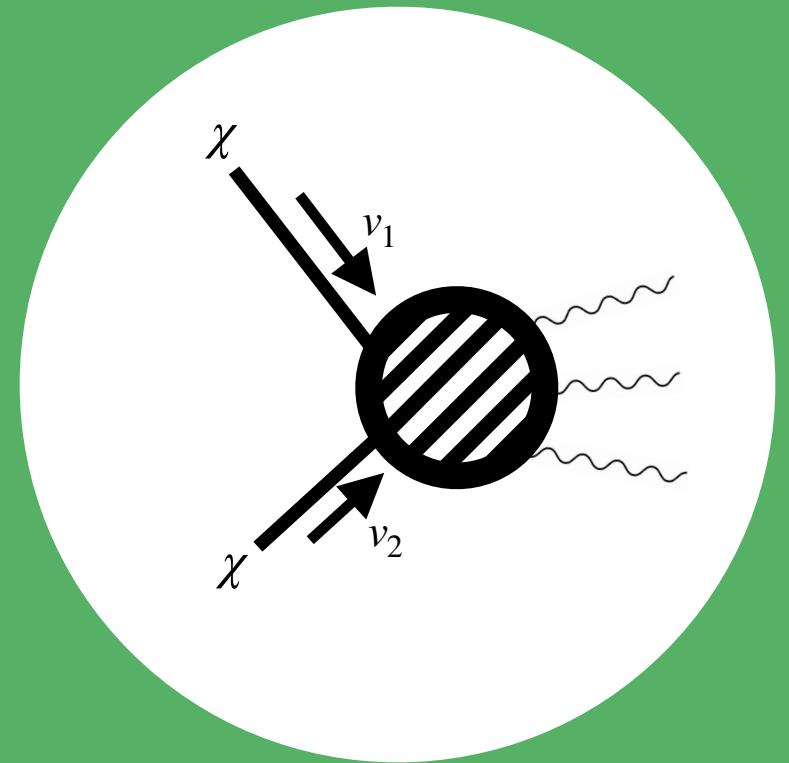
The constraint fixes the low-mass behavior

$$\frac{dN_1}{dm} \sim \gamma m^{-\alpha} \quad \text{with} \quad \alpha \sim 1.95$$

[GF and Lavalle (in prep.),  
see also e.g., Jiang+14, Giocolo+08]

**Analytical models are fast and easily adaptable  
to any host, at any redshift and to any cosmology**

**Almost the perfect tool to compute boost factors!  
(drawback: theoretical uncertainties)**

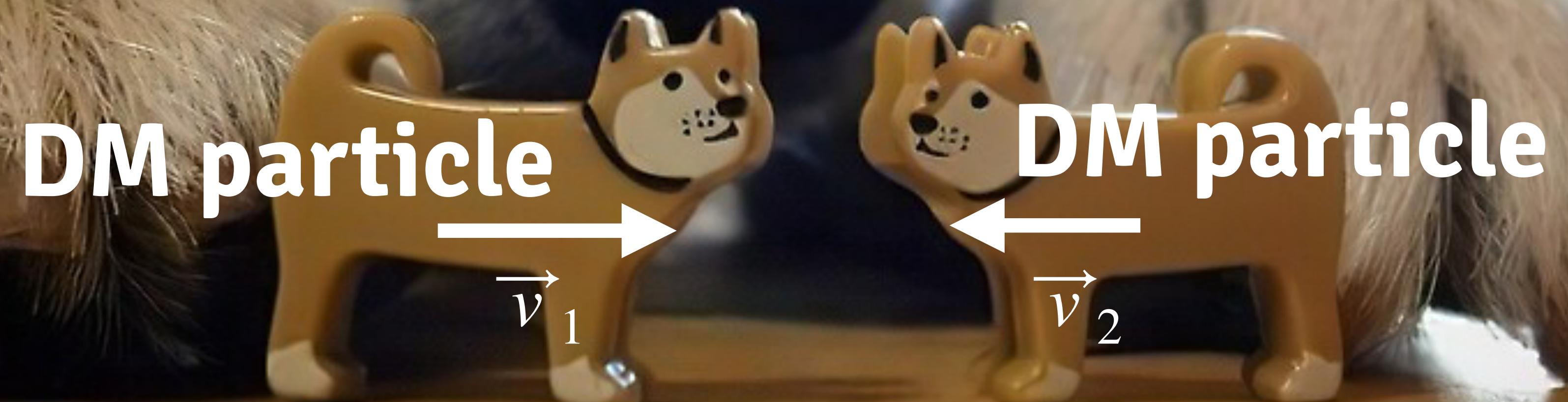


# DM Halo

## Part 2:

Sommerfeld enhanced  
and subhalo boosted  
J-factors

[2007.10392, 2203.16440]



We want

to classify astrophysical targets

In terms of their potential of DM detection  
(using our analytical model)...

... in scenarios with  
Sommerfeld enhancement  
and accounting for the subhalo boost

# Enhancement of the cross-section

Light mediators  $\phi$  can modify the cross-section in the non-relativistic limit

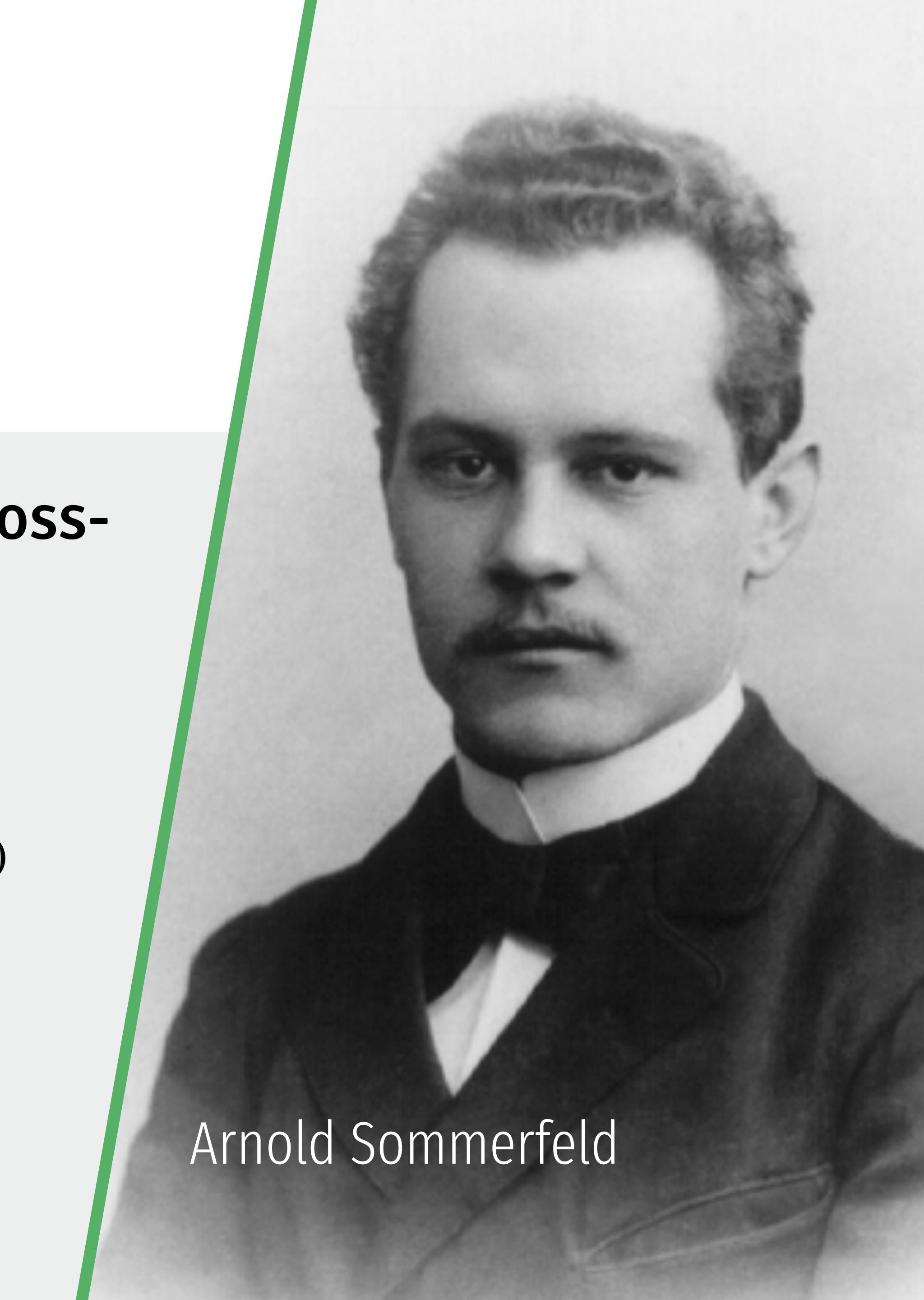
Need to solve Schrodinger's equation:

$$\left[ -\frac{1}{2mr} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{l(l+1)}{2mr^2} + V(r) \right] \psi_l(r) = E\psi_l(r)$$



An attractive (Yukawa) potential enhances the cross-section

Arnold Sommerfeld



# Enhancement of the cross-section

$$\sigma_{\text{ann}} v_{\text{rel}} = \sigma_0 + \sigma_1 v_{\text{rel}}^2$$



$$\sigma_{\text{ann}} v_{\text{rel}} = S_0(v_{\text{rel}}) \sigma_0 + S_1(v_{\text{rel}}) \sigma_1 v_{\text{rel}}^2$$

- To obtain analytic results approximate the Yukawa potential by the Hulthen potential:

$$V_H(r) = -\frac{\alpha m_* e^{-m_* r}}{1 - e^{-m_* r}} \rightarrow S_l(\epsilon_\phi, \epsilon_\nu) \quad m_* = \frac{\pi^2}{6} m_\phi$$

Arnold Sommerfeld



# Enhancement of the cross-section

Two important parameters:

$$\epsilon_\phi = \frac{m_\phi}{\alpha_\chi m_\chi}$$

Particle physics

$$\epsilon_\nu = \frac{\nu_{\text{rel}}}{2\alpha_\chi}$$

Astrophysics

With the coupling strength:

$$\alpha_\chi \equiv \frac{g_\chi^2}{4\pi}$$

Arnold Sommerfeld



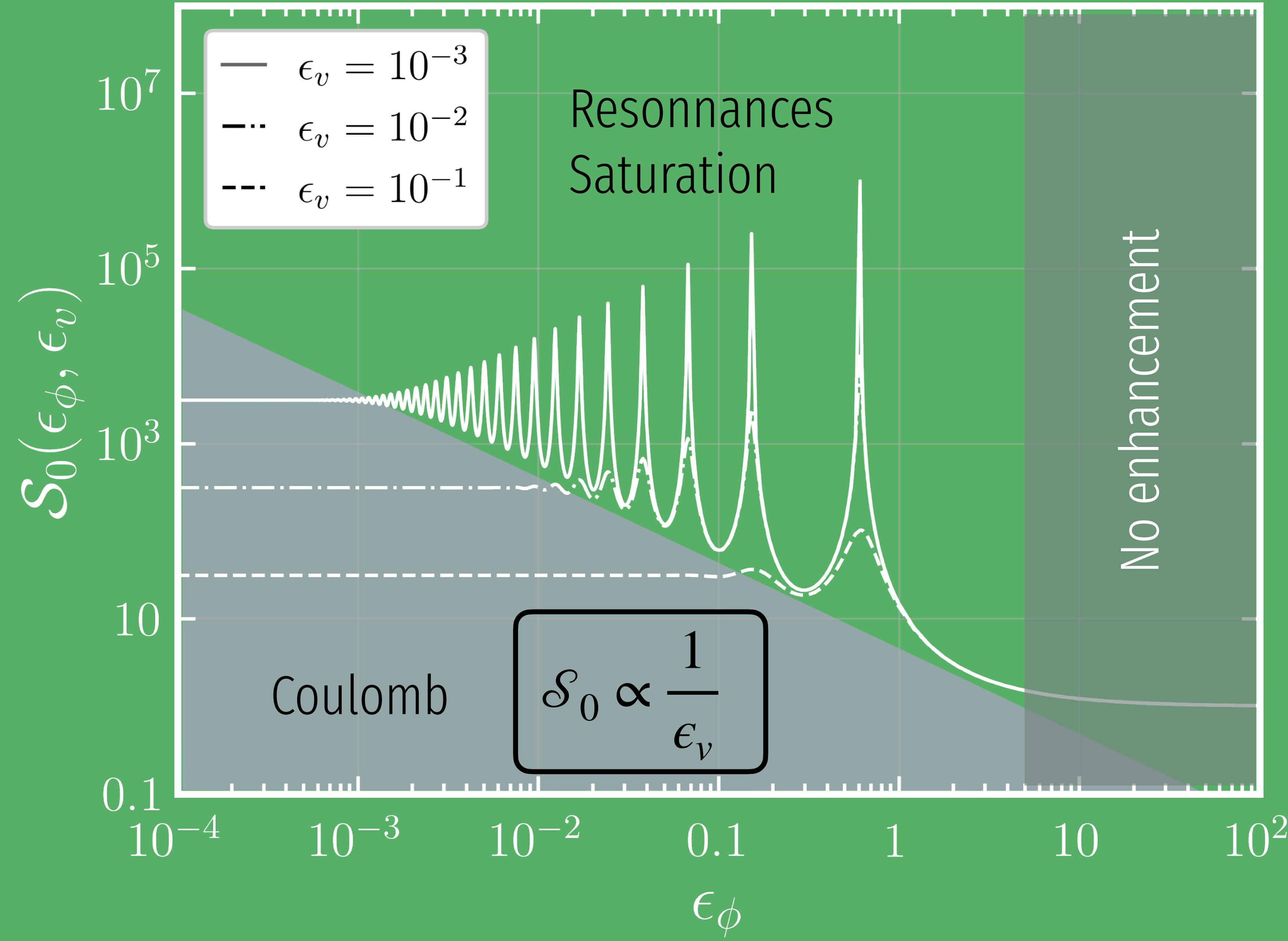
Regroup velocity dependent terms:

$$\mathcal{S}_l(\epsilon_\phi, v_{\text{rel}}) \equiv v_{\text{rel}}^{2l} S_l(\epsilon_\phi, \epsilon_v)$$

$$\sigma_{\text{ann}} v_{\text{rel}} = \sum_{l=1,2,\dots} \sigma_l \mathcal{S}_l(\epsilon_\phi, v_{\text{rel}})$$

$$\epsilon_\phi = \frac{m_\phi}{\alpha_\chi m_\chi}$$

$$\epsilon_v = \frac{v_{\text{rel}}}{2\alpha_\chi}$$



$$\epsilon_\phi = \frac{m_\phi}{\alpha_\chi m_\chi}$$

$$\epsilon_\nu = \frac{v_{\text{rel}}}{2\alpha_\chi}$$

What are the possible  
interesting targets?

# What are the possible interesting targets?

Dwarf spheroidal galaxies (dSphs)

$D \sim 10 - 100$  kpc and  $m \sim 10^7 - 10^9 M_{\odot}$

Sculptor  
[ESO/Digitized Sky Survey 2]



# What are the possible interesting targets?

Dwarf irregular galaxies (dIrrs)

$D \sim 1 \text{ Mpc}$  and  $m \sim 10^7 - 10^{10} \text{ M}_\odot$

IC10



# What are the possible interesting targets?

## Galaxy clusters

$D > 10 \text{ Mpc}$  and  $m \sim 10^{14} - 10^{15} \text{ M}_\odot$

Coma cluster  
[NASA / JPL-Caltech / L. Jenkins]



# dSphs

[Bonnivard+15]
Reticulum 2 (ultra-faint)
Sculptor (classical)
Draco (classical)

# dIrrs

[Gammaldi+21]		
	D [kpc]	$m$ ( $=M_{200}$ ) $[10^{10} M_\odot]$
NGC6822	480	3.16
IC10	790	3.98
WLM	970	0.40

# Clusters

[Sánchez-Conde+11]		
	D [Mpc]	$m$ ( $=M_{200}$ ) $[10^{14} M_\odot]$
Coma	102.18	8.77-13.16
Fornax	20.35	0.51-0.61
Perseus	80.69	5.14-7.71

# From the mass distribution to the phase-space/velocity distribution

We rely on Eddington's inversion method

Assuming isotropy & spherical symmetry

$$f(r, v) = \frac{1}{2\sqrt{2}\pi^2} \frac{d}{d\mathcal{E}} \int_0^{\mathcal{E}} \frac{d\rho}{d\Psi'} \frac{1}{\sqrt{\mathcal{E} - \Psi'}} d\Psi'$$

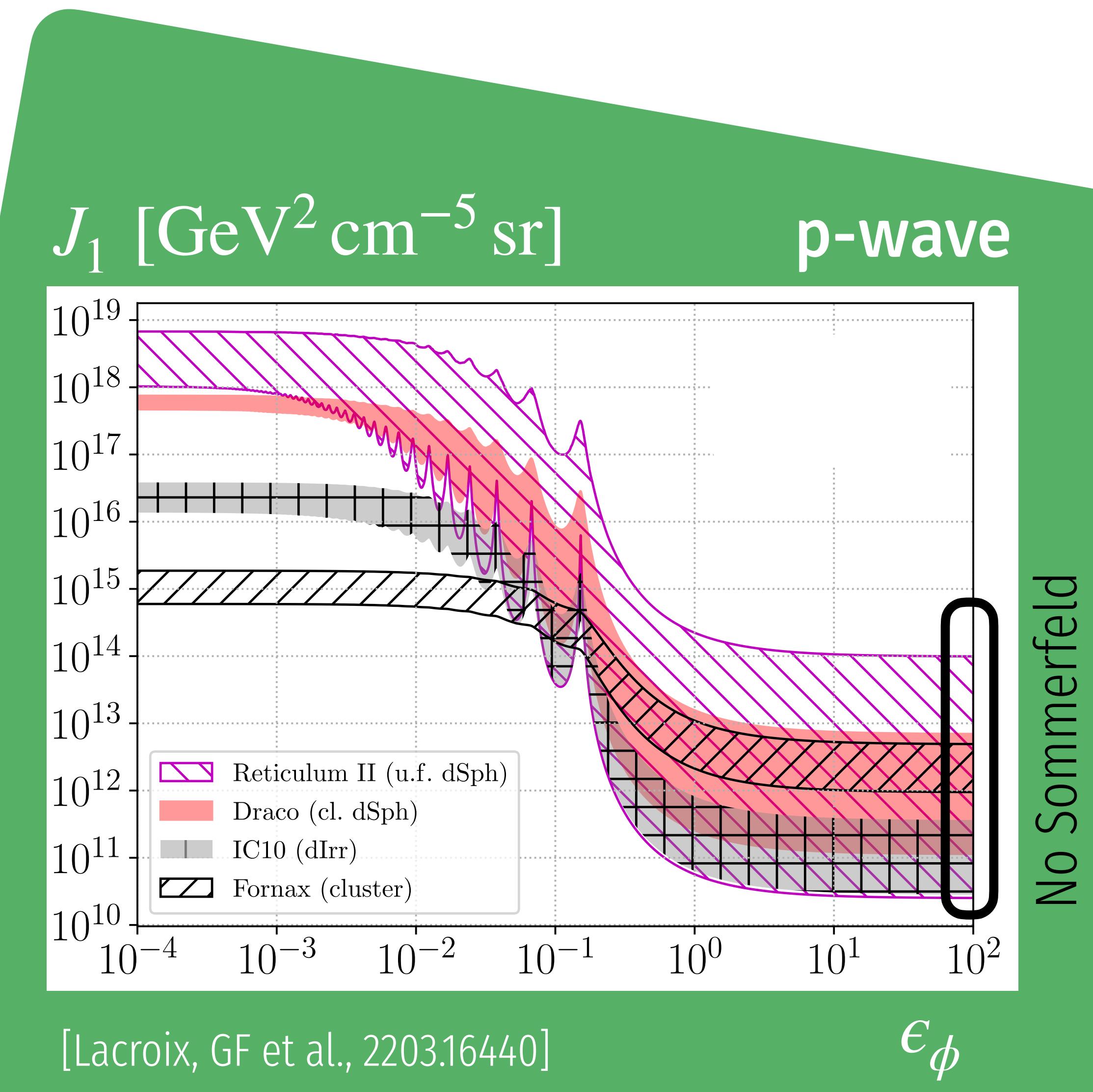
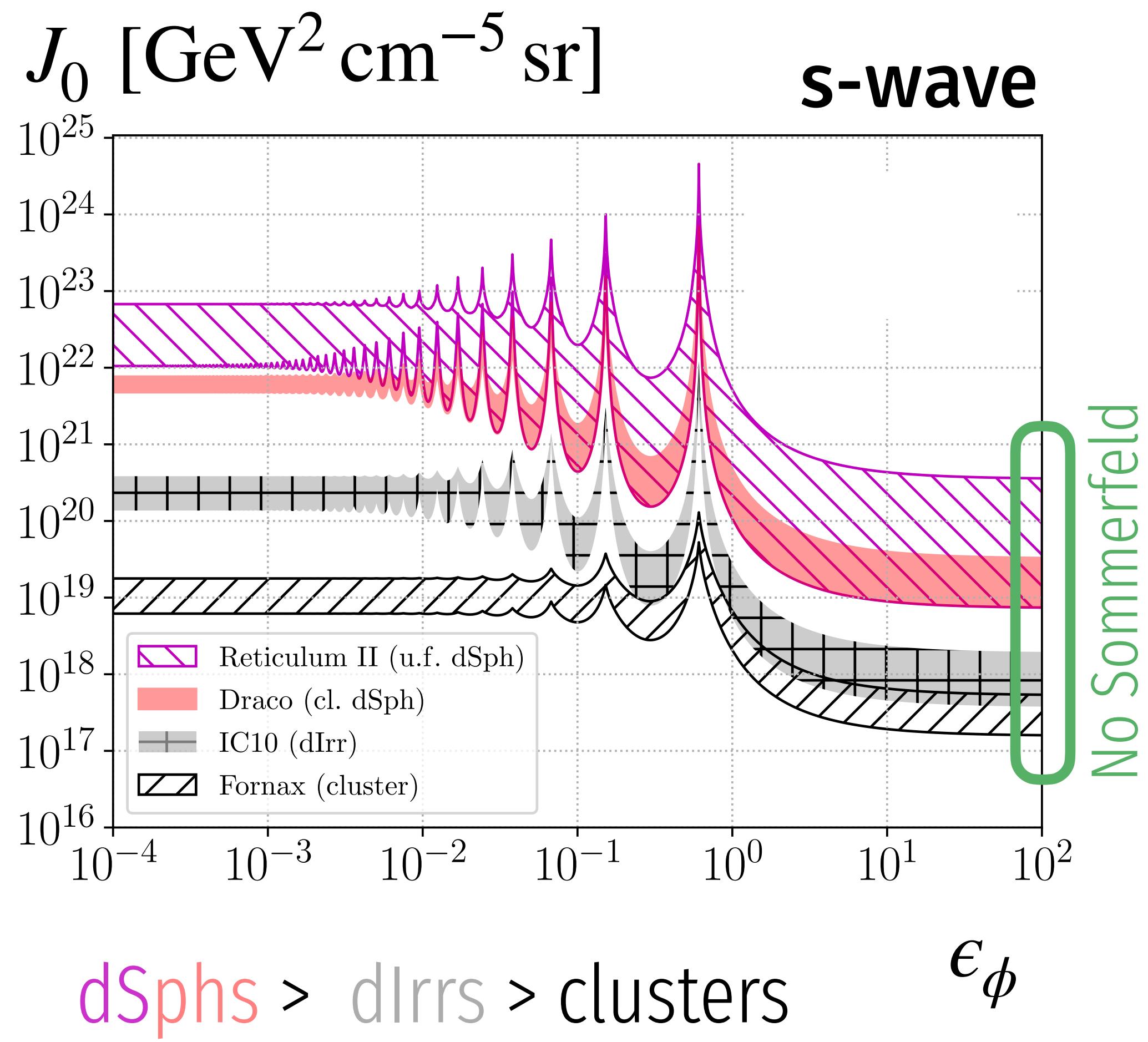
$$\mathcal{E}(r, v) = \Psi(r) - v^2/2$$

$$\Delta\Psi = -4\pi G_N \rho$$

Arthur Eddington



# Results for smooth halos



the smaller  
the structure,  
the lower  
the velocity

the impact of subhalos  
can be very important

# Boost factors

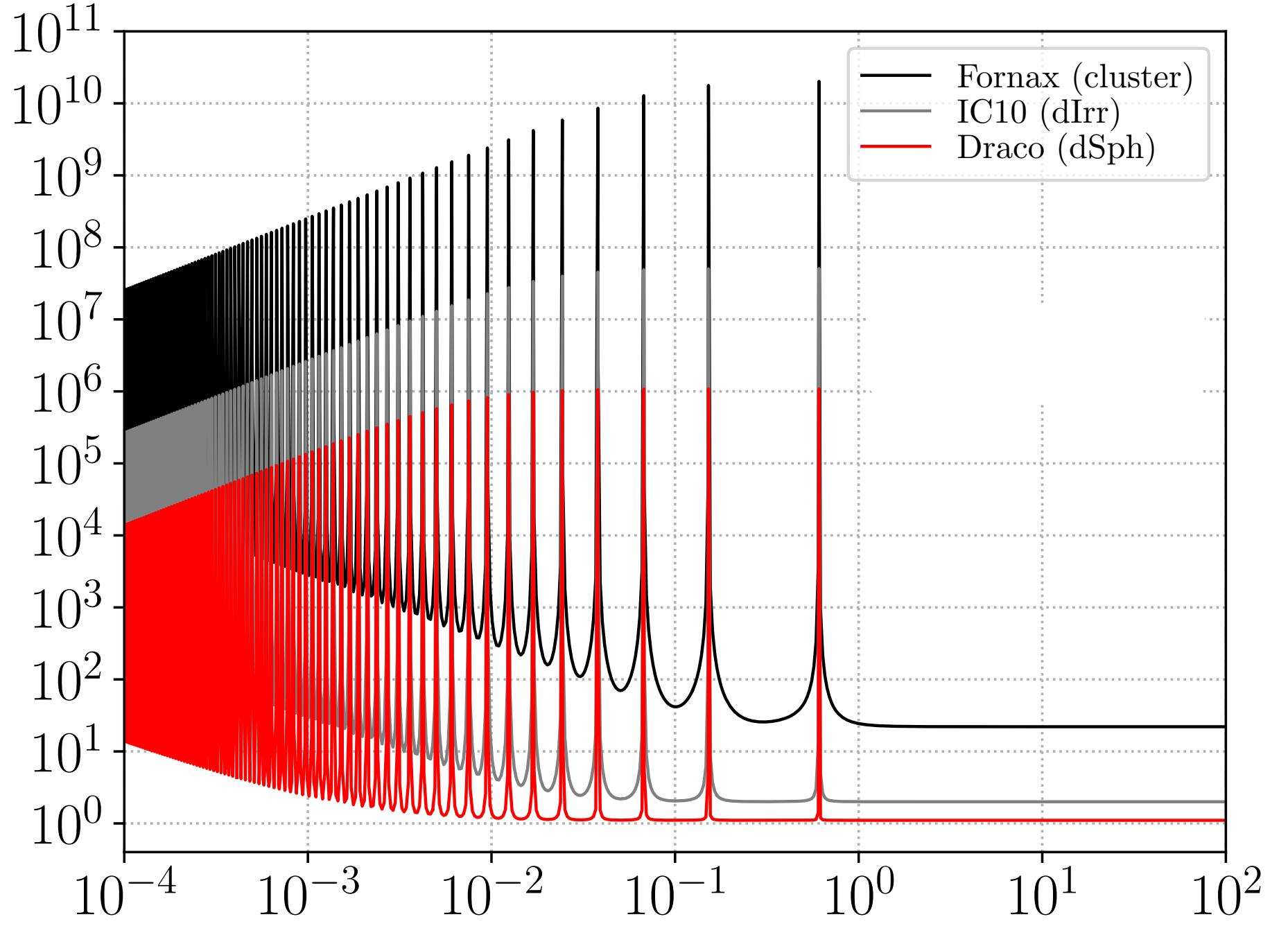
largest subhalo boost for clusters,  
lower subhalo boost for dSphs

cluster

dwarf spheroidal  
dwarf irregular

Boost

s-wave

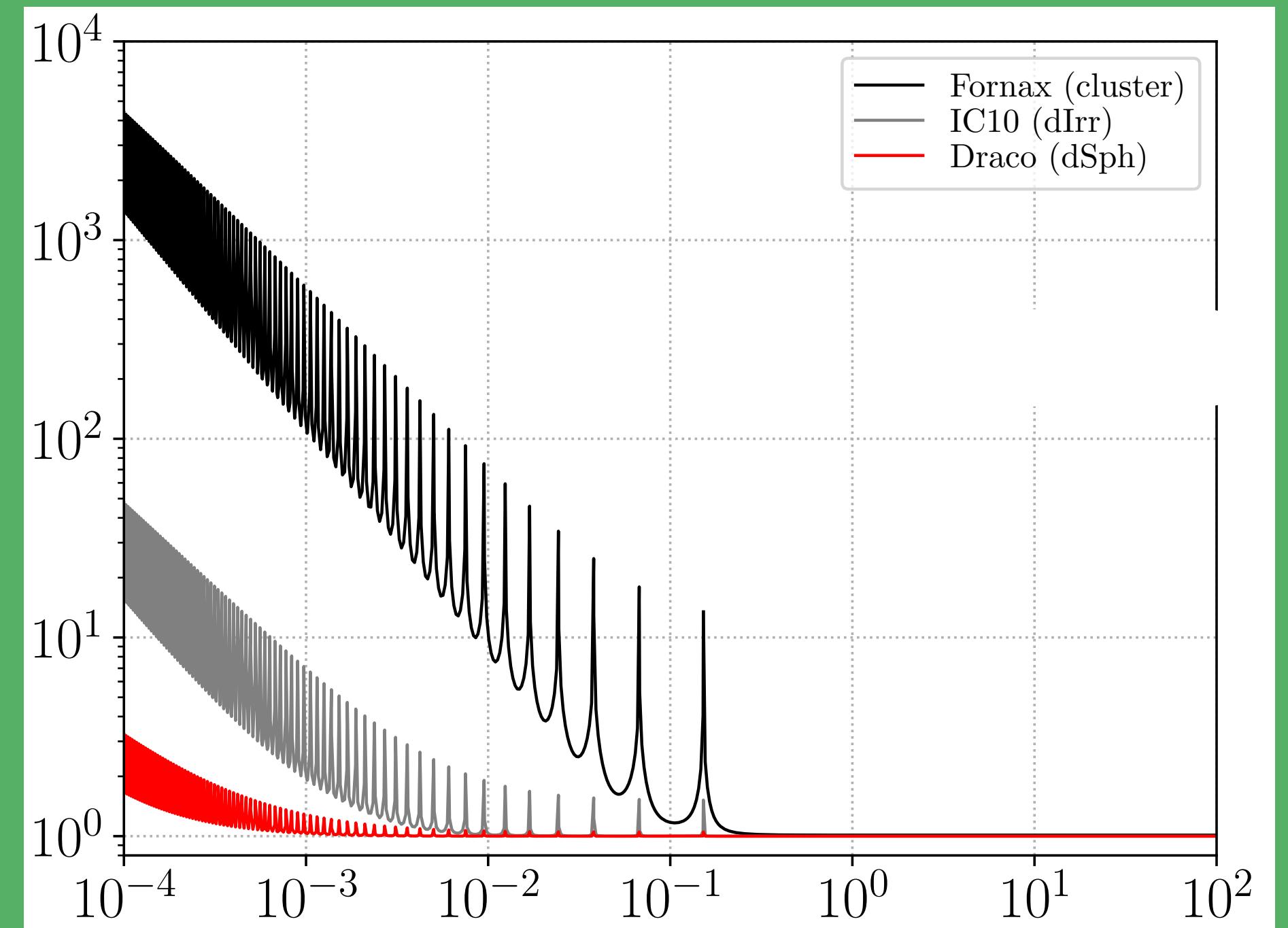


[Lacroix, GF et al., 2203.16440]

$\epsilon_\phi$

Boost

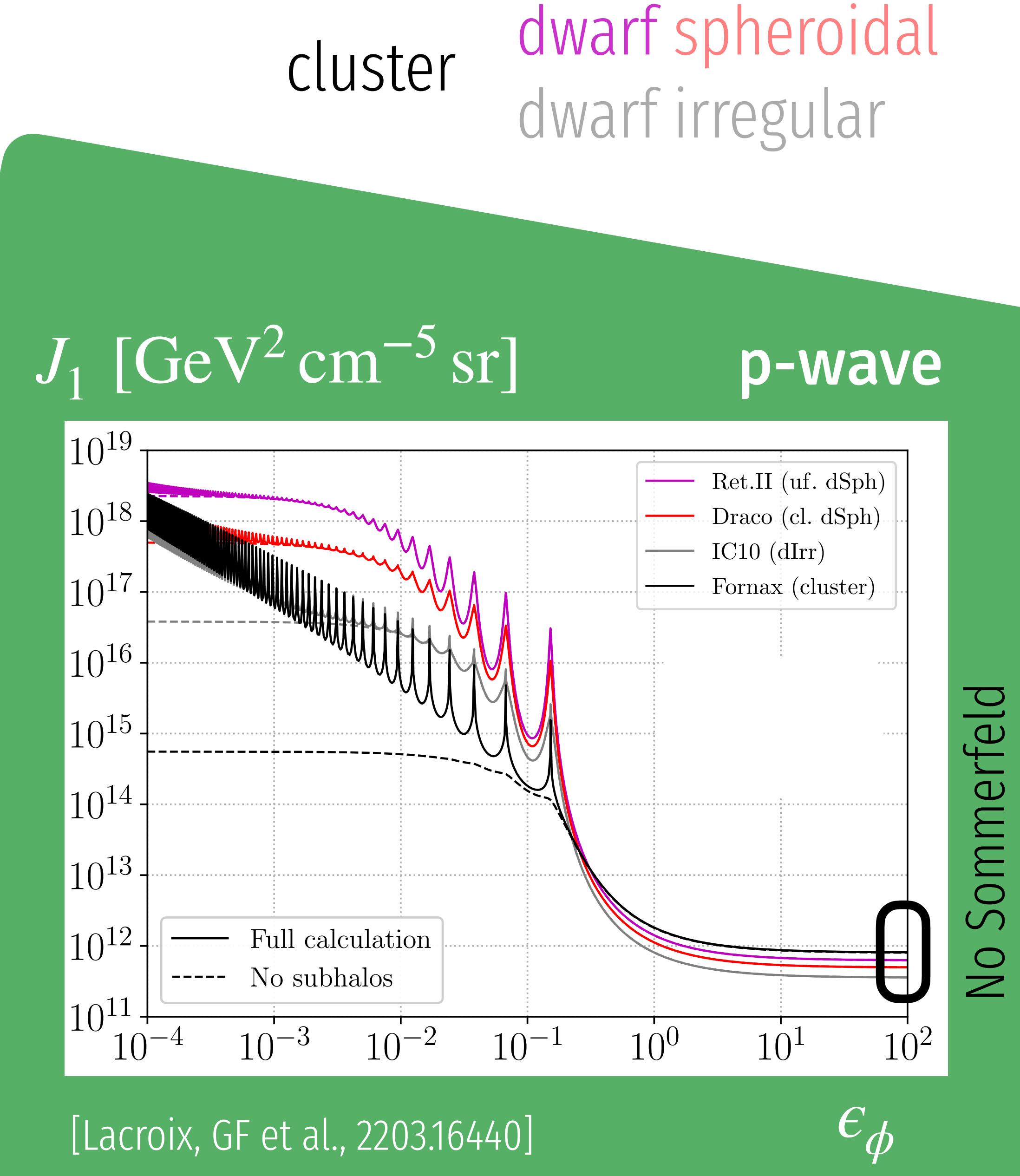
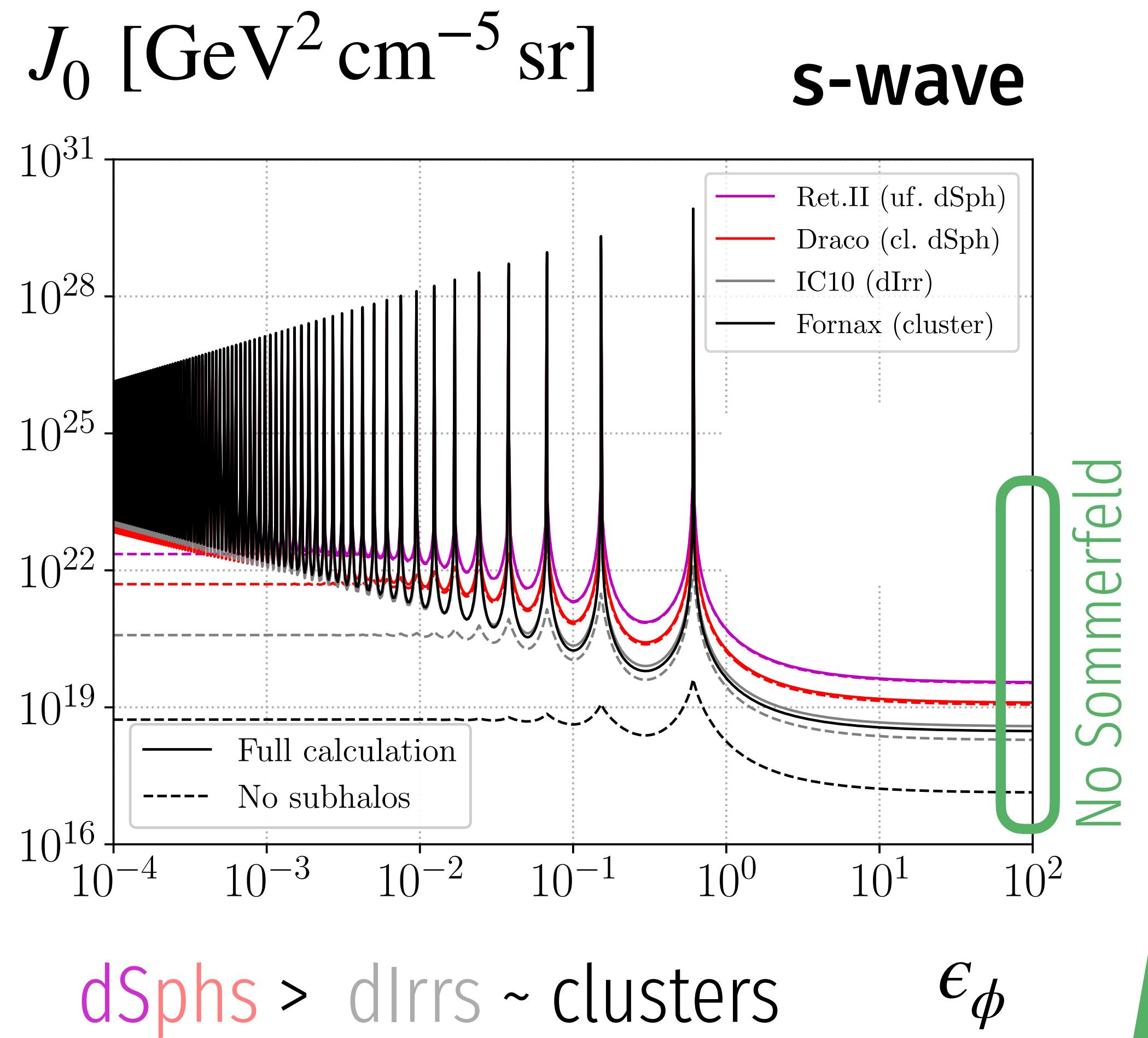
p-wave



[Lacroix, GF et al., 2203.16440]

$\epsilon_\phi$

# Results for full halos (with subhalos)



# Results for full halos (with subhalos)

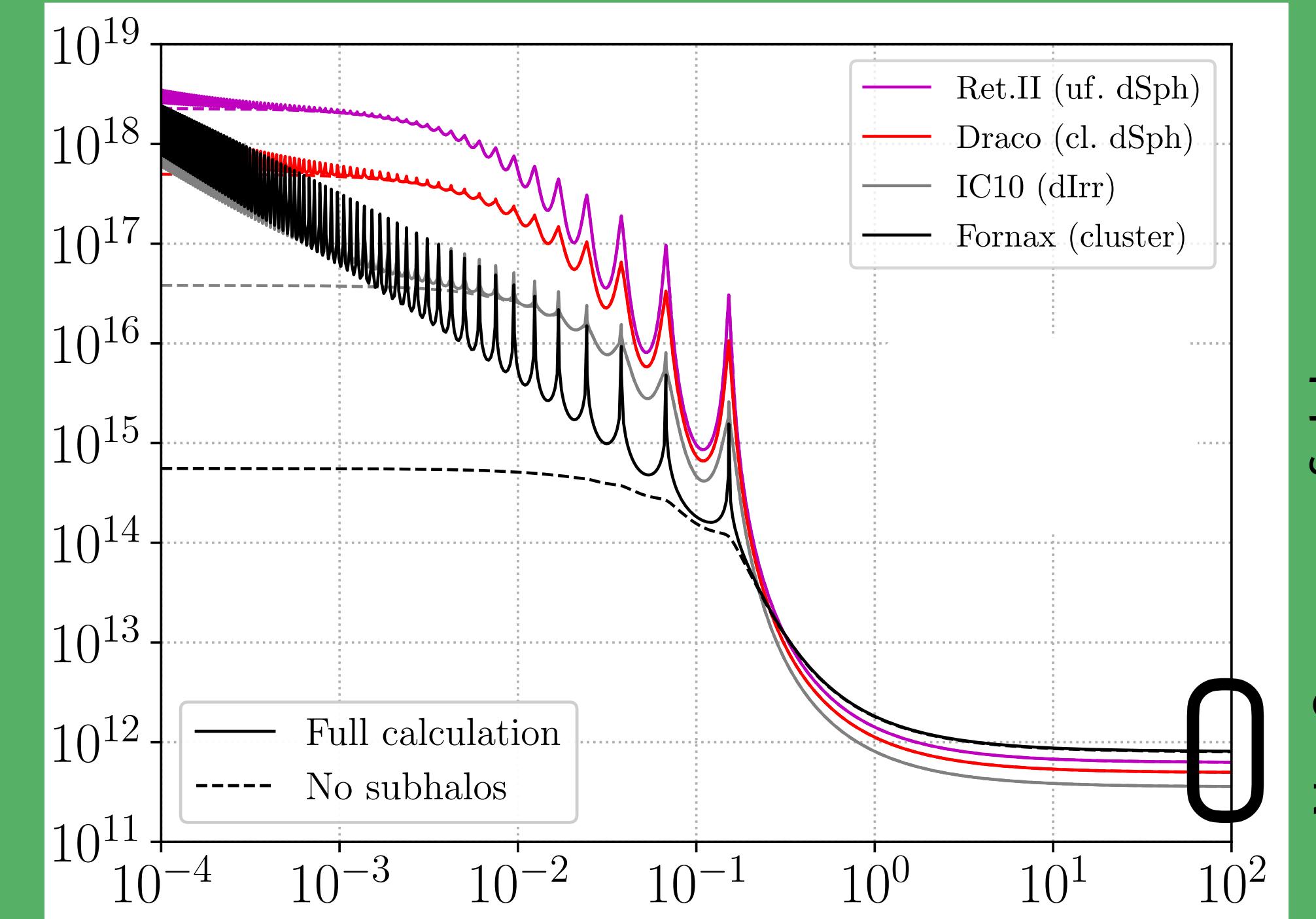
**p-wave:**  
Different hierarchies  
depending on  $\epsilon_\phi$

cluster

dwarf spheroidal  
dwarf irregular

$J_1$  [GeV $^2$  cm $^{-5}$  sr]

**p-wave**

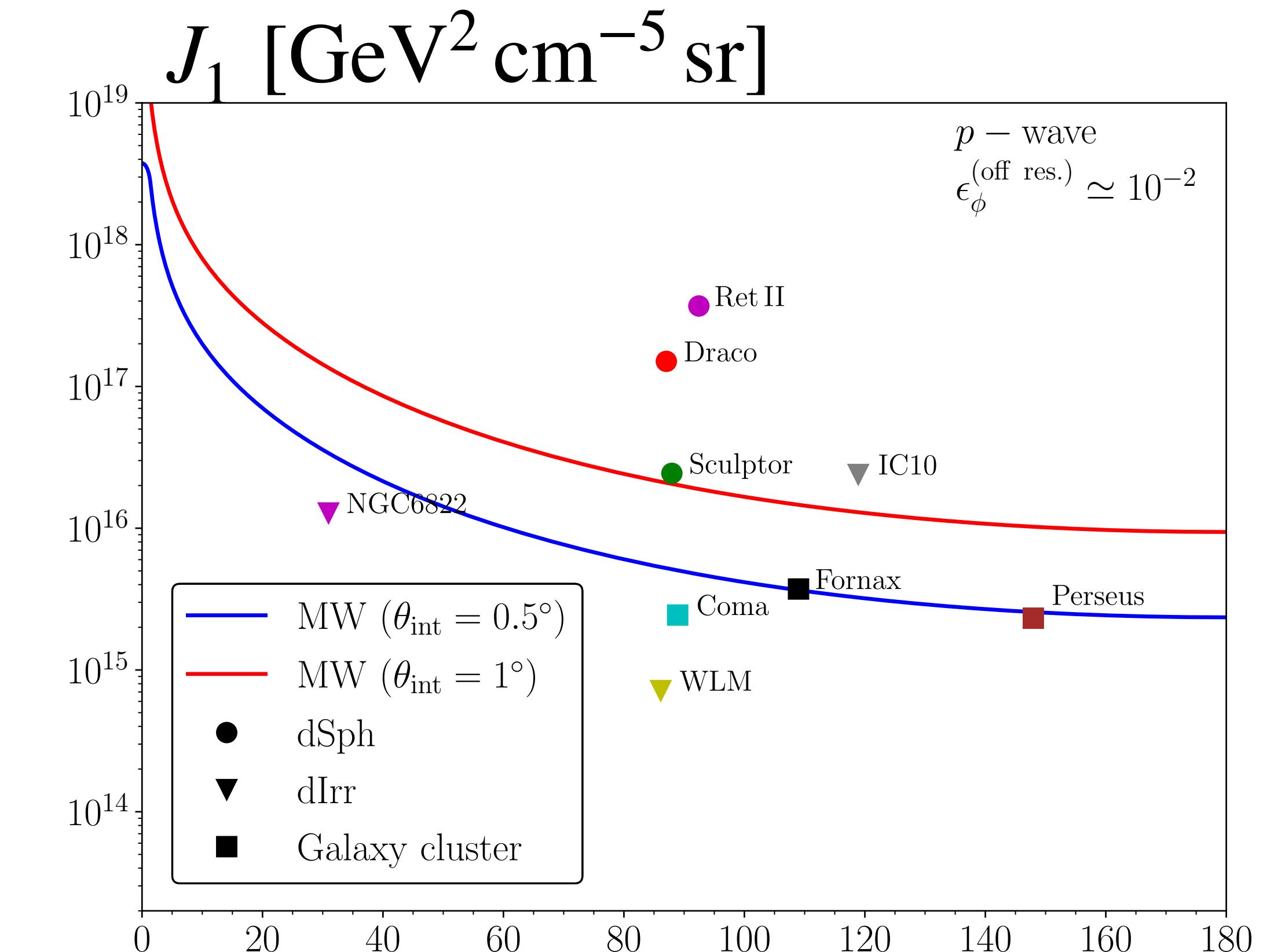
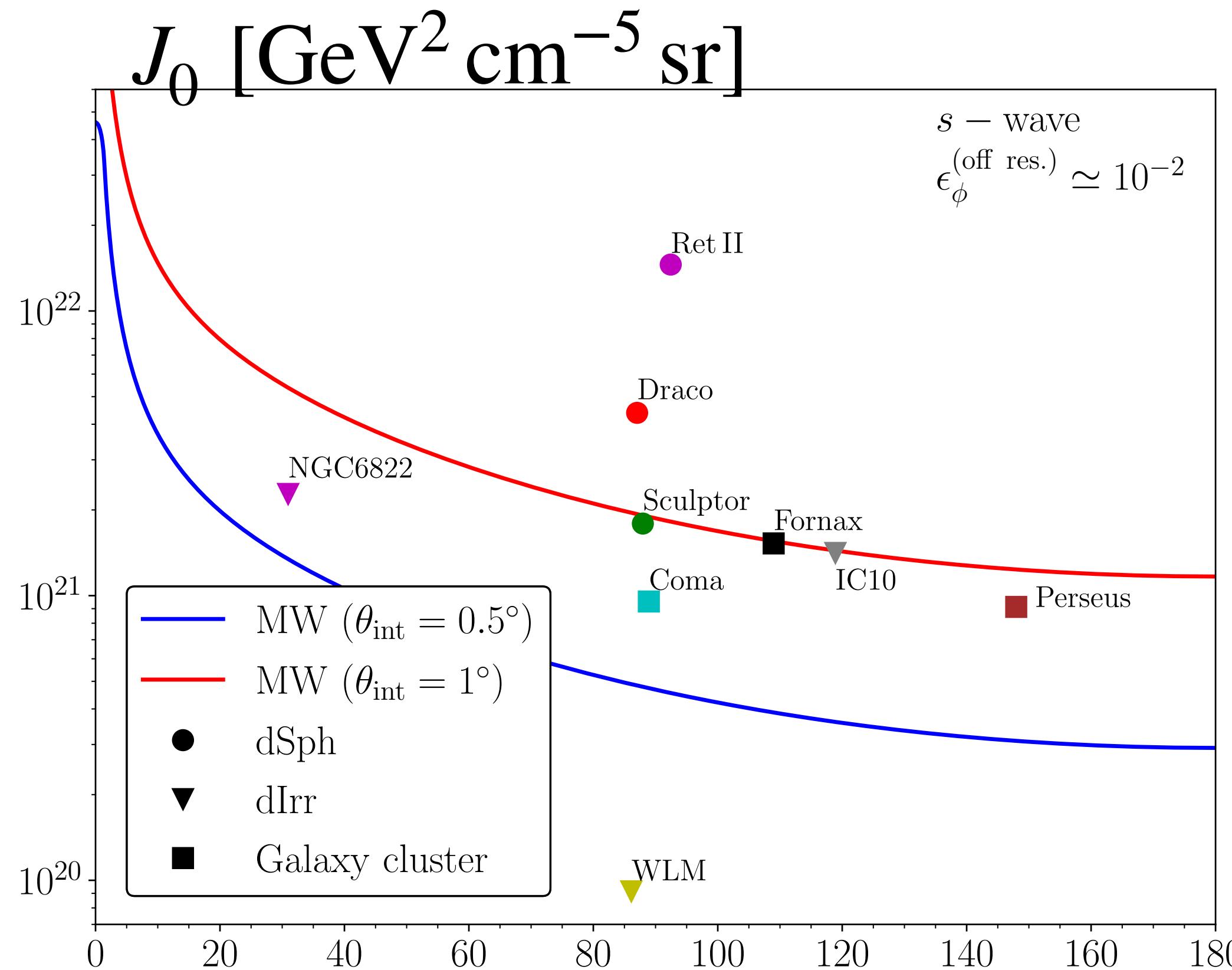


[Lacroix, GF et al., 2203.16440]

$\epsilon_\phi$

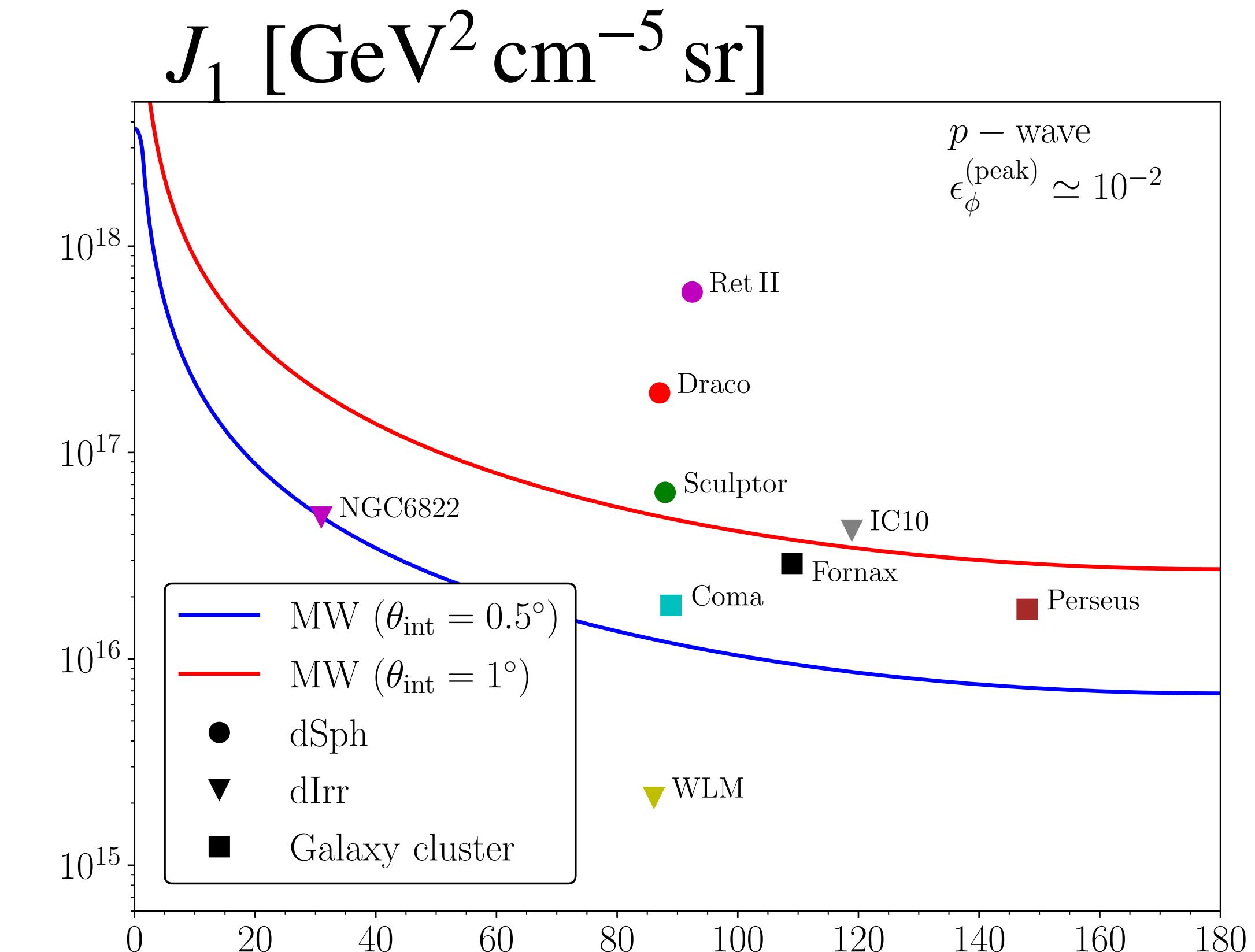
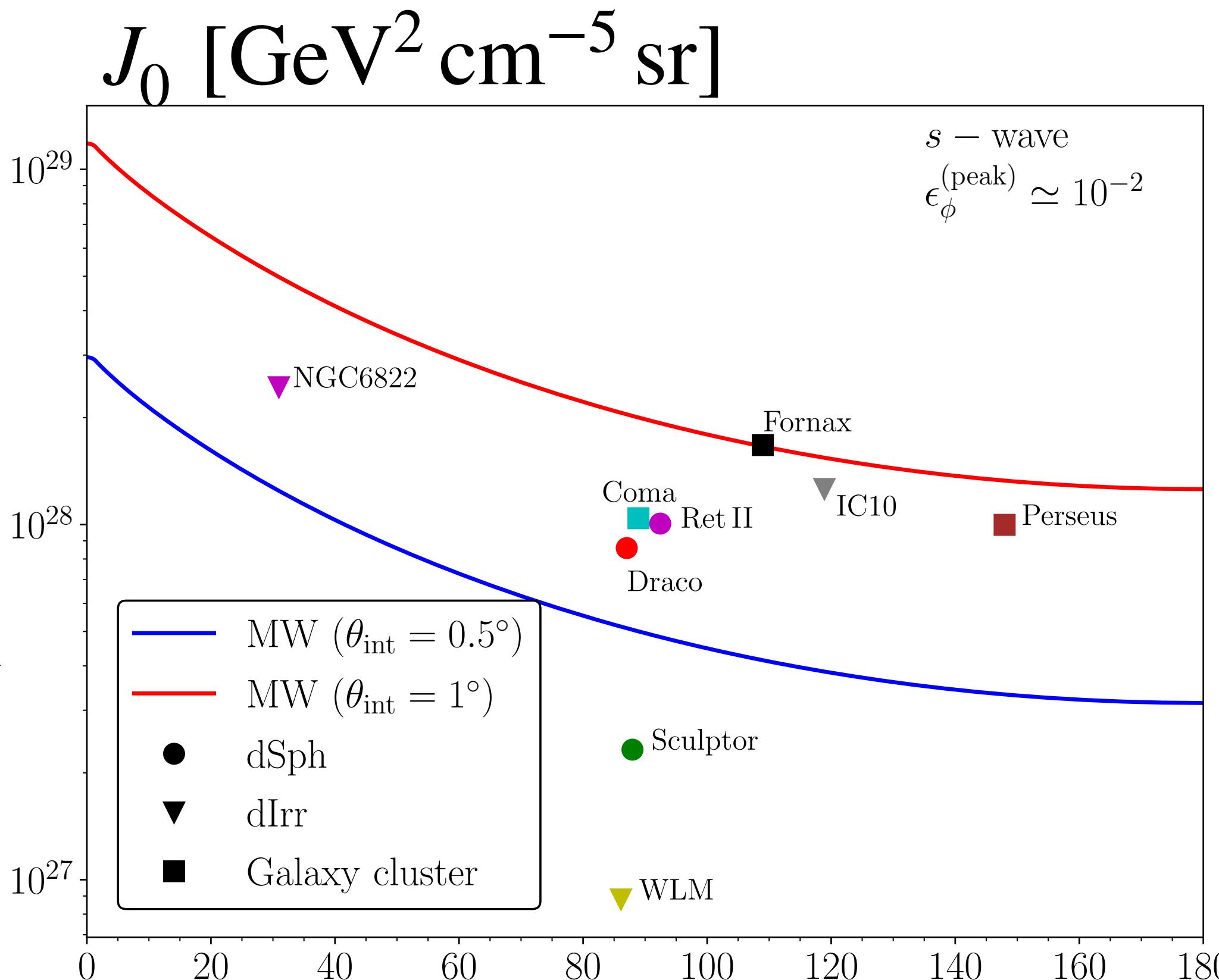
No Sommerfeld

# Comparison to the Milky-way background/foreground



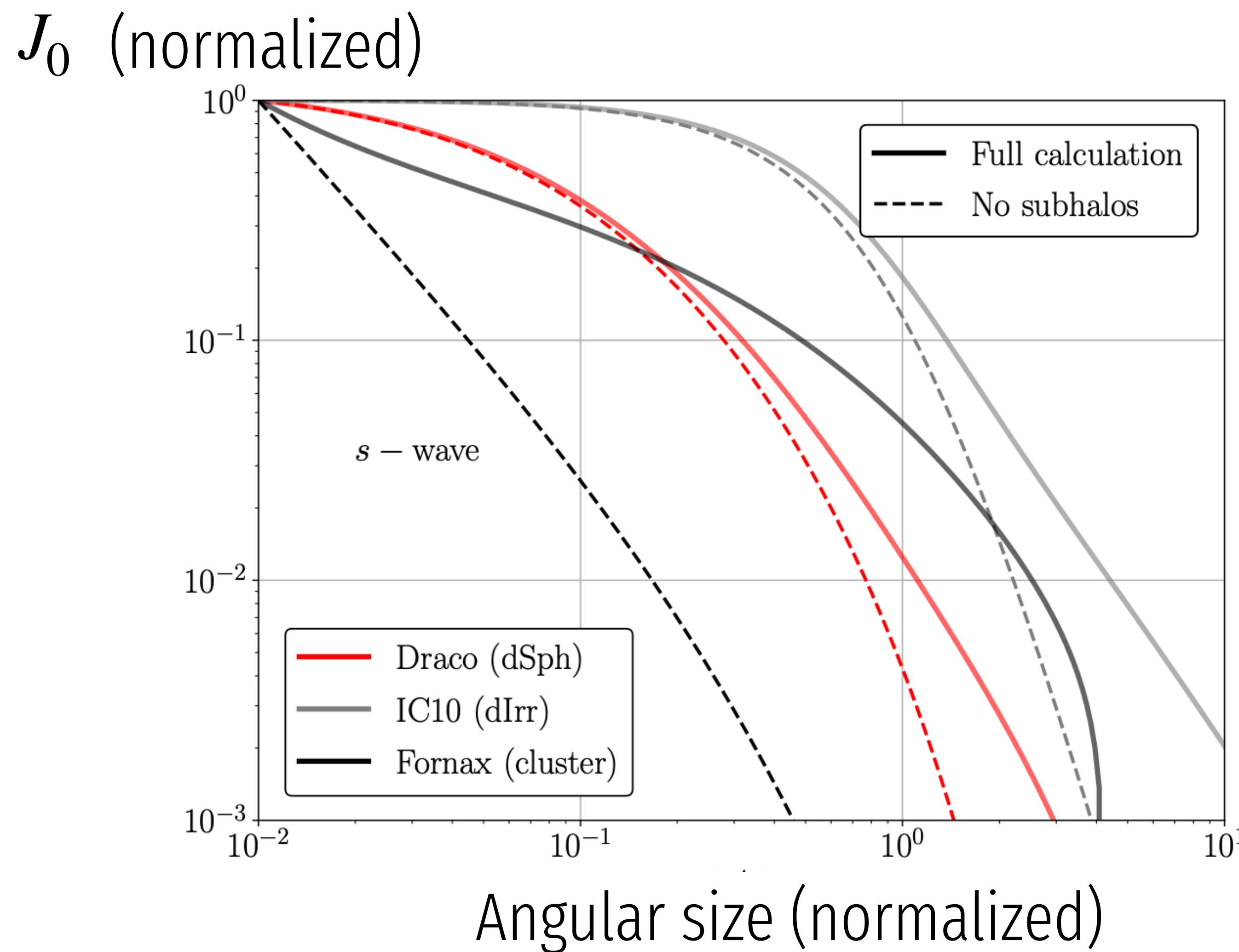
Angle from the Galactic center [deg]

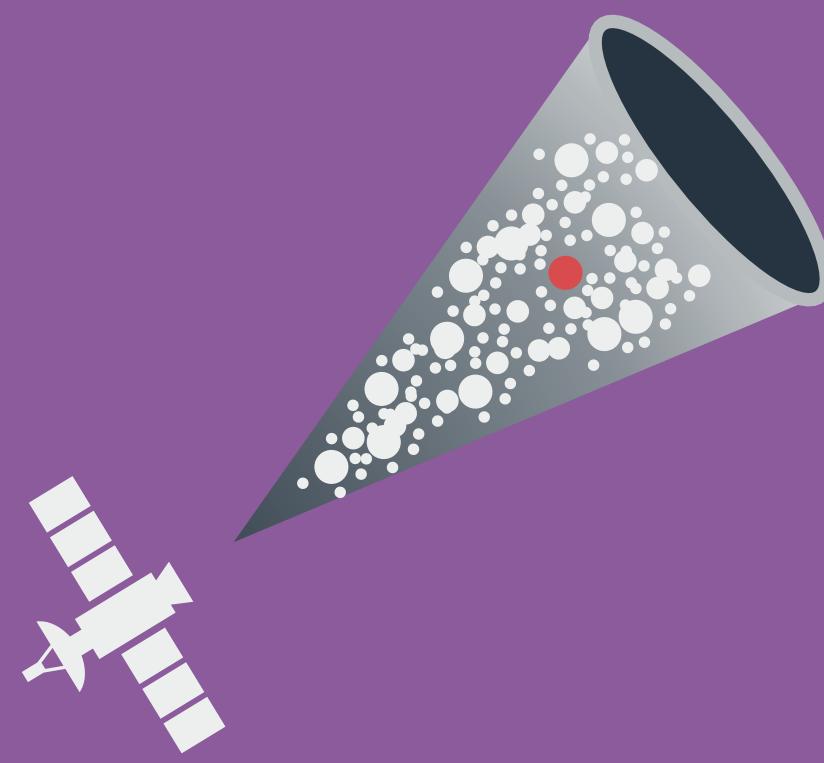
# Comparison to the Milky-way background/foreground



Angle from the Galactic center [deg]

# Subhalos can also enlarge the apparent size of the source





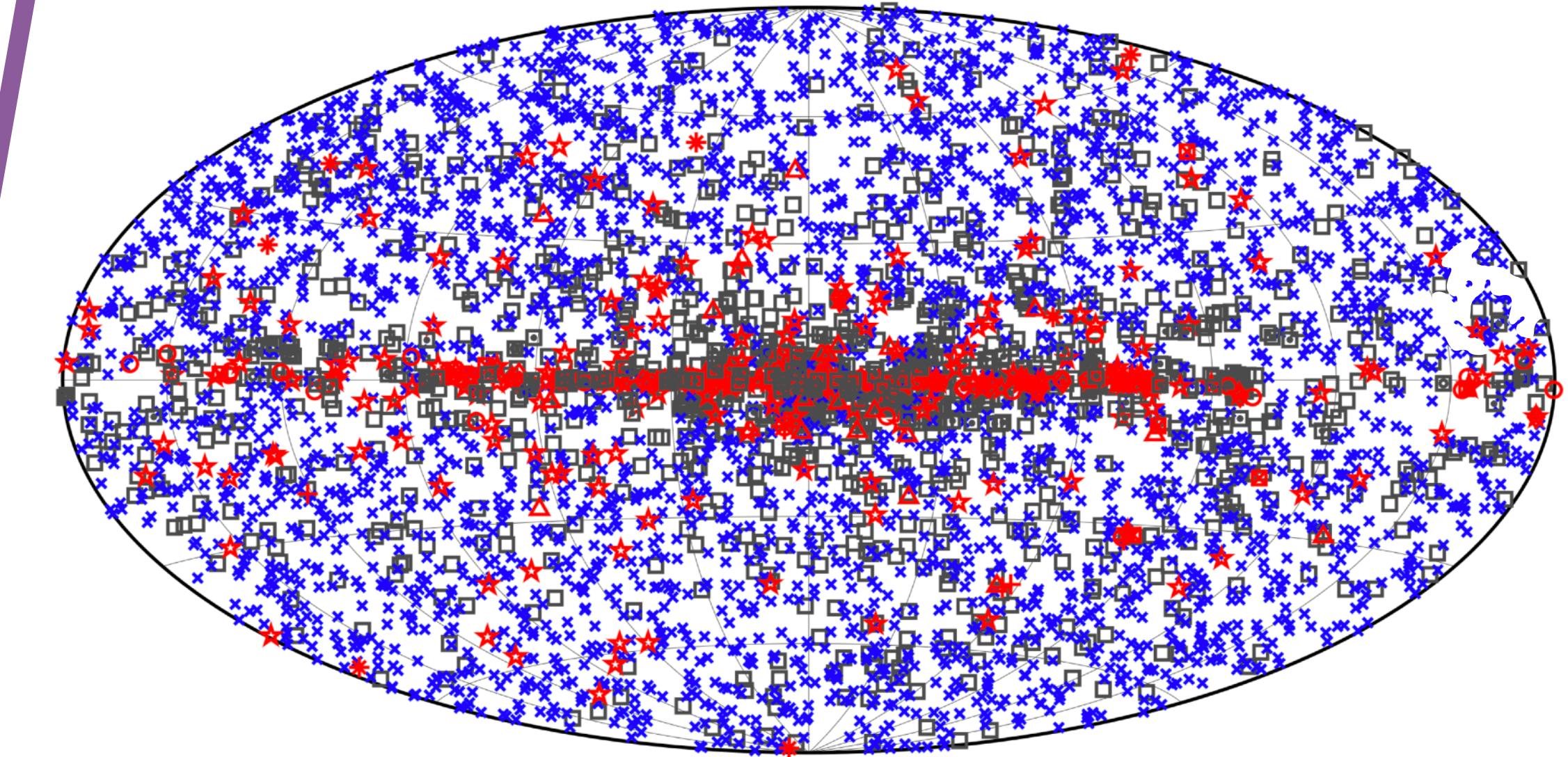
### Part 3:

Subhalos as point sources  
in the Milky Way

[2007.1039]



# Can dark matter subhalos be amongst the Fermi-LAT point sources?



[Fermi-LAT collaboration 19]

- **1525**

unassociated point sources  
in Fermi-LAT 4<sup>th</sup> catalog (4FGL)

[Fermi-LAT collaboration 19]

With our subhalo model + foreground/background model:  
Can some of these sources be DM halos?  
Could we detect them before the diffuse Galactic component?

# With our model we compute probabilities for the J-factors

- Probability to find a point-like subhalo with a J-factor above a threshold

$$\mathbb{P} ( > J, \psi, \delta\Omega ) = \frac{\delta\Omega}{N_{\text{sub}}} \iiint_{\textit{pt-like}} dm_t dc ds \left. \frac{\partial^2 n(m_t, c, s)}{\partial m_t \partial c} \right|_f \Theta(J_i(m_t, c, s) - J)$$

- Average number of visible subhalos:

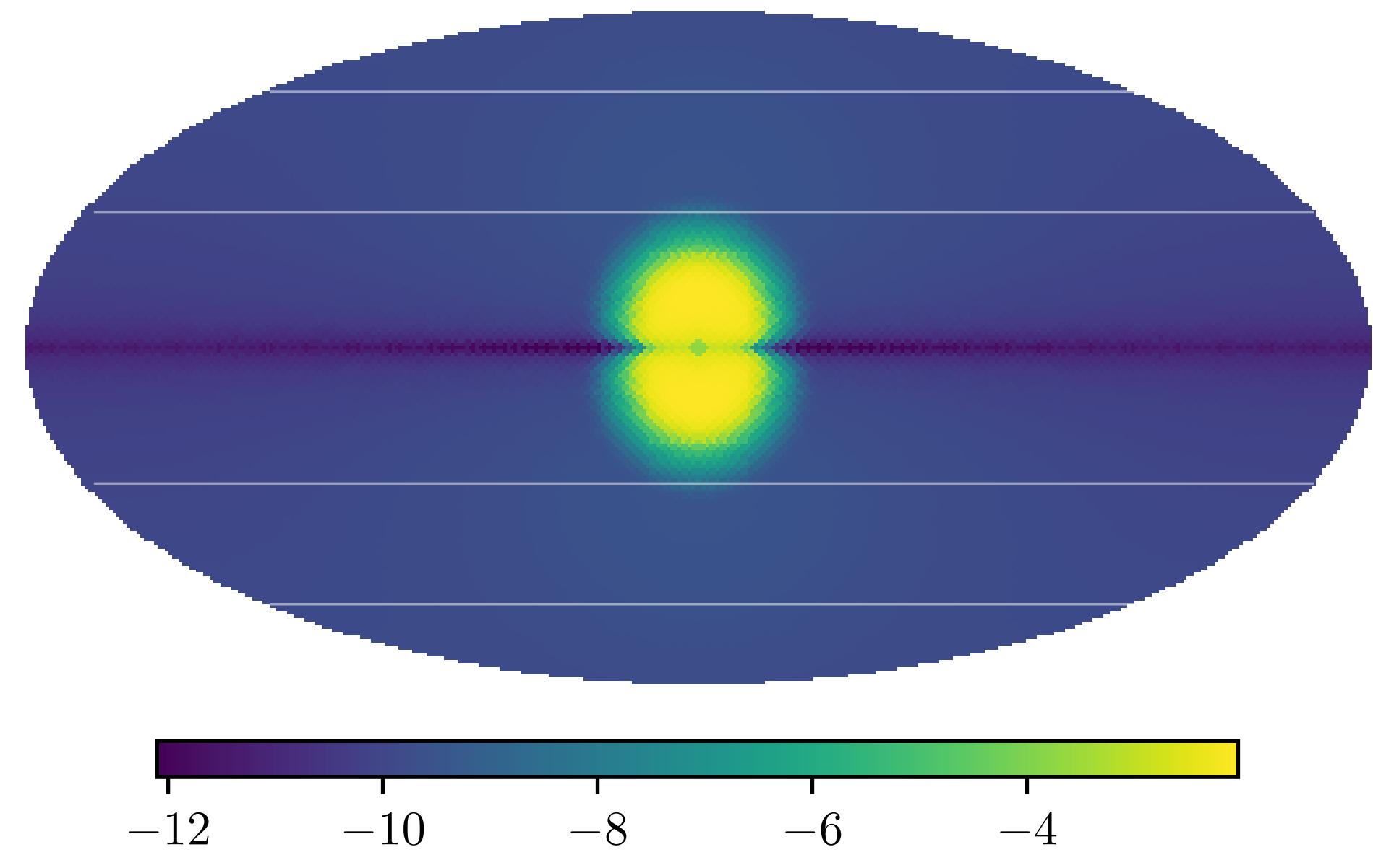
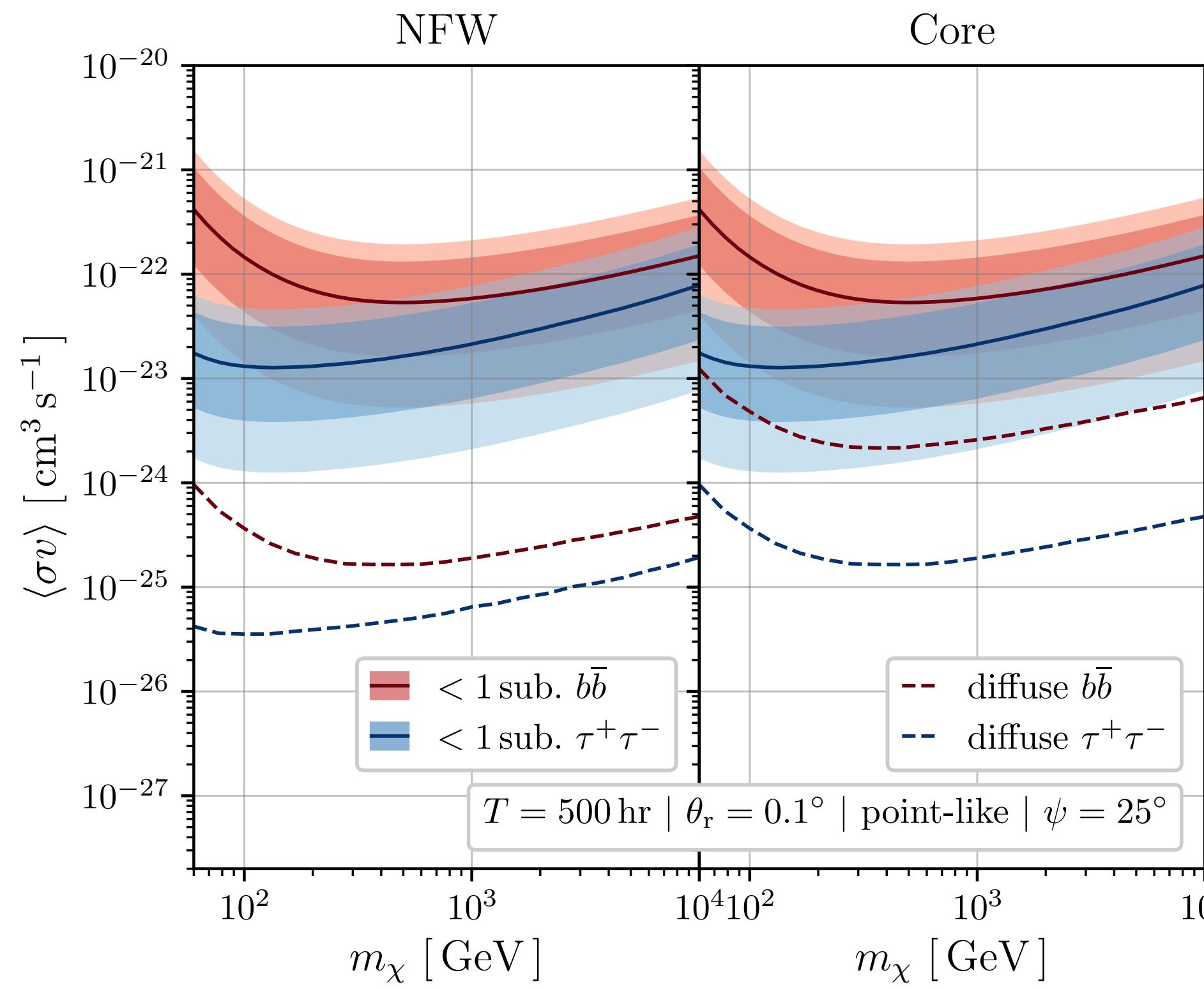
$$\langle N_{\text{vis}} \rangle = N_{\text{sub}} \mathbb{P}_J ( > J_{\min}, \psi, \delta\Omega )$$

# We add a background and perform a likelihood analysis

- Background model **compatible** with the baryonic distribution contributing to tidal stripping of the subhalos
- Likelihood analysis** and **mock data** to find the sensitivity to the diffuse halo and to subhalos (for Fermi-LAT and CTA)

$$\log_{10}(N_{\text{vis}}(l, b)/(\delta\Omega_r/\text{sr}))$$

Most « visible » sources  
are around  
the galactic center



For CTA and Fermi-LAT  
it is improbable to  
detect a subhalo before  
the diffuse emission  
(better chances if the MW halo is cored)

# Conclusions

- Subhalos modify/boost possible signals from dark matter annihilation
- Subhalos give a very significant boost with Sommerfeld enhancement (because smaller typical velocity)
- Subhalos could be searched as individual sources (even if hard to detect)