Latest Results from Daya Bay

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On behalf of the Daya Bay collaboration

Rencontres du Vietnam: Neutrino Physics, July 18, 2022

Outline

- Overview of the Daya Bay Experiment
- Details of recent results
  - Oscillation analysis with full dataset
  - Measurement of high-energy antineutrinos
- Highlights of other results (as time allows)

Overview of Daya Bay

Reactor antineutrino oscillation

$$\begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \cos\theta_{13} & 0 & \sin\theta_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{i\delta} & 0 & \cos\theta_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{-i\alpha_1/2} & 0 \\ 0 & 0 & e^{-i\alpha_2/2} \end{pmatrix}$$

Solar / Long baseline reactor	Short baseline reactor / Long baseline accelerator	Atmospheric / Long baseline accelerator	Neutrinoless double beta decay

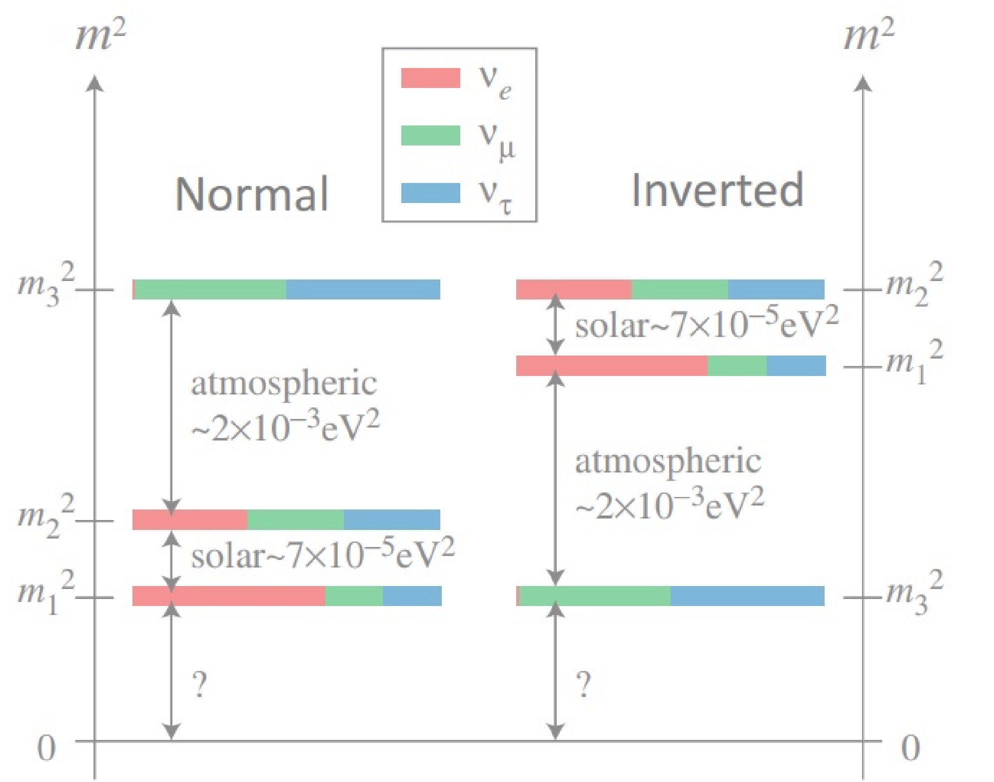
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{\text{PMNS}} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\begin{aligned} P_{\text{sur}} &= 1 - \cos^4\theta_{13}\sin^2 2\theta_{12}\sin^2\Delta_{21} \\ &\quad - \sin^2 2\theta_{13}(\cos^2\theta_{12}\sin^2\Delta_{31} + \sin^2\theta_{12}\sin^2\Delta_{32}) \\ &\equiv 1 - \cos^4\theta_{13}\sin^2 2\theta_{12}\sin^2\Delta_{21} \\ &\quad - \sin^2 2\theta_{13}\sin^2\Delta_{ee} \end{aligned}$$

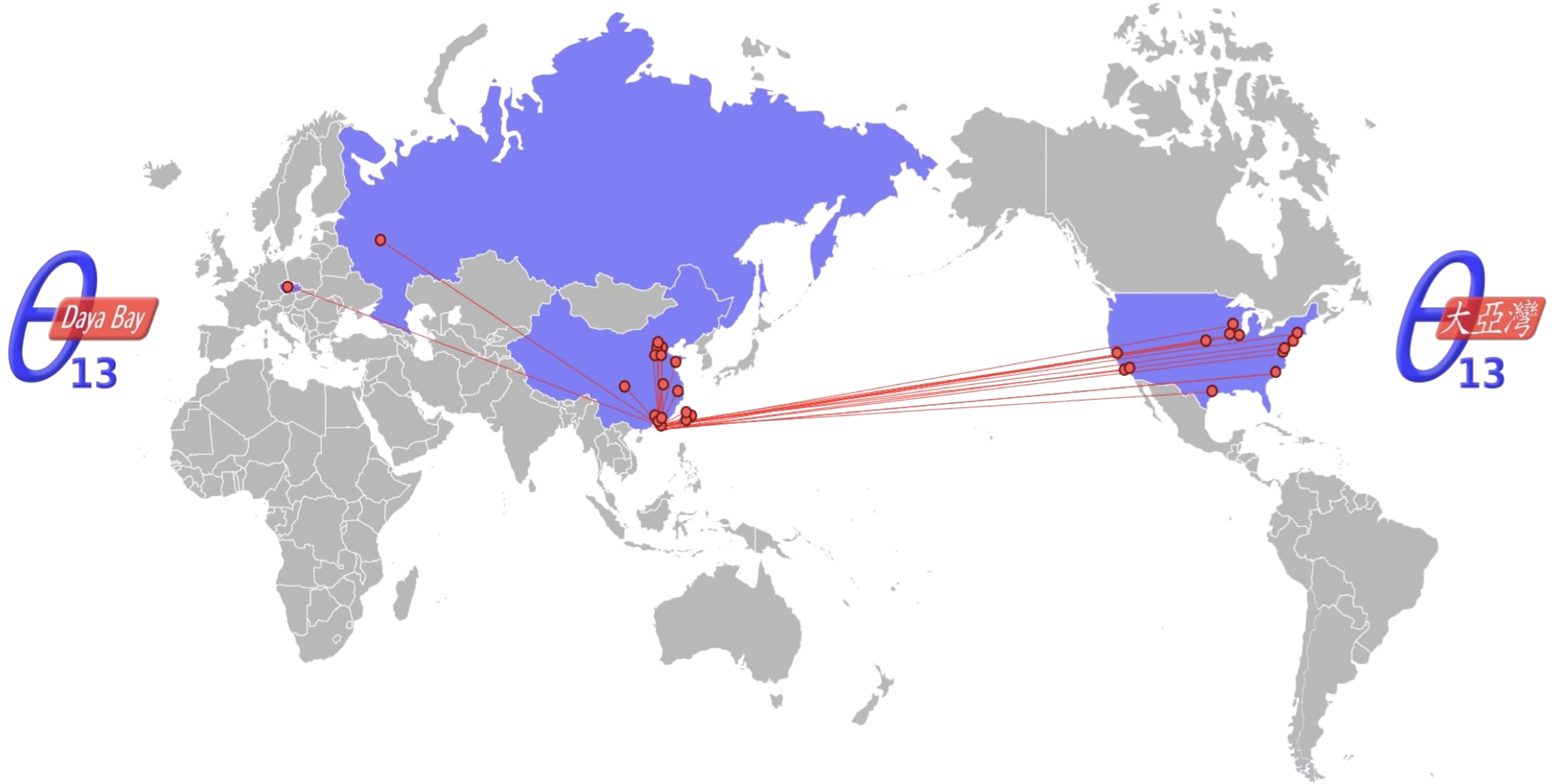
$$\Delta_{ji} \equiv \Delta m^2_{ji} L/4E$$

$$\Delta m^2_{ee} \simeq \cos^2\theta_{12}|\Delta m^2_{31}| + \sin^2\theta_{12}|\Delta m^2_{32}|$$

$$|\Delta m^2_{31}| \simeq \Delta m^2_{ee} \pm 2.3\times 10^{-5}\ \text{eV}^2$$
$$|\Delta m^2_{32}| \simeq \Delta m^2_{ee} \mp 5.2\times 10^{-5}\ \text{eV}^2$$



Daya Bay Collaboration



3 continents, ~200 collaborators, ~40 institutions

Daya Bay nuclear complex

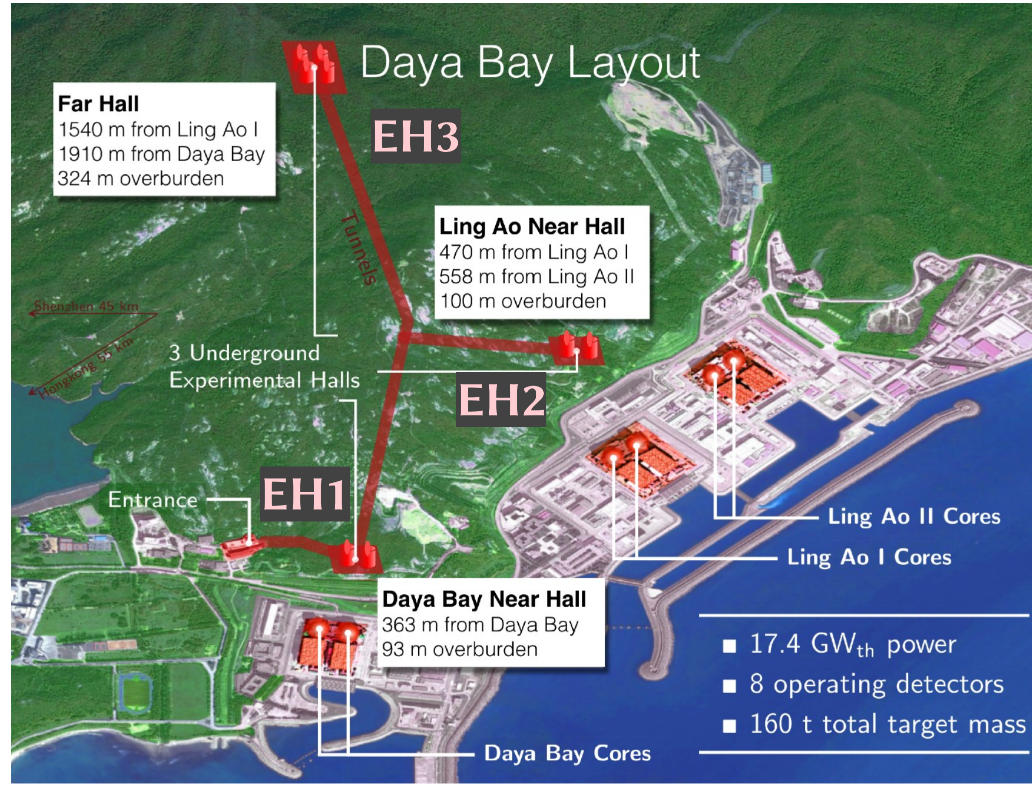
One of the world's most powerful nuclear power complexes and most intense antineutrino sources:

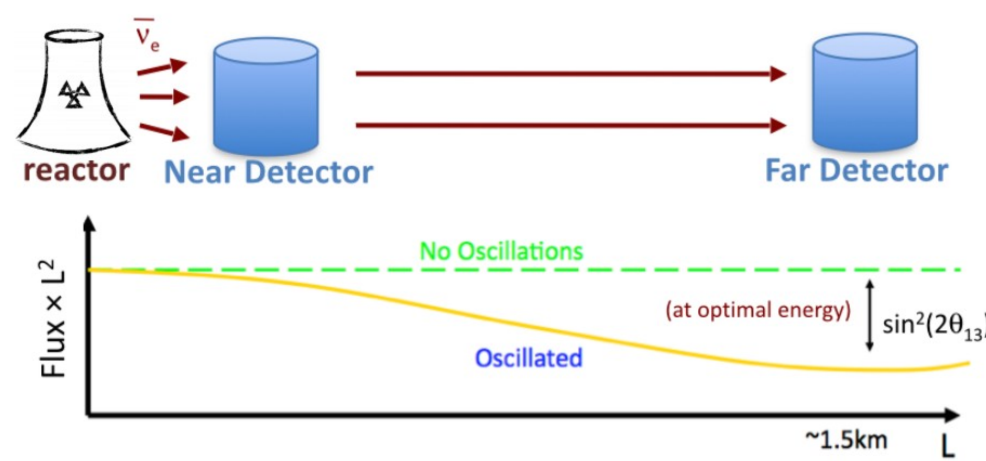
- Three plants (Daya Bay, Ling Ao I, Ling Ao II)
- Each plant: Two 2.9 GW_{th} pressurized water reactors
- Total power of 17.4 GW_{th} $\Longrightarrow$ **~4 × 10^{21} $\bar{\nu}_e$ / s**
- Power company provides time series data on power, burnup, fission fractions, etc.

Daya Bay experiment

Primary goal: Measure θ_{13} (and Δm^2_{ee}) via the short-baseline disappearance of reactor antineutrinos

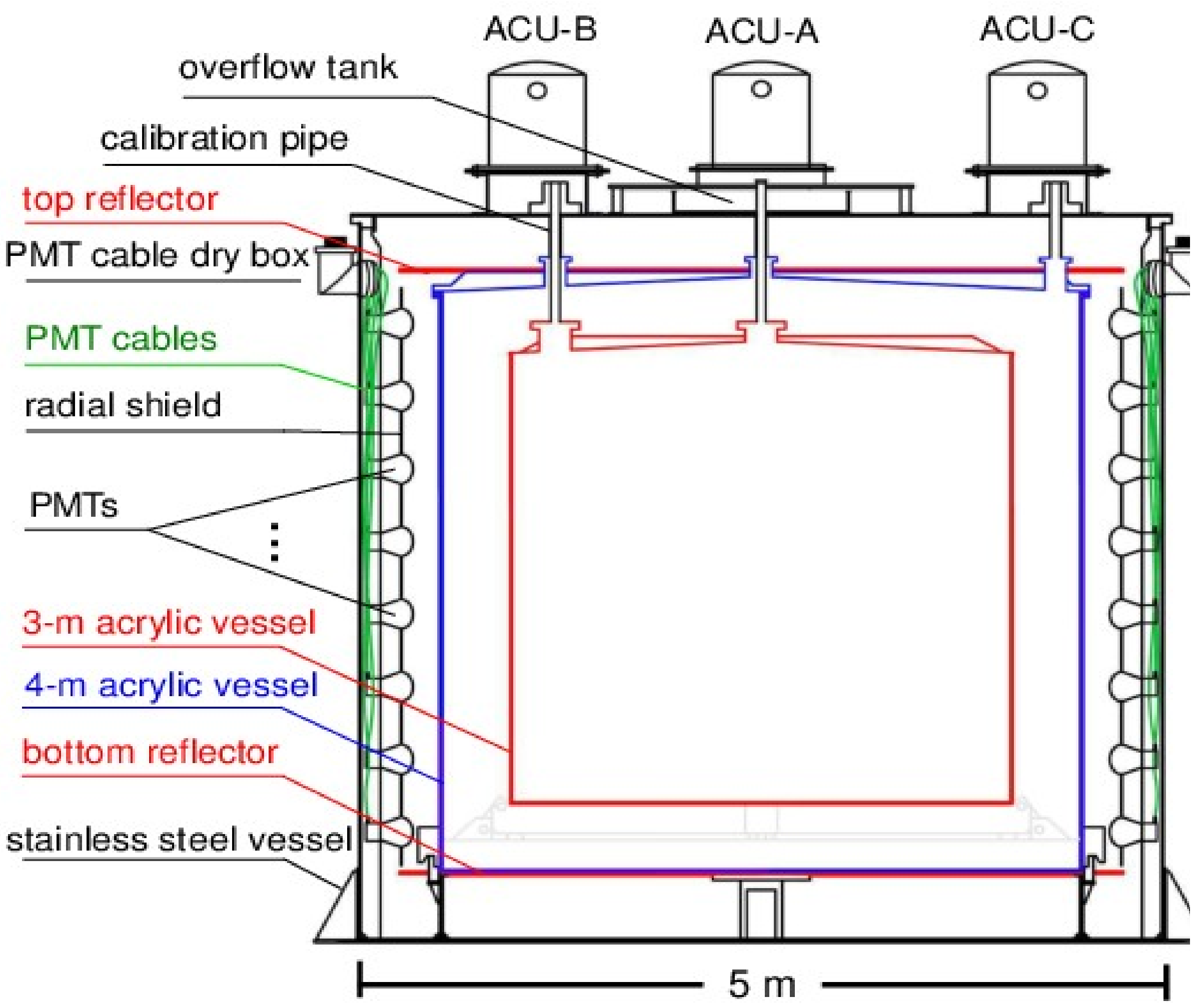
- **High statistics:** Powerful reactors, multiple large detectors
- **Low background:** Mountain overburden, radiopure materials, muon veto
- **Low systematics:** Near/far measurement essentially cancels (correlated) efficiency and reactor uncertainties
- **Optimal layout:** Far hall at disappearance maximum





*Nucl. Inst. Meth. A **811**, 133 (2016)*

Antineutrino detectors (ADs)



3 m inner acrylic vessel contains target mass: ~20 metric tons of 0.1% Gd-doped liquid scintillator

4 m outer acrylic vessel contains ~21 tons of undoped liquid scintillator to reduce inefficiency resulting from energy leakage from target

5 m stainless steel vessel contains ~40 tons of mineral oil for shielding, **192 8" PMTs** on walls for light detection

Three **Automated Calibration Units (ACUs)** on top for deployment of calibration sources

Nucl. Inst. Meth. A **773**, 8 (2015)

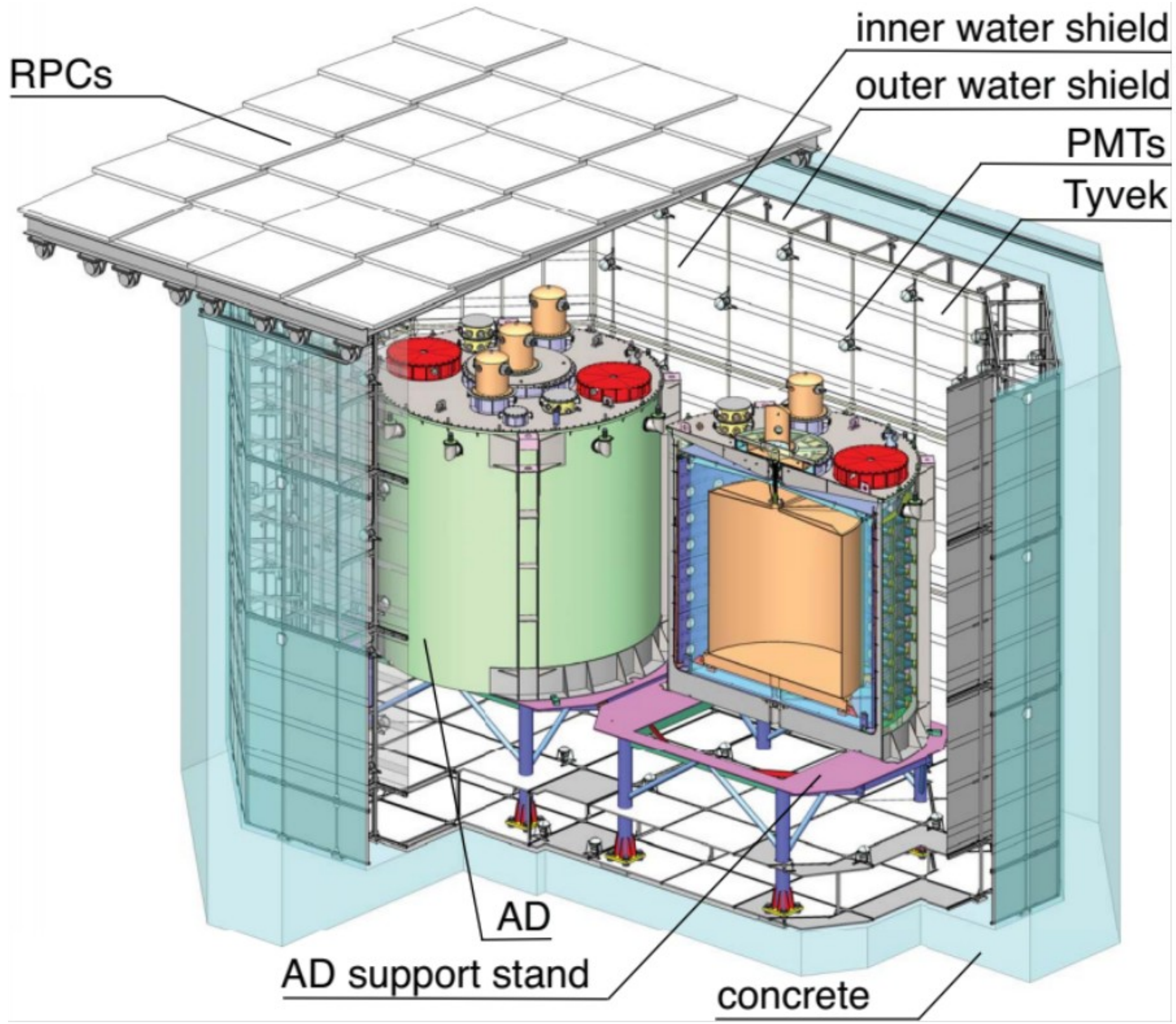
Muon system

Water pools (WPs) provide shielding against cosmic-ray muons, secondary neutrons, etc.

WPs are instrumented with PMTs, providing a muon veto system via detection of Cherenkov light

WPs are divided into inner and outer pools, optically and electronically isolated

Resistive plate chambers (RPCs) above WPs enable auxiliary studies of muons and associated backgrounds



Detecting antineutrinos

Antineutrinos are detected via **inverse β decay (IBD):**

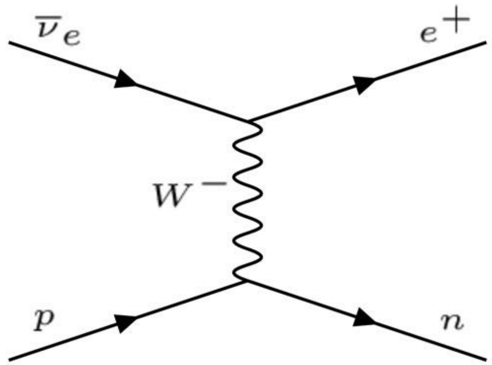
$$\bar{\nu}_e + p \to e^+ + n$$

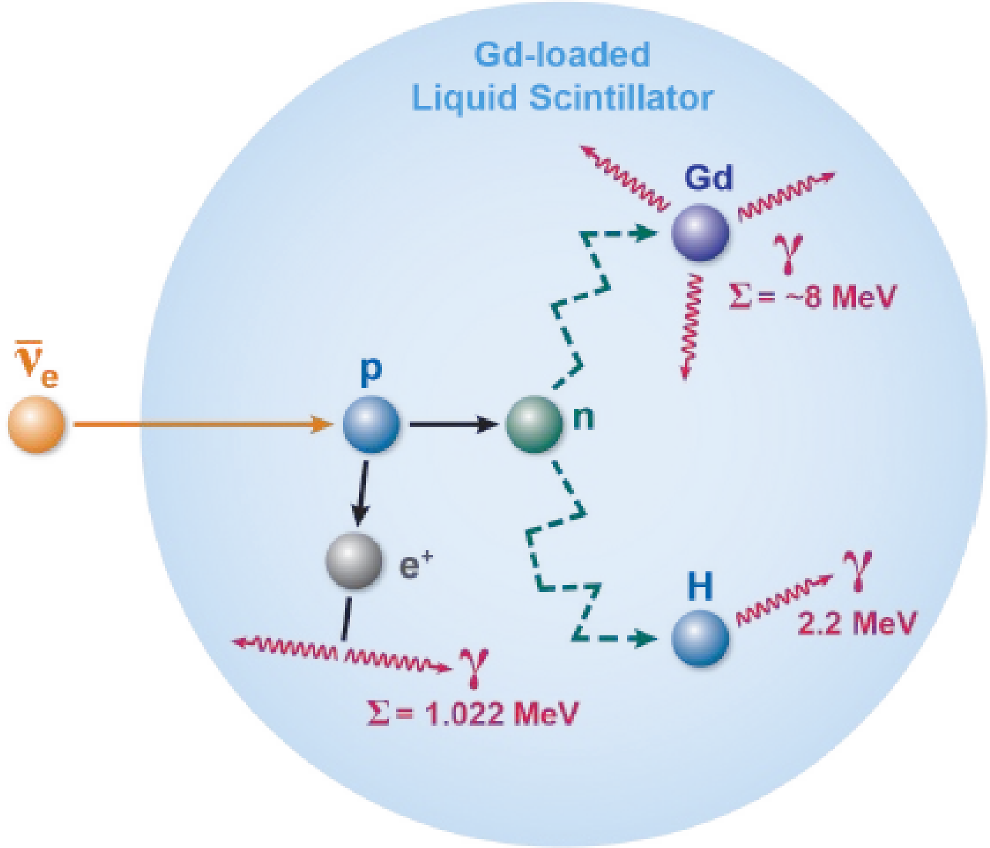
After an average of 28 μs, the neutron is captured on Gd, resulting in ~8 MeV of deexcitation gamma-rays. The coincident pulses provide a clean experimental signature, where

$$E_\nu \approx K_{e+} + 1.8\ \text{MeV}.$$

In terms of reconstructed (“prompt”) energy (including annihilation γ’s),

$$\mathbf{E_\nu \approx E_{prompt} + 0.8\ MeV}$$

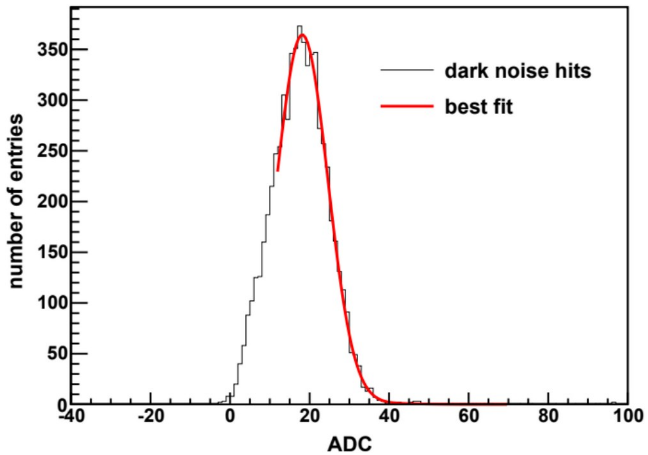




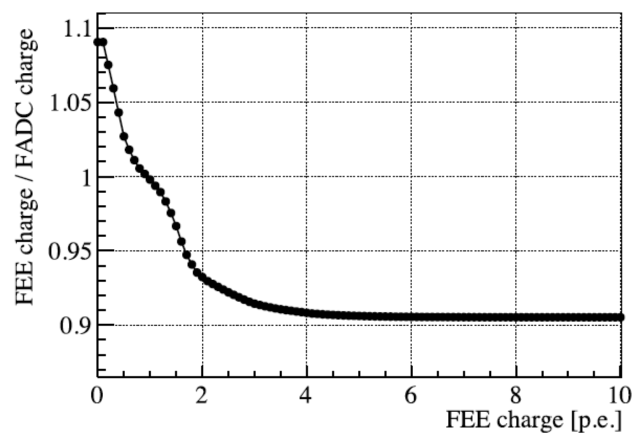
Phys. Rev. D **95**, 072006 (2016)

Reconstructed energy

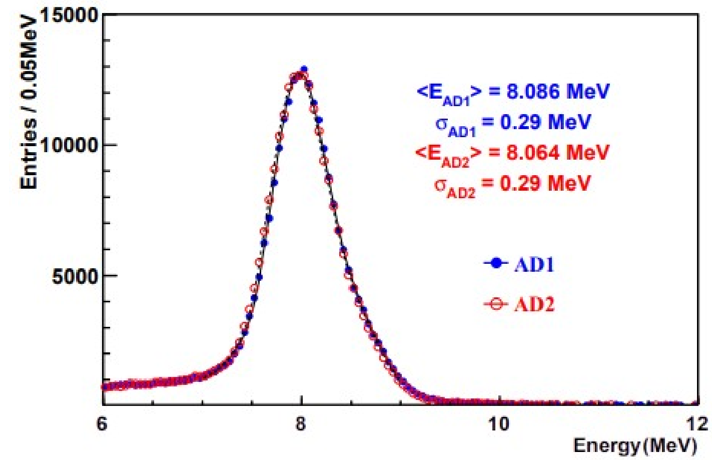
Gain calibration (ADC counts per photoelectron) from dark noise



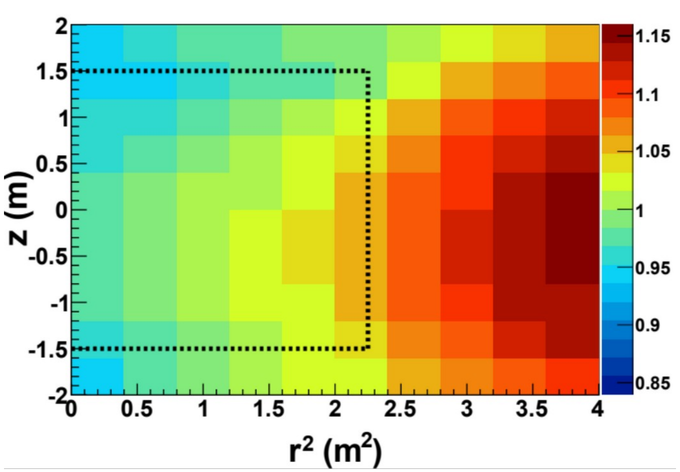
Channel-by-channel correction for **electronics nonlinearity** (NEW)



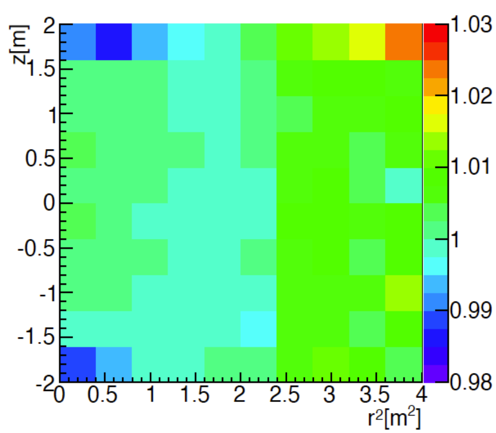
Energy scale calibration (photoelectrons per MeV) from spallation neutrons



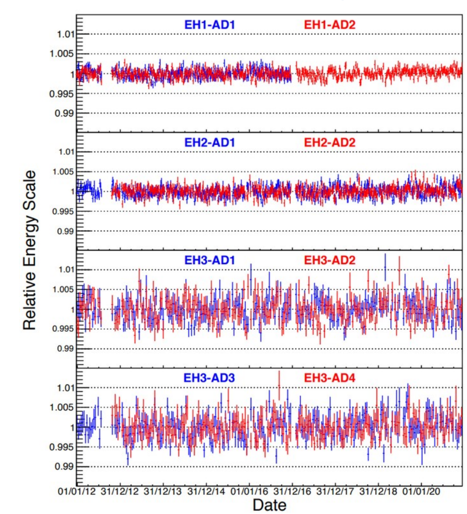
Nonuniformity correction from spallation neutrons



Additional correction based on α's (NEW)



End result: Stable and consistent reconstructed energy



Nucl. Inst. Meth. A ***940***, *230 (2019)*

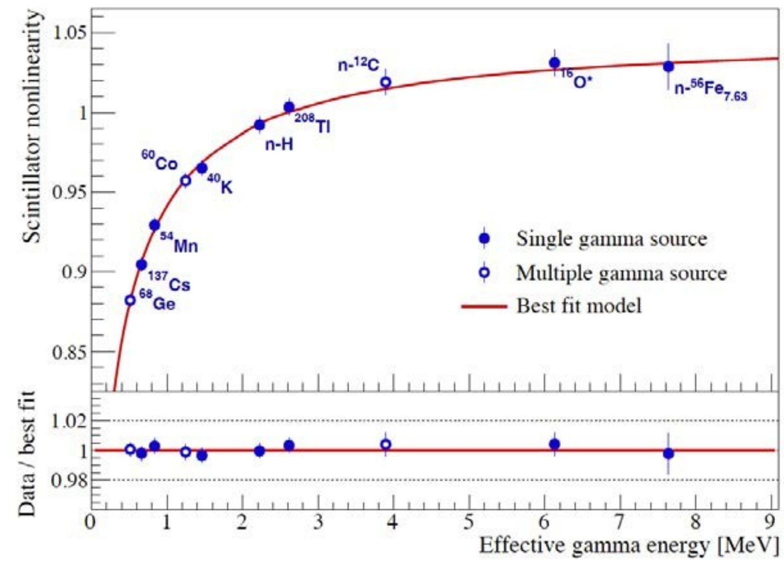
Antineutrino energy

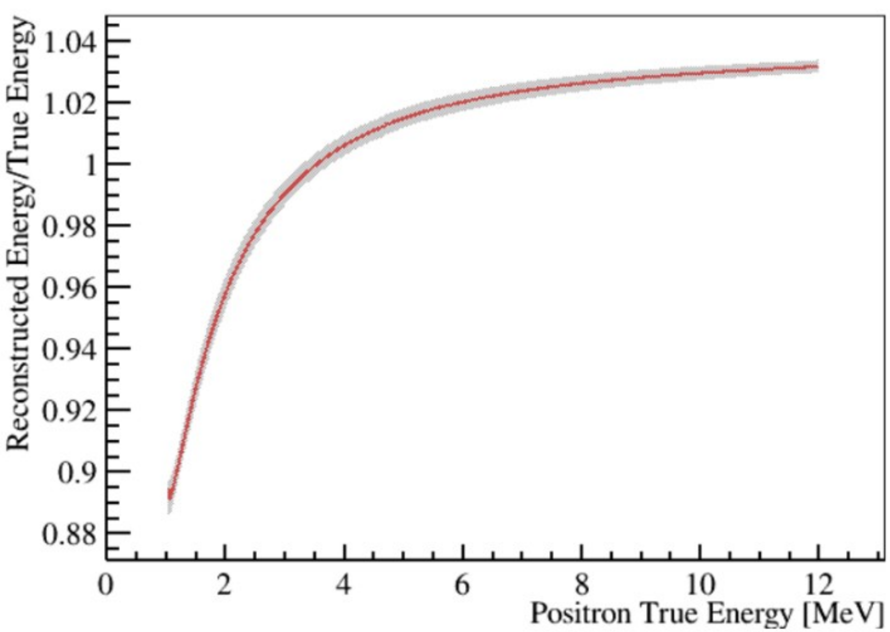
Take reconstructed energy, apply correction for scintillator nonlinearity (NL), get “true” deposited energy (e^+ ionization + annihilation)

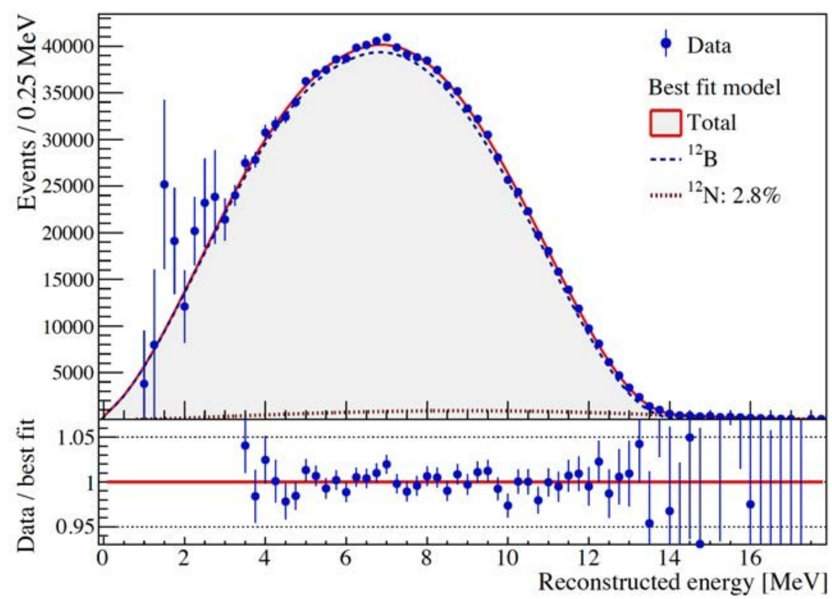
NL correction derived from calibration data, cross-checked with cosmogenic beta spectra

Antineutrino energy (for oscillation analysis) follows from kinematics:

$$E_\nu \approx E_{true} + 0.8\ \text{MeV}$$



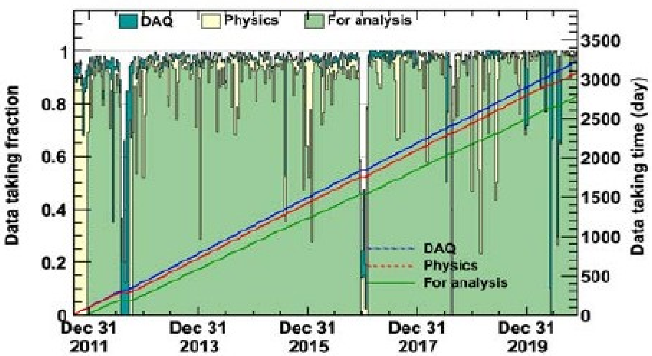




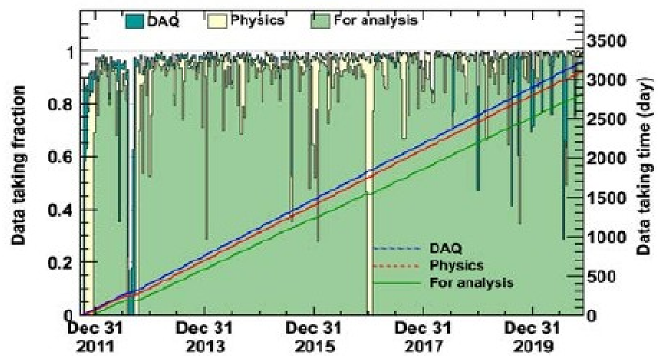
Oscillation analysis with full dataset

Full dataset!

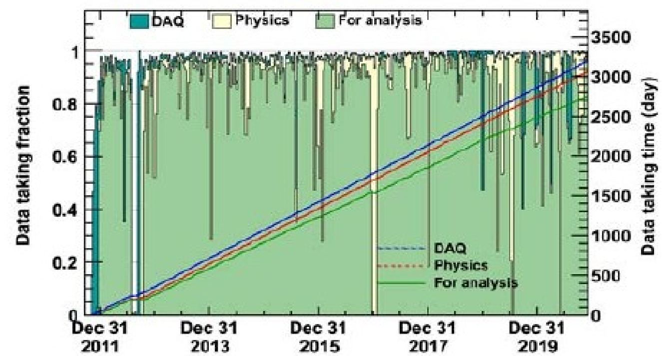
EH1



EH2



EH3



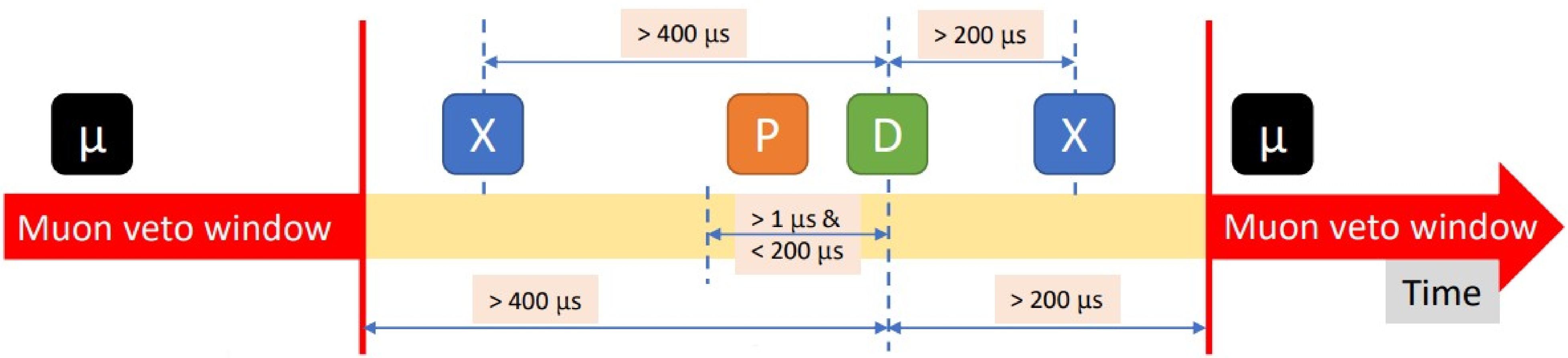
Configuration	EH1	EH2	EH3	Start date — end date	Duration (days)
6-AD	2	1	3	24 Dec 2011 — 28 July 2012	217
8-AD	2	2	4	19 Oct 2012 — 21 Dec 2016	1524
7-AD	1	2	4	26 Jan 2017 — 12 Dec 2020	1417
Total					**3158**

~2700 days of data in “good run list”; ~5.5 million antineutrinos

IBD selection

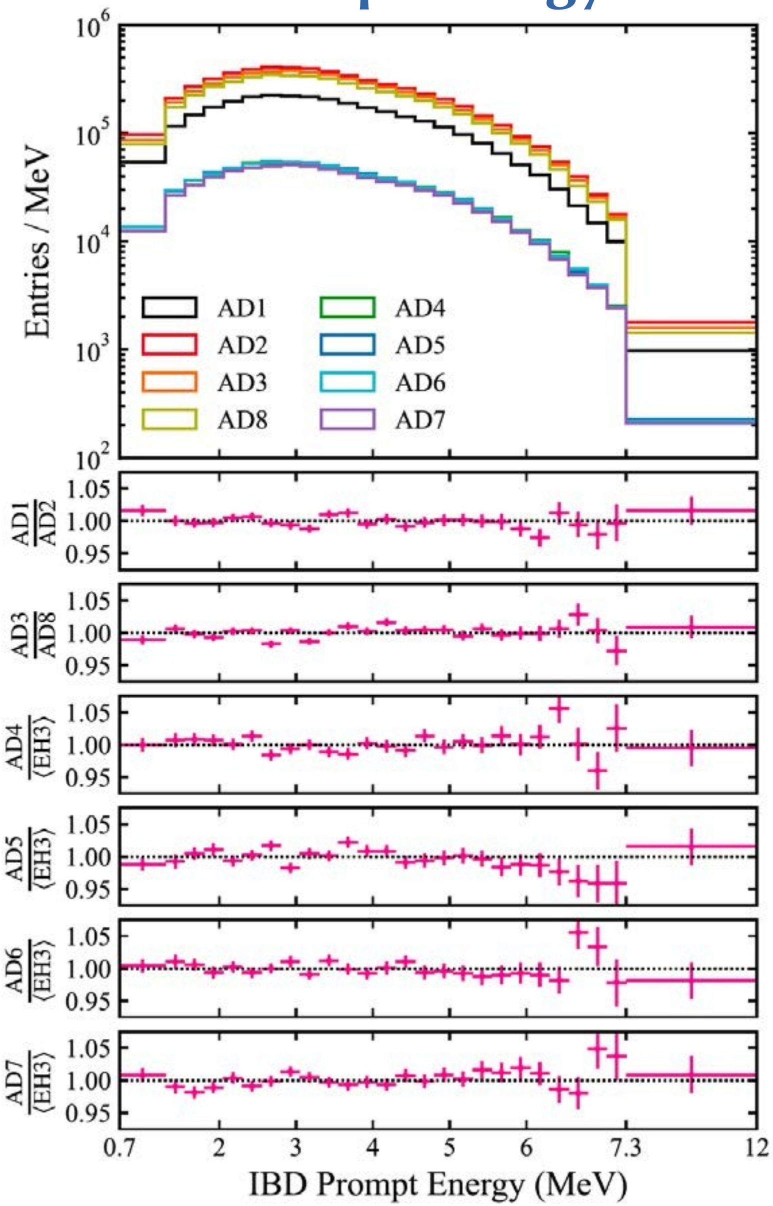
P = Prompt-like signal; D = Delayed-like signal; X = Non-muon signal

Muon veto	Water pool muon	Veto [-2 μs, 200 μs] after NHIT > **12** in OWS or IWS			
	AD muon	Veto [-2 μs, 1 ms] after > **20 MeV** signal in AD			
	AD shower	Veto [-2 μs, 0.4 s] after > **2 GeV** signal in AD			
	IWS muon veto	Veto [-2 μs, 10 μs] after **6** < NHIT ≤ **12** in IWS			(NEW)
Flasher cut		Standard, DocDB-7424		Residual, DocDB-12462	
Energy cut		0.7 ≤ P ≤ 12 MeV	6 ≤ D ≤ 12 MeV	0.7 ≤ X ≤ 20 MeV	
Decoupled Multiplicity Cut (DMC)	Full DMC for post-6AD period	Each D	One P within (-200 μs, -1 μs) && no X within (-400 μs, 200 μs)		
			Time to last muon veto window > 400 μs		
			Time to next muon veto window > 200 μs		

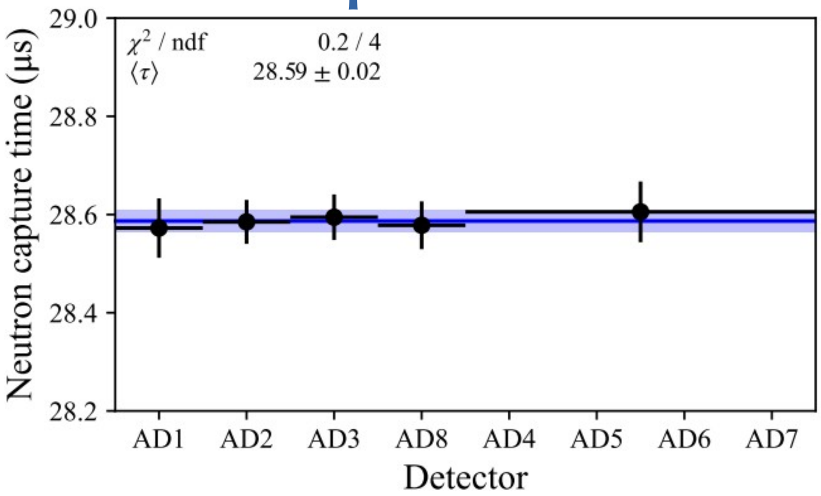


AD consistency

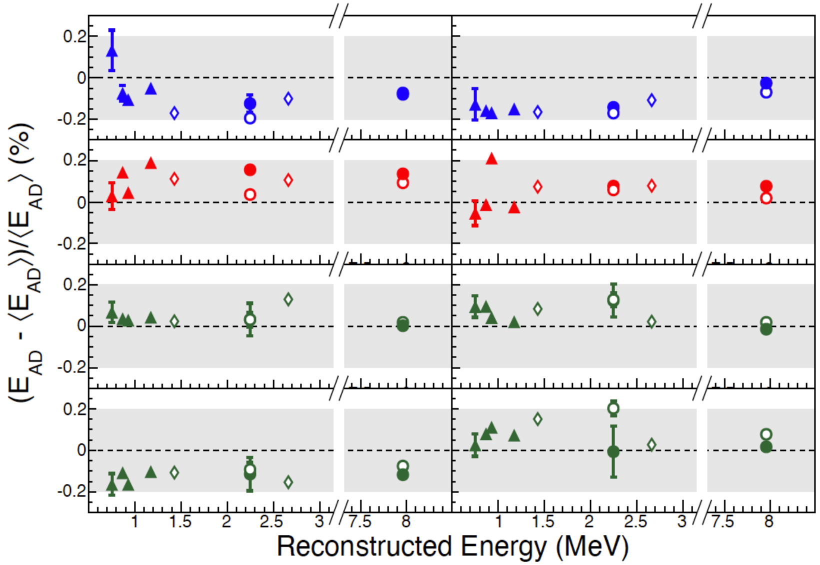
Prompt energy



Capture time



Energy scale



Phys. Rev. D ***95***, *072006 (2016)*

Summary of backgrounds

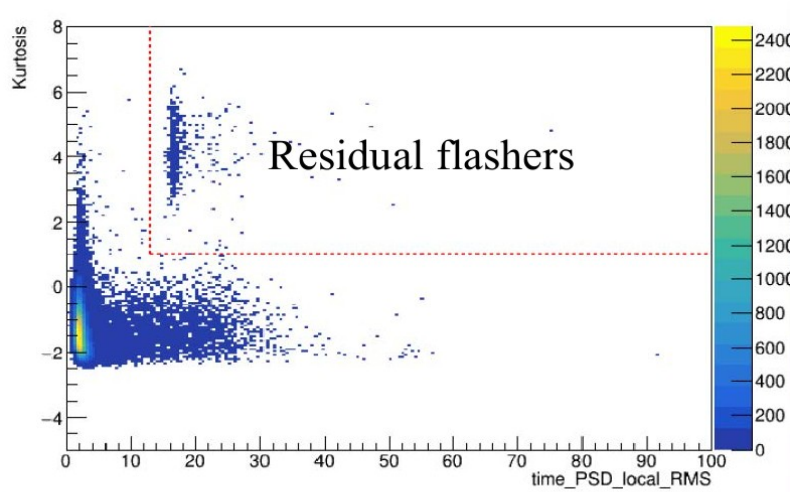
- **Flashers:** Light emission from PMTs
- **Accidentals** (uncorrelated pairs)
- **^{9}Li/^{8}He:** Cosmogenic isotopes ($\tau \sim 100$ ms). Can emit neutron after β decay.
- **Fast neutrons:** Ejected by muons from surroundings
- **"Muon-X" (NEW):** Prompt signal from low-energy muons. Delayed signals:
  - **Michel electron** from muon decay
  - **Neutron capture** from spallation
  - **Retrigger** from electronic ringing
- **^{241}Am^{13}C calibration source:** Pairs of gammas from neutron inelastic scattering followed by capture in detector materials
- **^{13}C$(\alpha,n)^{16}$O:** α from natural radioactivity + capture of ejected neutron

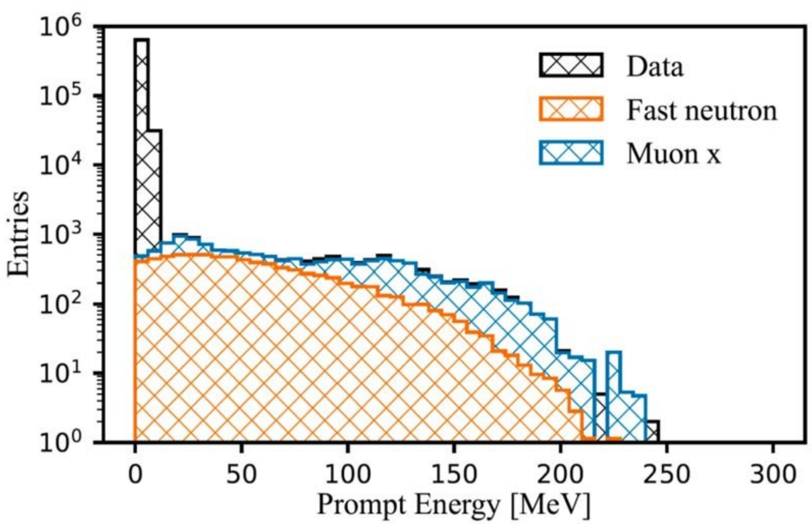
Details in backup slides!

Background updates

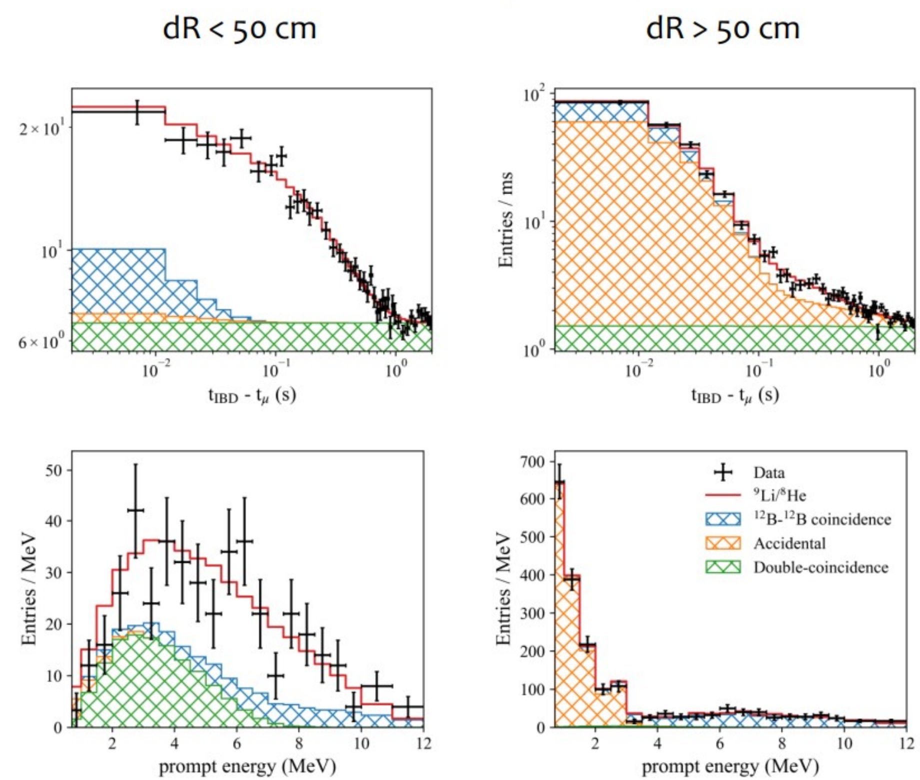
Residual flashers discovered in data added for this analysis, efficiently removed with new cut:

time_PSD_local_RMS > 13 && Kurtosis > 1





New **multidimensional fit** technique reduces uncertainty of **^{9}Li/^{8}He** background by ~2



New **muon-X** background due to failures of water pool PMTs. Largely eliminated by additional 10 μs veto on water pool muons with 7-12 PMT hits.

Efficiencies

Muon veto efficiency

$$T_\mu = T_{un-vetoed} - N_{un-vetoed} \cdot 600\mu s \qquad \epsilon_\mu = \frac{T_\mu}{T_{before}}$$

Total time outside veto windows — # of unvetoed periods > 600 μs — Total livetime ("before" vetoing)

Decoupled multiplicity cut (DMC) efficiency

$$\epsilon_{DMC} = P(0; R_x \cdot 400\mu s) \cdot P(0; R_x \cdot 200\mu s) = e^{-R_x \cdot 600\mu s}$$

Other efficiencies:

- Equal among ADs in Asimov case
- AD-uncorrelated uncertainties included in fit via pull terms.
- Correlated uncertainties do not affect the oscillation analysis

Description	Efficiency	Uncertainty	
		Correlated	Uncorrelated
Protons per AD		0.0092	0.0003
Gd capture ratio	0.8417	0.0095	0.001
Spill-in	0.0486	0.01	0.00019
Delayed energy cut	0.9271	0.0097	0.00072
Prompt energy cut	0.9981	0.001	0.0001
Capture time cut	0.987	0.0012	0.0002
Flasher cut	0.9998	0.0001	0.00013
Total	0.802484	0.01927	0.00132

IBD selection results

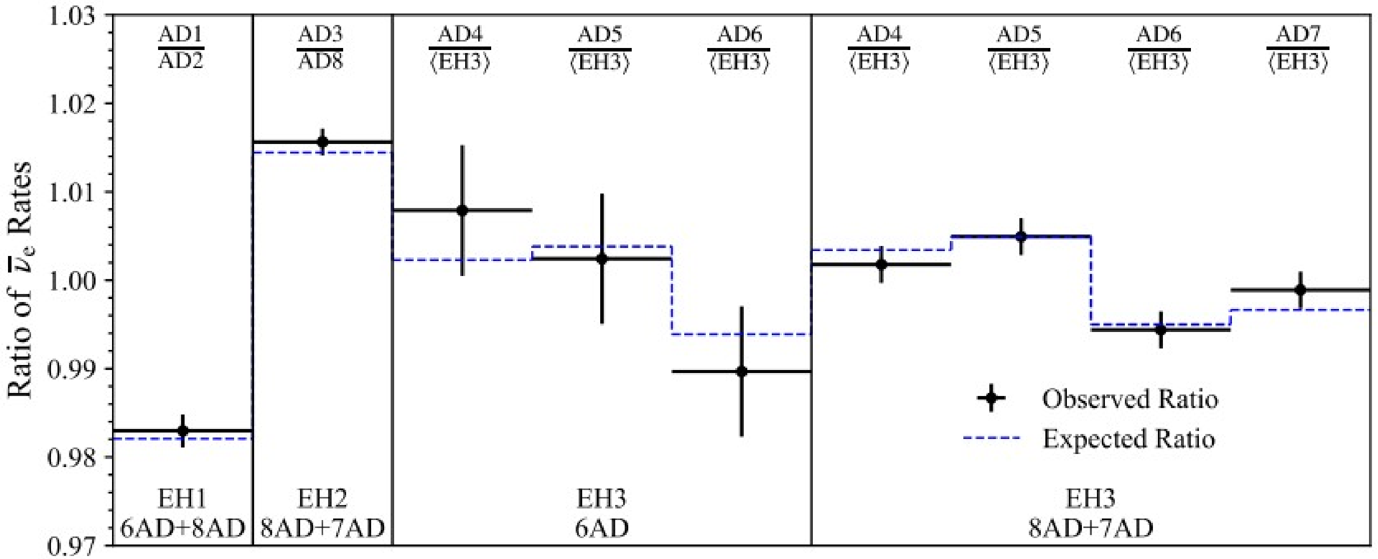
Summary table

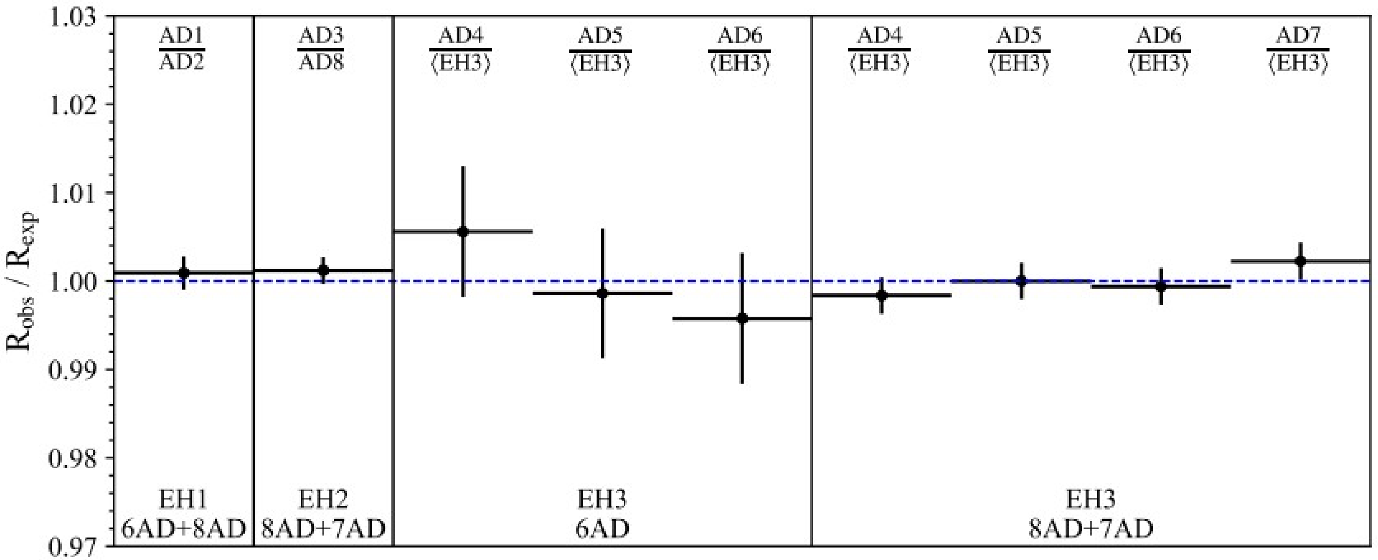
	EH1		EH2		EH3			
	AD1	AD2	AD3	AD8	AD4	AD5	AD6	AD7
$\bar{\nu}_e$ candidates	794335	1442475	1328301	1216593	194949	195369	193334	180762
DAQ live time [days]	1535.111	2686.11	2689.88	2502.816	2689.156	2689.156	2689.156	2501.531
ϵ_μ	0.8006	0.7973	0.8387	0.8366	0.9815	0.9815	0.9814	0.9814
$\bar{\epsilon}_m$	0.9671	0.9678	0.969	0.9688	0.9693	0.9693	0.9692	0.9693
Accidentals [day^{-1}]	7.11 ± 0.01	6.76 ± 0.01	5.00 ± 0.00	4.85 ± 0.01	0.80 ± 0.00	0.77 ± 0.00	0.79 ± 0.00	0.66 ± 0.00
Fast neutron & muon-x [day^{-1}]	0.83 ± 0.17	0.96 ± 0.19	0.56 ± 0.11	0.56 ± 0.11	0.05 ± 0.01	0.05 ± 0.01	0.05 ± 0.01	0.05 ± 0.01
^{9}Li, ^{8}He [AD^{-1} day^{-1}]	2.97 ± 0.53		2.09 ± 0.36		0.25 ± 0.03			
^{241}Am-^{13}C [day^{-1}]	0.16 ± 0.07	0.13 ± 0.06	0.12 ± 0.05	0.11 ± 0.05	0.04 ± 0.02	0.04 ± 0.02	0.04 ± 0.02	0.03 ± 0.01
^{13}C(α, n)^{16}O [day^{-1}]	0.08 ± 0.04	0.06 ± 0.03	0.04 ± 0.02	0.06 ± 0.03	0.04 ± 0.02	0.04 ± 0.02	0.03 ± 0.02	0.04 ± 0.02
$\bar{\nu}_e$ rate, $R_{\bar{\nu}_e}$ [day^{-1}]	657.11 ± 0.94	685.09 ± 0.81	599.83 ± 0.65	592.07 ± 0.67	75.03 ± 0.18	75.22 ± 0.18	74.42 ± 0.18	74.94 ± 0.18

Backgrounds at sub-2% level

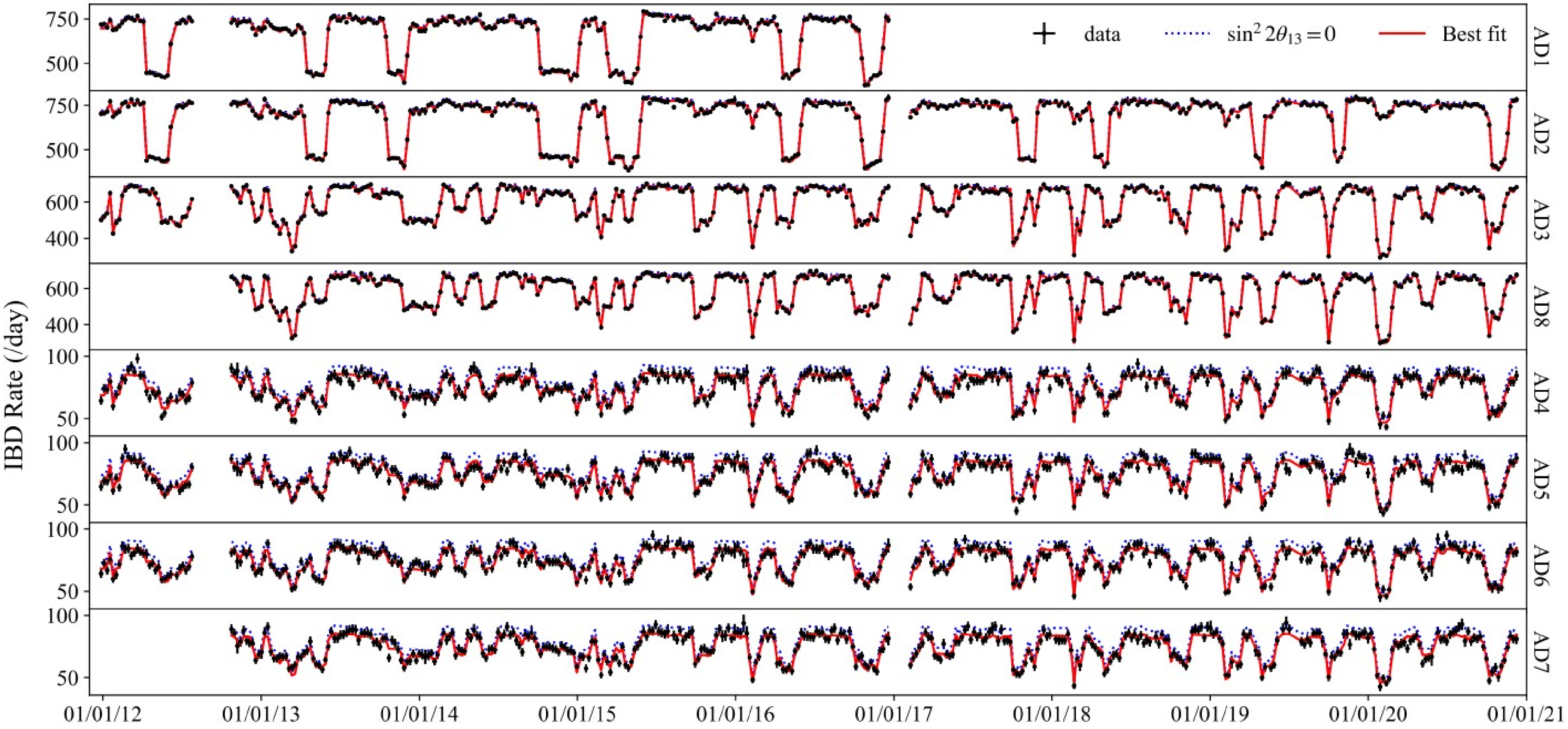
Side-by-side IBD rates

(Background-subtracted, efficiency-corrected)





IBD rates over time



Oscillation fit

Fitter's prediction

- Use Huber-Mueller isotope spectra
  - With non-equlibirium correction
  - Plus spent nuclear fuel contribution
  - Shape/normalization allowed to vary (pull terms)
- Fold with reactor power and fission fractions (from power company) to predict antineutrino flux
- Propagate to ADs ($1/R^2$, oscillation)
- Convert to positron energy (kinematics)
- Apply effects of energy leakage, nonlinearity, resolution (w/ pulls)
- Add predicted backgrounds (w/ pulls)

Prediction for fitting

$$F_d^{\ i} = T_d^{\ i}\Bigg(1+\varepsilon_D+\varepsilon_d+\frac{0.072\varepsilon_d^E}{0.2}+b^i\varepsilon_d^E+\varepsilon_i+\sum_{l=1}^{4} f_l^{\ i}\varepsilon_l+g^i\varepsilon_d^{IAV}$$

$$+a^i\varepsilon_R+\begin{cases}\sum_{r=1}^{6}\omega_r^d(\varepsilon_r+\varepsilon_r^n S_n^i+\varepsilon_r^s S_s^i+\varepsilon_r^f S_f^i) & \text{(6AD period)}\\ \sum_{r=1}^{6}\omega_r^d(\varepsilon_r'+\varepsilon_r^n S_n^i+\varepsilon_r^s S_s^i+\varepsilon_r^f S_f^i) & \text{(8AD period)}\\ \sum_{r=1}^{6}\omega_r^d(\varepsilon_r''+\varepsilon_r^n S_n^i+\varepsilon_r^s S_s^i+\varepsilon_r^f S_f^i) & \text{(7AD period)}\end{cases}\Bigg)$$

$$+B_{di}^{Li9}(1+\eta_k^{Li9})(1+\eta_i^{Li9Shape})+B_{di}^{Fn}(1+\eta_k^{Fn})(1+\eta_i^{FnShape})$$

$$+B_{di}^{AmC}(1+\eta^{AmC})(1+\eta_i^{AmCShape})+B_{di}^{acc}(1+\eta_d^{acc})+B_{di}^{alphaN}(1+\eta_d^{alphaN})$$

$$+B_{di}^{Md}(1+\eta_k^{Md})\ \text{(7AD period)}$$

Indices:

i = Energy bin
d = Detector
r = Reactor
k = Experimental hall

F_d^i = Best-fit prediction
T_d^i = Nominal prediction w/ oscillation
ε_D (ε_d) = Correlated (uncorr.) det. eff. pulls
ε_d^E = Relative energy scale pulls
ε_i = Bin-uncorrelated shape pulls

ε_l = Nonlinearity pulls
ε_d^{IAV} = Energy leakage pulls
ε_R = Correlated reactor flux pull
ε_r (ε_r^n, ε_r^s, ε_r^f,) = Uncorr. reactor flux pulls
(non-equilibrium, spent fuel, fission frac.)
η = Background rate/shape pulls

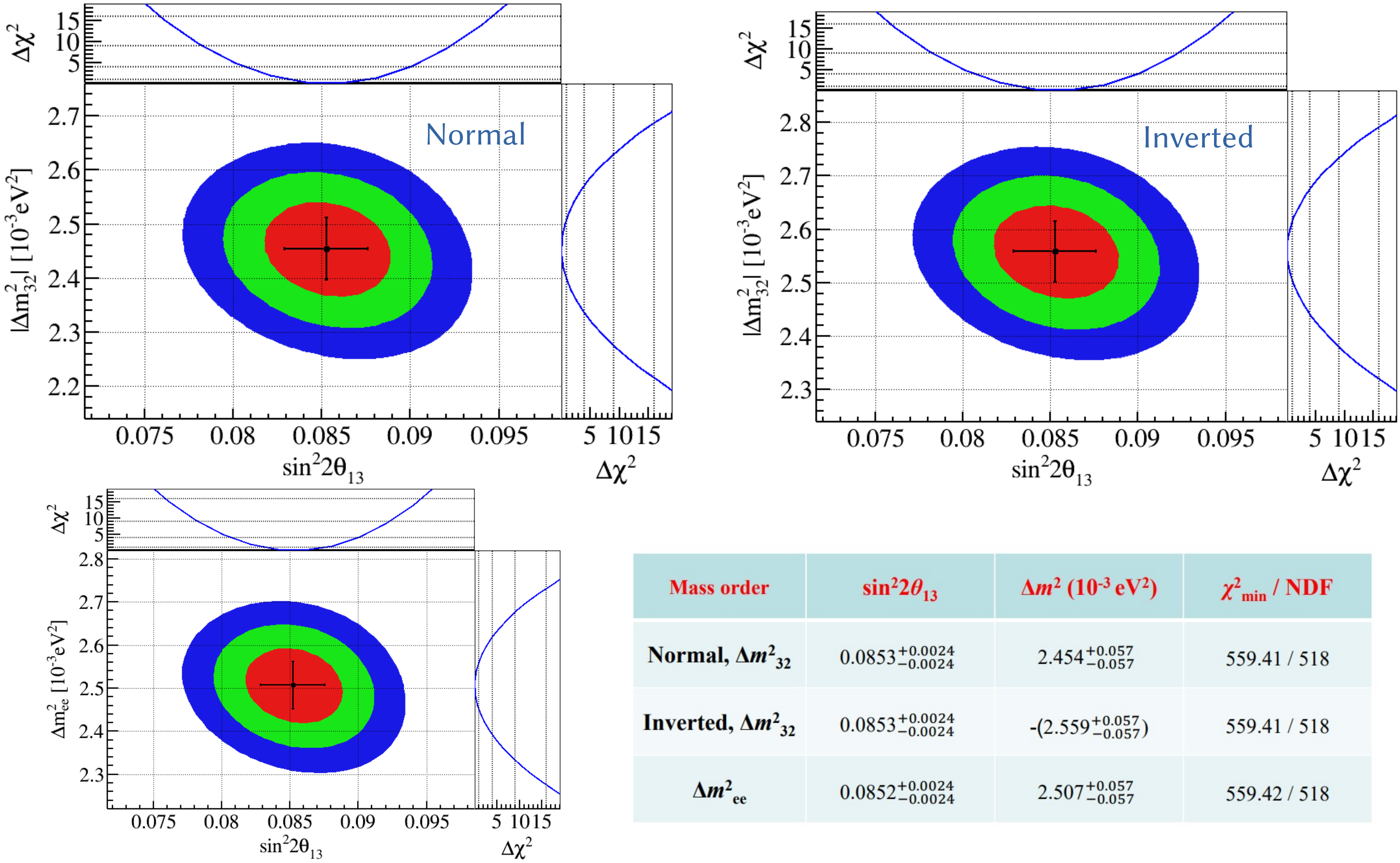
χ^2 expression

M_d^i = Measurement

$$\chi^2 = \min_{\gamma} \sum_{d=1}^{6}\sum_{i=1}^{29} \frac{[M_d^i - F_d^i]^2}{F_d^i} + \sum_{d=1}^{8}\sum_{i=1}^{29} \frac{[M_d^i - F_d^i]^2}{F_d^i} + \sum_{d=1}^{7}\sum_{i=1}^{29} \frac{[M_d^i - F_d^i]^2}{F_d^i}$$
$$+ \sum_{d=1}^{8}\left[(\frac{\varepsilon_d}{\sigma_d})^2 + (\frac{\eta_d^{acc}}{\sigma_d^{acc}})^2 + (\frac{\eta_d^{alphaN}}{\sigma_d^{alphaN}})^2 + (\frac{\varepsilon_d^E}{\sigma_d^E})^2 + (\frac{\varepsilon_d^{IAV}}{\sigma_d^{IAV}})^2\right]$$
$$+ \sum_{k=1}^{3}\left[(\frac{\eta_k^{Fn}}{\sigma_k^{Fn}})^2 + (\frac{\eta_k^{Li9}}{\sigma_k^{Li9}})^2 + (\frac{\eta_k^{Md}}{\sigma_k^{Md}})^2\right] + (\frac{\eta^{AmC}}{\sigma^{AmC}})^2$$
$$+ \sum_{i=1}^{29}(\frac{\eta_i^{Li9Shape}}{\sigma_i^{Li9Shape}})^2 + \sum_{i=1}^{29}(\frac{\eta_i^{FnShape}}{\sigma_i^{FnShape}})^2 + \sum_{i=1}^{14}(\frac{\eta_i^{AmCShape}}{\sigma_i^{AmCShape}})^2 + \sum_{l=1}^{4}(\frac{\varepsilon_l}{\sigma_l})^2$$
$$+ \sum_{r=1}^{6}\left[(\frac{\varepsilon_r}{\sigma_r})^2 + (\frac{\varepsilon_r^n}{\sigma_r^n})^2 + (\frac{\varepsilon_r^s}{\sigma_r^s})^2 + (\frac{\varepsilon_r^f}{\sigma_r^f})^2 + (\frac{\varepsilon_r'}{\sigma_r'})^2 + (\frac{\varepsilon_r''}{\sigma_r''})^2\right]$$

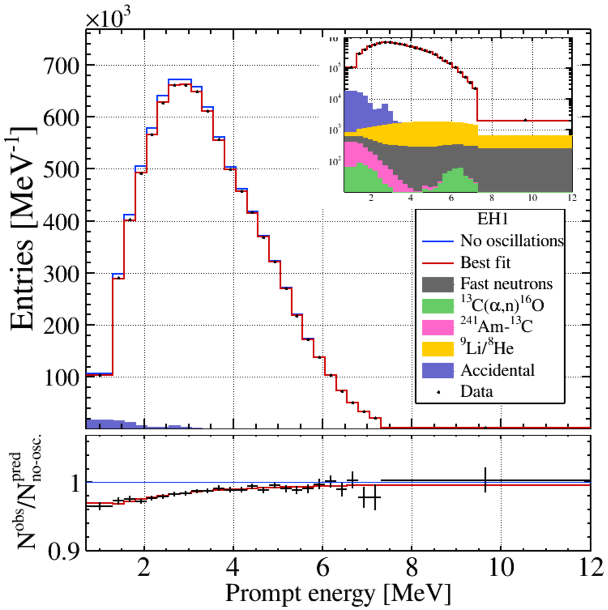
Minimize this χ^2 with respect to oscillation parameters and all pulls
Result: Best-fit $\sin^2 2\theta_{13}$ and Δm^2!
Individual 1σ bounds where $\Delta\chi^2$ crosses 1 (w/ other osc. par. fixed)
1σ (2σ, 3σ) 2D contours where $\Delta\chi^2$ crosses 2.30 (6.18, 11.83)

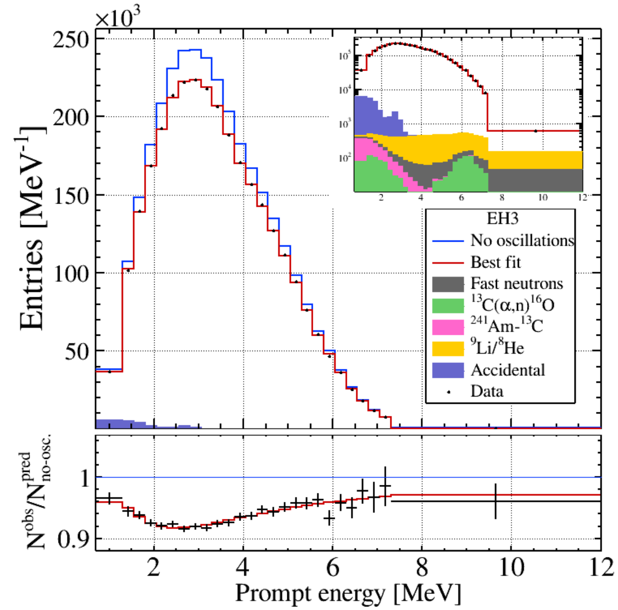
Oscillation parameters

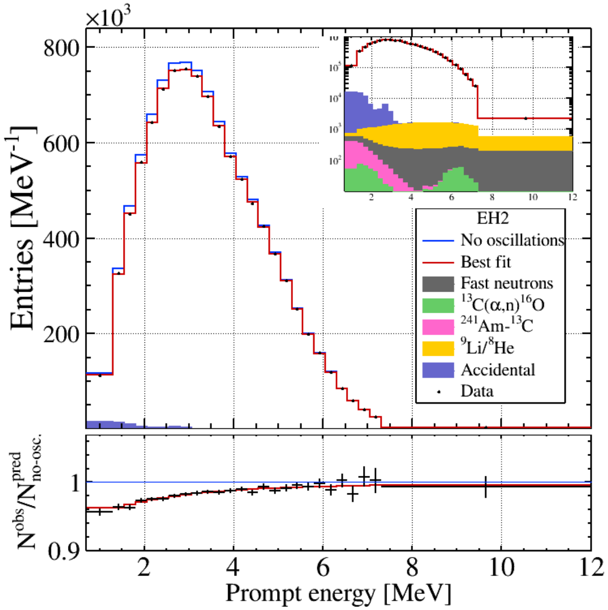


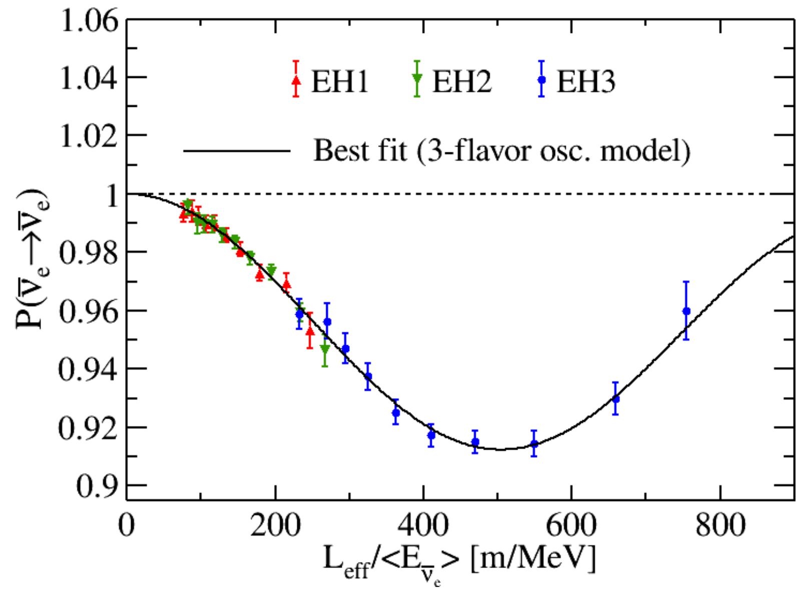
Mass order	$\sin^2 2\theta_{13}$	Δm^2 (10^{-3} eV^2)	χ^2_{min} / NDF
Normal, Δm^2_{32}	$0.0853^{+0.0024}_{-0.0024}$	$2.454^{+0.057}_{-0.057}$	559.41 / 518
Inverted, Δm^2_{32}	$0.0853^{+0.0024}_{-0.0024}$	$-(2.559^{+0.057}_{-0.057})$	559.41 / 518
Δm^2_{ee}	$0.0852^{+0.0024}_{-0.0024}$	$2.507^{+0.057}_{-0.057}$	559.42 / 518

Prompt spectra



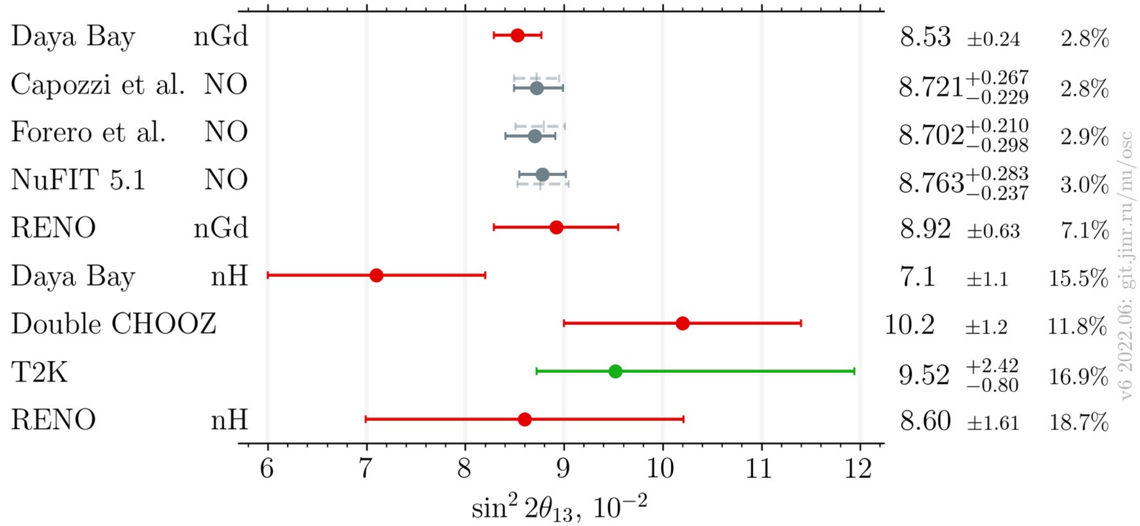




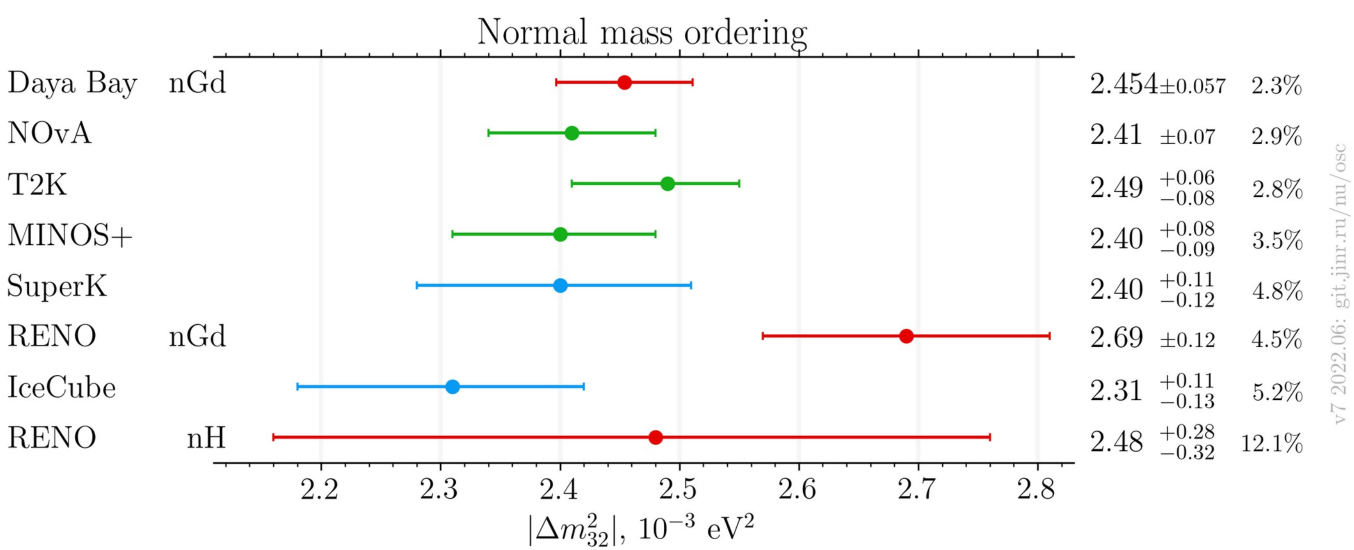


Global comparison

$\sin^2 2\theta_{13}$



Δm^2_{32}



Measurement of high-energy reactor antineutrinos

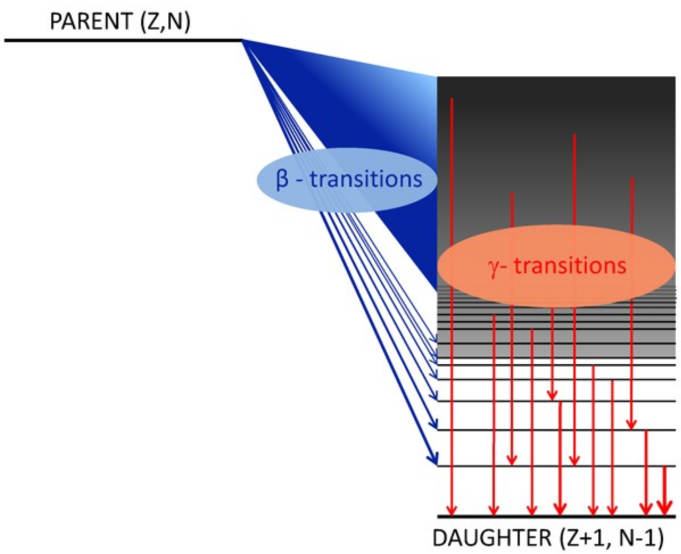
High-energy reactor $\bar{\nu}$

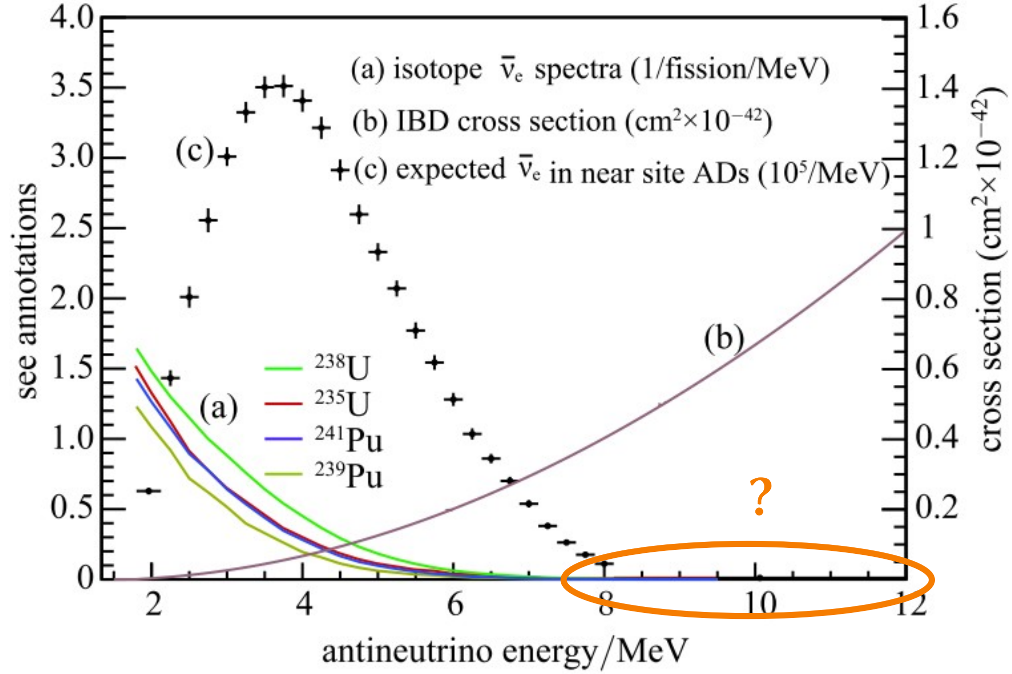
Why study high-energy reactor antineutrinos **(HERA)** having **prompt energy > 8 MeV**?

- Shed light on discrepancies between data and nuclear models
- Characterize backgrounds for future experiments (diffuse νSNB, etc.)

Expected to arise from short-lived, high-Q_β isotopes (e.g. ^{88,90}Br, ^{94,96,98}Rb)

- Previous measurements biased by "pandemonium effect":





Challenge: Low rate, large contamination of cosmogenic backgrounds (isotopes, fast neutrons)

This study:

- Multivariate analysis to extract HERA prompt spectrum
- HERA $\bar{\nu}$ spectrum unfolded from prompt spectrum

Dataset

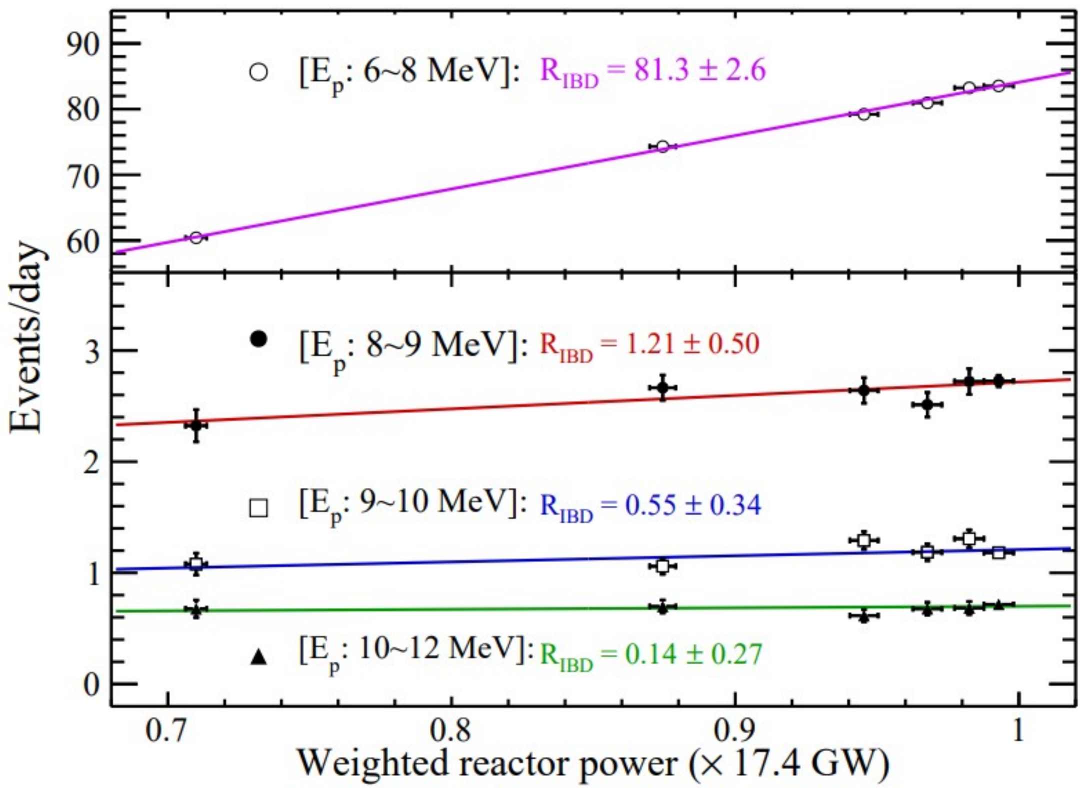
- Use 1958 days of data (Dec 24, 2011 – Aug 31, 2017)
- Standard ("A") selection with shortened (1 ms) shower veto
  - More information for constraining cosmogenic isotope bkg.

Criterion	Selection A
Calibration	^{60}Co and ^{241}Am-^{13}C method
Reconstruction	Corrected center-of-charge
8-inch PMT light emission	Reject $f_{\rm ID} \geq 0$
2-inch PMT light emission	Reject $Q_{\rm max}$(2-inch PMTs) $>$ 100 p.e.
Prompt energy	(0.7, 12.0) MeV
Delayed energy	(6, 12.0) MeV
Prompt-delayed Δt	(1, 200) μs
Multiplicity veto *(pre)*	No signal >0.7 MeV 200 μs before prompt
Multiplicity veto *(post)*	No signal >0.7 MeV 200 μs after delayed
Water Shield muon veto	Veto (-2, 600) μs after NHIT $>$ 12 in OWS or IWS
AD muon veto	Veto (0, 1) ms after >20 MeV signal
AD shower veto	Veto ~~(0, 1) s~~ (0, 1) ms after >2.5 GeV signal

- 4M IBD candidates, ~9000 in 8-12 MeV prompt energy region

The challenge

$$R = R_{\mathrm{IBD}} P_{\mathrm{reactor}} + R_{\mathrm{bkg}}$$



- Rate in 8-12 MeV 1/20th of rate in 6-8 MeV
- $R_{\mathrm{IBD}} > 0$ at $>30\sigma$ for 6-8 MeV; $<2.5\sigma$ above 8 MeV
- Backgrounds: Cosmogenic isotopes (^{9}Li/^{8}He correlated; ^{12}B+^{12}B accidental), fast neutrons

Strategy:

- **Get the fractions of IBDs and backgrounds** in each energy bin with an unbinned fit over all events in the bin
- **Get the IBD prompt spectrum** by scaling the total spectrum by the IBD fraction in each bin
- **Get the uncertainty** both from statistics and and from systematics (pull term) in background model
- **Validate the model** by defining a discriminator P_{IBD}. Compare its distribution to a toy MC prediction.
- **Unfold the prompt spectrum** to get the antineutrino spectrum

Model PDF

In each energy bin, let r_p be the fraction of events of type *p:*

- Cosmogenic isotopes
- Fast neutrons
- IBDs

Then in each energy bin, we want the best fit of the vector **r** (of r_p for all three *p*)

For each IBD candidate, determine:

- **Δt**, vector of time differences to most-recent muon in each of eight muon categories (four energy ranges × with/without associated neutron)
- *z*, vertical position of event
- *w*, baseline-weighted reactor power

Then define event-by-event probability distribution function (PDF), parameterized by **r**:

$$F(\boldsymbol{r}; \boldsymbol{\Delta t}, z, w) = \sum_p r_p f_p(\boldsymbol{\Delta t}) h_p(z) k_p(w)$$

where f_p, h_p, and k_p are the 1D PDFs of **Δt**, *z*, and *w* for events of type *p*, respectively

1D Sub-PDFs

Time-to-last-muon f(Δt):

$$f(\Delta t) = \kappa \cdot e^{-\kappa \Delta t}$$

where $\kappa = R_\mu$ for muon-uncorrelated events (IBDs and fast neutrons) and $\kappa = R_\mu + \frac{n}{\tau}$ for cosmogenic isotopes (τ = lifetime; n = 1 for ^{9}Li/^{8}He, n = 2 for ^{12}B+^{12}B)

- Eight muon categories in **Δt**: (four energy ranges × with/without water pool neutron)
- For lowest-energy muons *without WP neutrons* in EH1 and EH2, high muon rate leads to degenerate time constants
  - Use result for neutron-tagged muons, scaled by the neutron tagging efficiency from EH3 (validated with Geant4, constrained by pulls in fit)

Vertical position h(z):

- Uniform for IBDs and cosmogenic isotopes
- Top-dominated for fast neutrons (measured from IBD candidates in 12-20 MeV)

Weighted reactor power k(w):

- Proportional to *w* for IBDs
- Flat for cosmogenic isotopes and fast neutrons

$$F(\boldsymbol{r}; \boldsymbol{\Delta t}, z, w) = \sum_p r_p f_p(\boldsymbol{\Delta t}) h_p(z) k_p(w)$$

Fitting event type ratios

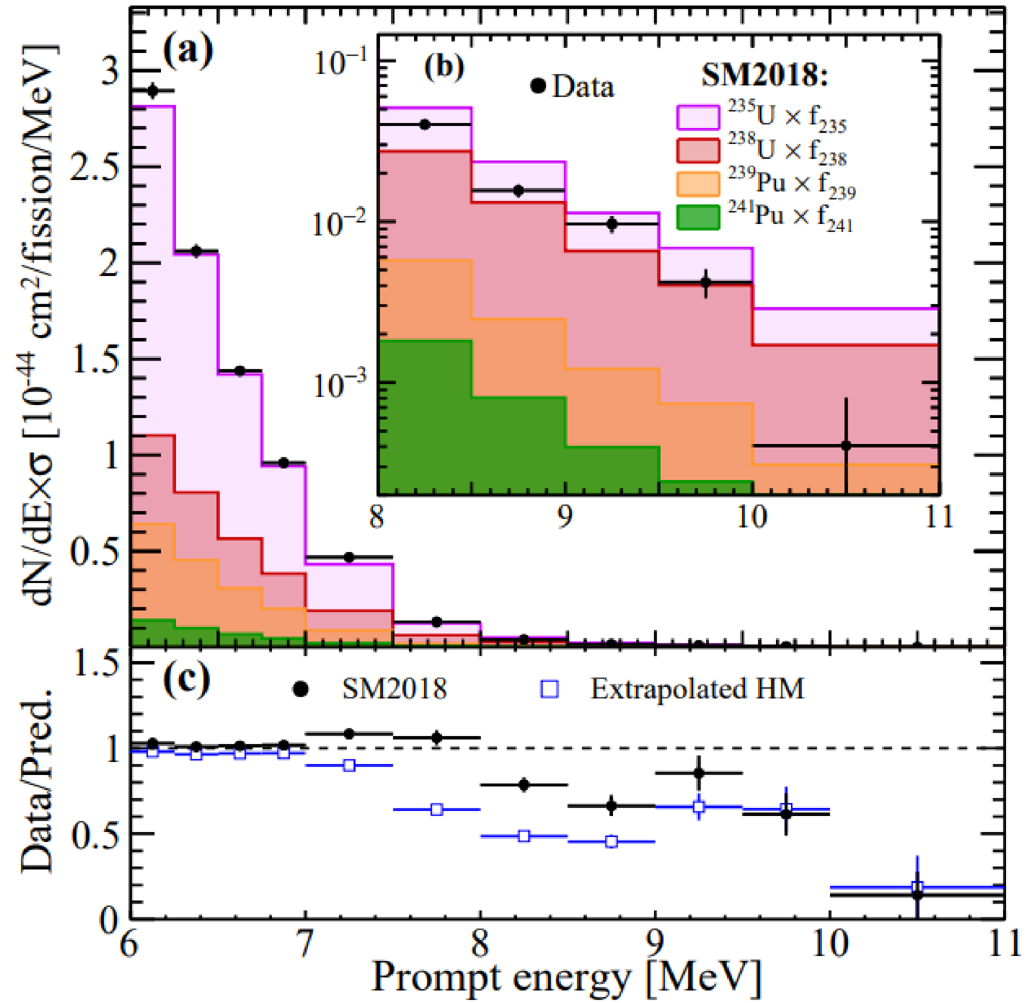
Fitting: Find **r** (for each energy bin) that minimizes

$$\chi^2(\boldsymbol{r}) = -2\sum\left[\log F(\boldsymbol{r}; \boldsymbol{\Delta t}, z, w)\right] + g(\boldsymbol{\epsilon})$$

where the sum runs over all IBD candidates and $g(\boldsymbol{\varepsilon})$ is the penalty for the systematic pulls $\boldsymbol{\varepsilon}$ (which modify the 1D PDFs f_p etc.)

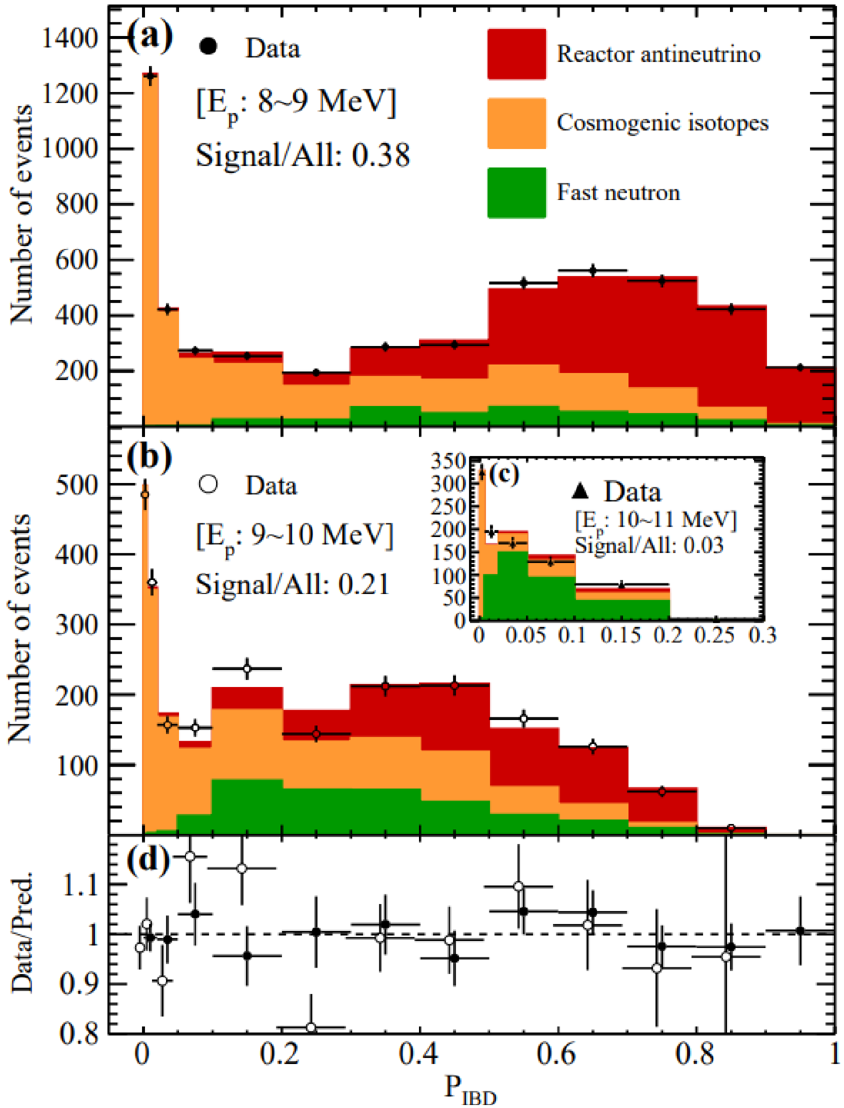
Observations:

- Measured IBD yield 3% larger than SM2018 (recent summation prediction) in 6-8 MeV, but 29% smaller for 8-11 MeV
- Above 10 MeV, expected 5σ significance based on SM2018, but only observed 1σ
- Deficits vs SM2018 consistent with pandemonium effect
- Extrapolated Huber-Mueller model strongly disagrees with data ⟶ not recommended



Prompt spectrum $S_{\mathrm{IBD}}(E) = r_{\mathrm{IBD}}(E)\, S_{\mathrm{all}}(E)$

Validation of model



P_{IBD} shows that reactor antineutrinos can be statistically separated from backgrounds (worse at higher energy).

Define, for each event, the discriminator P_{IBD} (probability of being an IBD):

$$P_{\mathrm{IBD}} = \frac{r_{\mathrm{IBD}} f_{\mathrm{IBD}}(\boldsymbol{\Delta t}) h_{\mathrm{IBD}}(z) k_{\mathrm{IBD}}(w)}{F(\boldsymbol{r}; \boldsymbol{\Delta t}, z, w)}$$

Use toy MC to predict P_{IBD} distribution for each event class and energy bin

Compare measured distribution to sum (over event classes) of predictions. Good agreement.

Cross check: Obtain consistent results when ignoring (one of) **Δt**, *z*, *w*, or when only including a single hall

Background model contributes bulk of uncertainty:

- Below 9.5 MeV, ~60% of total uncertainty comes from isotopes
- Above 10 MeV, ~70% comes from fast neutrons
- Statistics only ~20% of total

Unfolded antineutrino spectrum

Unfold the prompt spectrum using a fitting procedure. Minimize:

$$\chi^2 = (\boldsymbol{P} - \boldsymbol{M})\mathbf{Cov}^{-1}(\boldsymbol{P} - \boldsymbol{M})^T$$

M is the measured prompt spectrum, Cov is its covariance, and

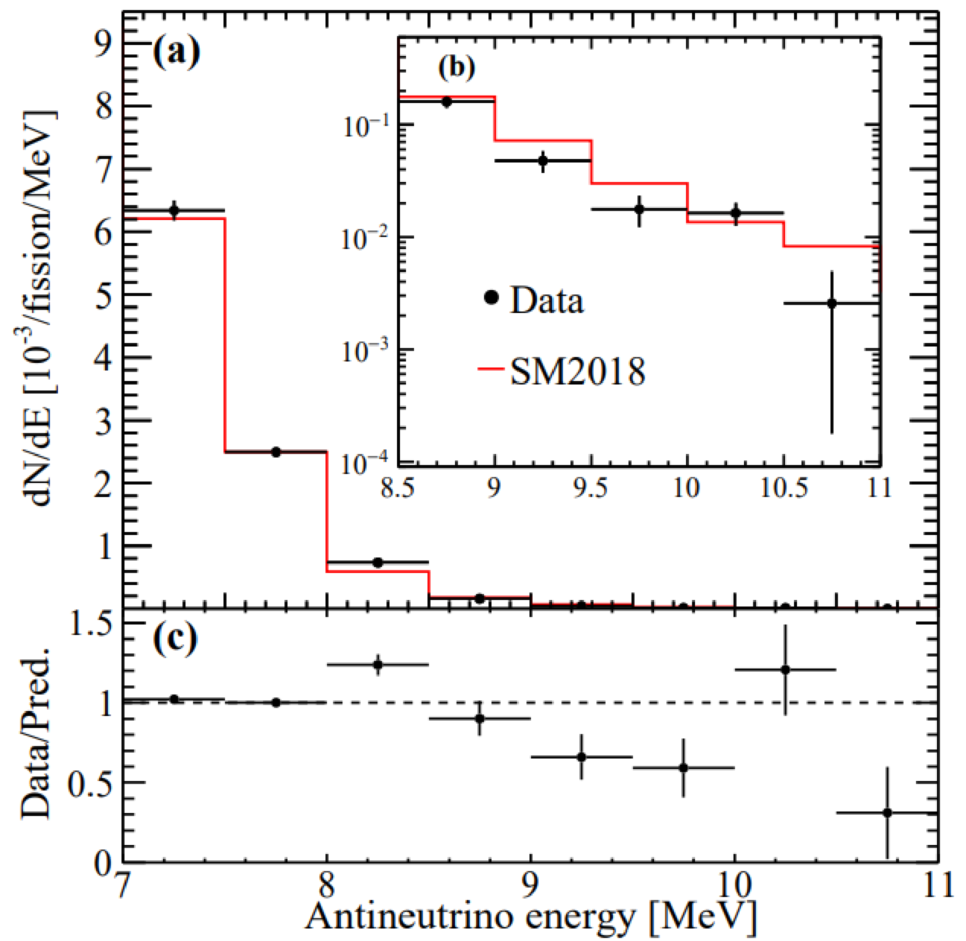
$$\boldsymbol{P} = \boldsymbol{R}\boldsymbol{S}^{\text{fit}}$$

is the predicted prompt spectrum, based on the detector response matrix **R** and the predicted antineutrino spectrum $\mathbf{S}^{\text{fit}}$ (varied in the minimization)

Unfolding procedure is equivalent to matrix inversion, but allows more statistical tests

Use wide bins to avoid large statistical fluctuations and bin-to-bin anticorrelations (common pitfall in unfolding)

Spectrum for immediate use by the community



- Assuming no HERA above (10, 10.5, 11) MeV, the χ^2 is found to be (38.3, 1.6, 0.03)
- **6.2σ rejection of no-HERA above 10 MeV**
- Above 10.5 MeV, no-HERA hypothesis not rejected

Wrapping up

More physics from Daya Bay

- Some highlights:
  - Search for events associated with gravitational waves
  - Fuel evolution measurement
  - Unfolded antineutrino spectra
  - Joint measurement of antineutrino spectra w/ PROSPECT
  - Joint sterile neutrino search w/ MINOS/MINOS+
  - Measurement of seasonal variation of muon flux
  - Oscillation analysis using neutron capture on hydrogen
  - ...and more!
- See backup slides!
- http://dayabay.ihep.ac.cn/twiki/bin/view/Public/DybPublications

Conclusion

- With full 3158-day dataset, Daya Bay has measured $\sin^2 2\theta_{13}$ and $\Delta m^2_{32}/\Delta m^2_{ee}$ to precisions of 2.8% and 2.3%, respectively
  - Likely to remain the most precise measurement of $\sin^2 2\theta_{13}$ for the foreseeable future
- Using 1958 days of data, Daya Bay has provided the first detailed measurement of high-energy reactor antineutrinos
  - Reactor antineutrinos above 10 MeV observed at 6.2σ
- Many other significant recent results
- Some upcoming results
  - Spectral oscillation analysis with neutron capture on hydrogen
  - Updated sterile neutrino search
  - Updated fuel evolution measurement