Rencontres du Vietnam 2018 August 10th, 2018 CP Violation in Heavy-Flavour Hadrons





Jörg Marks, Heidelberg University on behalf of the LHCb collaboration



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- CP symmetry describes the invariance under charge conjugation C and parity operation P
- ➤ Diagonalizing Yukawa matrices lead to quark mass eigenstates which are mixed by virtue of the weak interactions → CKM matrix



Photo: R. Hahn, Sandbox Studio

$$\begin{pmatrix} |d'\rangle\\|s'\rangle\\|b'\rangle \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub}\\V_{cd} & V_{cs} & V_{cb}\\V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{\text{CKM}}} \cdot \begin{pmatrix} |d\rangle\\|s\rangle\\|b\rangle \end{pmatrix}$$

$$d$$
 V_{ud} u W



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- Mixing between families is CKM suppressed
- In case of 3 families elements are complex



- CP symmetry describes the invariance under charge conjugation C and parity operation P
- ➤ Diagonalizing Yukawa matrices lead to quark mass eigenstates which are mixed by virtue of the weak interactions → CKM matrix



Photo: R. Hahn, Sandbox Studio

$$V_{\rm CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho(i\eta)) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho(i\eta)) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

Wolfenstein parametrization $\lambda \approx 0.22$

Only source of CPV in SM



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Only source of CPV in SM

Decays of heavy flavoured particles are well suited to test the CKM mechanism and search for non-SM contributions to CP violation.
→ focus is here on beauty and charm hadrons



Mixing in Neutral Mesons

Neutral mesons (K, D, B, B_s) are created as flavour eigenstates of the strong interaction. They can mix through weak interactions. $|P^0\rangle \xrightarrow[d, \overline{u}, \overline{d}, \overline{s}]{} |\bar{P}^0\rangle \xrightarrow[d, \overline{u}, \overline{d}, \overline{s}]{} |\bar{P}^0\rangle$

The time evolution is obtained by

$$i\frac{\partial}{\partial t} \begin{pmatrix} P^0(t) \\ \bar{P}^0(t) \end{pmatrix} = \begin{bmatrix} \begin{pmatrix} M_{11} & M_{12} \\ M_{12}^* & M_{22} \end{pmatrix} - \frac{i}{2} \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{12}^* & \Gamma_{22} \end{pmatrix} \end{bmatrix} \begin{pmatrix} P^0(t) \\ \bar{P}^0(t) \end{pmatrix}$$

> The physical eigenstates are P_L and P_H :

$$|P_{L,H}\rangle = p|P^{0}\rangle \mp q|\bar{P}^{0}\rangle$$
$$|P_{L,H}(t)\rangle = e^{-i(m_{L,H} - i\Gamma_{L,H}/2)t}|P_{L,H}(t=0)\rangle$$

> Define mass and lifetime differences of P_L and P_H :

$$x = \frac{\Delta m}{\Gamma} = \frac{m_H - m_L}{\Gamma}$$
 $y = \frac{\Delta \Gamma}{2\Gamma} = \frac{\Gamma_H - \Gamma_L}{2\Gamma}$

$$\Gamma = \frac{\Gamma_L + \Gamma_H}{2}$$



Mixing in Neutral Mesons

Neutral mesons (K, D, B, B_s) are created as flavour eigenstates of the strong interaction. They can mix through weak interactions. $|P^0\rangle \xrightarrow[d, \overline{u}, \overline{d}, \overline{s}]{} |\bar{P}^0\rangle \xrightarrow[d, \overline{u}, \overline{d}, \overline{s}]{} |\bar{P}^0\rangle$

 $> B_s^0 - \bar{B}_s^0$ oszillation frequency Δm_s measurement in $B_s^0 \to D_s^- \pi^+$ decays



 $\Delta m_s = 17.768 \pm 0.023 \, (stat) \pm 0.006 \, (syst) \, ps^{-1}$



Classification of CP Violation

For CP violation need 2 interfering amplitudes with different strong $~\delta~$ and weak $\phi~$ phases

Amplitudes of the decay $P \to f$ and their conjugates are $\mathcal{A}_f = \langle P | \mathcal{H} | f \rangle$, $\bar{\mathcal{A}}_{\bar{f}} = \langle \bar{P} | \mathcal{H} | \bar{f} \rangle$, $\mathcal{A}_{\bar{f}} = \langle P | \mathcal{H} | \bar{f} \rangle$ and $\bar{\mathcal{A}}_f = \langle \bar{P} | \mathcal{H} | f \rangle$

- ➤ CP violation in decay (direct CPV) $\Gamma(P \to f) \neq \Gamma(\bar{P} \to \bar{f}) \qquad |A_f| \neq |\bar{A}_{\bar{f}}|$
- CP violation in mixing (indirect CPV)
 - $|\frac{q}{p}| \neq 1$
- ➤ CP violation in interference of mixing and decay (indirect CPV) $arg(\lambda_f) \neq arg(\lambda_{\bar{f}}) , \qquad \lambda_f \equiv \frac{q}{p} \frac{\bar{\mathcal{A}}_f}{\mathcal{A}_f}$



Selected results

CP violation in Mixing



Introduction - CPV in $B^0_{(s)}$ Mixing

Semileptonic $B_{(s)}^0$ decays are tree-dominated transitions, CP-conserving and flavour specific.

➤ Provide theoretically clean measurements of CP violation in $B^0_{(s)}$ mixing by comparing probabilities $\mathcal{P}(B^0_{(s)} \to \bar{B}^0_{(s)})$ and $\mathcal{P}(\bar{B}^0_{(s)} \to B^0_{(s)})$

•
$$a_{sl}^q = \frac{\Gamma(\bar{B}_{(s)}^0 \to f) - \Gamma(B_{(s)}^0 \to \bar{f})}{\Gamma(\bar{B}_{(s)}^0 \to f) + \Gamma(B_{(s)}^0 \to \bar{f})} \approx \frac{\Delta\Gamma_{(s)}}{\Delta m_{(s)}} \tan \phi_{12}^q$$

- Standard Model predictions are small $a_{sl}^s = (2.2 \pm 0.27) \cdot 10^{-5}$ $a_{sl}^d = (-4.7 \pm 0.6) \cdot 10^{-4}$
- PRD 83 (2011) 036004 arXiv:1102.4274

Two experimental methods

Same-sign dilepton asymmetry

$$a_{sl} = \frac{\mathcal{P}_{ll}^{++} - \mathcal{P}_{ll}^{--}}{\mathcal{P}_{ll}^{++} + \mathcal{P}_{ll}^{--}}$$

Untagged t dependent asymmetry

$$\frac{N(\bar{B},t) - N(B,t)}{N(\bar{B},t) + N(B,t)} = \frac{a_{sl}}{2} \left(1 - \frac{\cos(\Delta mt)}{\cosh(\frac{1}{2}\Delta\Gamma t)}\right)$$



BABAR - CP Violation in B^0 Mixing

Measure the asymmetry A_{CP} in the dilepton same-sign rates using the full BABAR dataset
Phys. Rev. Lett. 114, 081801 (2015)

•
$$A_{CP} = \frac{\mathcal{P}_{ll}^{++} - \mathcal{P}_{ll}^{--}}{\mathcal{P}_{ll}^{++} + \mathcal{P}_{ll}^{--}}$$

• Allow for 4 lepton comb. $l_1 l_2 = \{ee, e\mu, \mu e, \mu \mu\}$ $\int_{-20}^{0} -10$ $\int_{-20}^{-10} -20$ $\int_{-30}^{0} -30$ -40

10

• Asymmetry in $B^0-ar{B}^0$ mixing from semileptonic decays in $\,\Upsilon(4S)\,$

$$A_{CP} = (-3.9 \pm 3.5 \pm 1.9) \cdot 10^{-3}$$

HFLAV $A_{CP}^{\Upsilon(4S)} = (-1.9 \pm 2.7) \cdot 10^{-3}$

Subsample

ee

 $e\mu$

 μe

 $\mu\mu$

 $\underbrace{l^+}_{\mathbf{Y}} \gg \bar{B} \underbrace{\uparrow}_{\mathbf{X}} \underbrace{\uparrow}_{\mathbf{X}} B \xrightarrow{I}_{\mathbf{X}} \underbrace{I^+}_{\mathbf{X}} \underbrace{I^+}_{\mathbf{$

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LHCb - CP Violation in B_s^0 Mixing

Phys. Rev. Lett. 117, 061803 (2016)

> Measure a_{sl}^s in semileptonic $B_s^0 \to D_s^- (\to K^+ K^- \pi^-) \mu^+ \nu_{\mu} X$ decays

• No determination of the initial flavour and integrate over the B_s^0 decay time

$$A_{raw} = \frac{N(D_s^-\mu^+) - N(D_s^+\mu^-)}{N(D_s^-\mu^+) + N(D_s^+\mu^-)} = \frac{a_{sl}}{2}$$

- No dependency on the production rate of B_s^0 and \bar{B}_s^0 due to the high mixing frequency of Δm_s
- Fit D⁻_s mass for various subsamples (magnet polarity, Dalitz region, data taking period)
- Asymmetry in B_s^0 mixing

 $a_{sl}^s = (0.39 \pm 0.26 \pm 0.20)\%$





Summary of CPV in $B_{(s)}^0$ Mixing

> Measurements of a_{sl}^d and a_{sl}^s are fully compatible with SM predictions





Selected results

Direct CP violation



Direct CPV in Charmless $B_{(s)}^0$ Decays

- ➤ The interference between $b \to u$ tree and $b \to s(d)$ loop processes gives access to direct CP violation in $B_{(s)}^0 \to K\pi$ decays
 - Measure the t-integrated asymmetries $A_{CP} = \frac{N_{\bar{B} \to \bar{f}} - N_{B \to f}}{N_{\bar{B} \to \bar{f}} + N_{B \to f}}$ $B \to f = \begin{pmatrix} B^0 \to K^+ \pi^- \\ B_s^0 \to K^- \pi^+ \end{pmatrix}$
 - Test Standard Model prediction in a model independent way by

$$\Delta = \frac{A_{CP}(B^0 \to K^+\pi^-)}{A_{CP}(B^0_s \to K^-\pi^+)} + \frac{\mathcal{B}(B^0_s \to K^-\pi^+)}{\mathcal{B}(B^0 \to K^+\pi^-)} \frac{\tau_d}{\tau_s} \underbrace{=}_{\mathsf{SM}} 0 \qquad \begin{array}{c} \mathsf{H.J. \ Lipkin, \ Phys. \ Lett.}\\ \mathsf{B \ 621 \ (2005) \ 126} \end{array}$$

$$\mathsf{LHCb \ measurements} \qquad \qquad \mathsf{external \ averages} \qquad \qquad \begin{array}{c} \mathsf{SM} \end{array}$$



CP Violation in $B^0 \to K^+\pi^-$ Decays

First CP asymmetry measurements by B factory experiments in 2004

• First evidence CPV in $B^0 \to K^+ \pi^-$



 Results after updates to the full datasets compared to CDF and LHCb





Direct CP Violation in B_s^0 Decays





Direct CP Violation in $B_{(s)}^0$ Decays



Results after correcting for detection and production asymmetries

• $A_{CP}(B_s^0 \to K^- \pi^+) = 0.213 \pm 0.015 \, (stat) \pm 0.007 \, (sys)$

5.8

•
$$A_{CP}(B^0 \to K^+\pi^-) = -0.084 \pm 0.004 \,(stat) \pm 0.003 \,(sys)$$

HFLAV averages + JHEP 04 (2013) 001 arXiv:1612.07233

 $\Delta = -0.11 \pm 0.04 \pm 0.03 \rightarrow$ compatible with SM expectation



Candidates / ($5 \text{ MeV}/c^2$

8

10

Decay time [ps]

12

2

Selected results

CP Violation in the Interference of Mixing and Decay



First Observation of CPV in B System

CP violation in the B^0 System was observed in 2001 at the B factories in B^0/\bar{B}^0 decays to a CP eigenstate q/p

 \blacktriangleright The interference between $B^0 \rightarrow J/\Psi K_s^0$ decays and $B^0 \rightarrow \bar{B}^0 \rightarrow J/\Psi K_s^0$ mixed decays gives access to CP violation in the Interference of mixing and decay



PRL 87 (2001) 091802

Measure time-dependent asymmetry to get the CP violating phase

$$\mathcal{A}(t) \equiv \frac{\Gamma(\bar{B}^0(t) \to J/\Psi K_s^0) - \Gamma(B^0(t) \to J/\Psi K_s^0)}{\Gamma(\bar{B}^0(t) \to J/\Psi K_s^0) + \Gamma(B^0(t) \to J/\Psi K_s^0)}$$
$$\approx \sin(2\beta)\sin(\Delta m_{(s)}t))$$

Requires flavor tagging







Summary - $sin(2\beta)$ Measurements

World average as compiled by the Heavy FLavour AVeraging group





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Summary - $sin(2\beta)$ Measurements

World average as compiled by the Heavy FLavour AVeraging group

- Compatible with indirect prediction from CKMfitter(2016) $sin(2\beta_{indirect}) = 0.740^{+0.020}_{-0.025}$
- Uncertainty of the LHCb result with $\int \mathcal{L} = 3fb^{-1}$ is competitive with Belle and BABAR JHEP 11 (2017) 170
- Current average allows for two solutions in $[0, \pi]$ $\sin(2\beta_{HFLAV}) = 0.699 \pm 0.017$

$$\beta = (22.2 \pm 0.7)^{\circ} \quad \beta = (67.8 \pm 0.7)^{\circ}$$

$$\sin(2\beta) \equiv \sin(2\phi_1) \overset{\text{HFLAV}}{\underset{\text{PRELIMINARY}}{\text{Moriond 2018}}}$$





Summary - $sin(2\beta)$ Measurements

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- Current average provides two solutions in $[0, \pi]$ $\sin(2\beta_{HFLAV}) = 0.699 \pm 0.017$ $\beta = (22.2 \pm 0.7)^{\circ}$ $\beta = (67.8 \pm 0.7)^{\circ}$ 7.3σ recent combined Belle and BABAR analysis measures $\cos(2\beta) > 0$ arXiv:1804.06152, arXiv:1804.06153

$$\sin(2\beta) \equiv \sin(2\phi_1) \underset{\text{PRELIMINARY}}{\text{HFLAV}}$$

$$0.69 \pm 0.03 \pm 0.01$$





Similar to the B^0 system the interference of mixing and decay gives access to the CP violating phase $\phi_s \equiv -2\beta_s$ in the B_s^0 system



then by LHCb, Atlas and CMS

Access new physics contributions by deviations from SM

•
$$\phi_s = \phi_s^{SM} + \Delta \phi_s$$
, $\Delta \phi_s = \arg(M_{12}/M_{12}^{SM})$

• Get ϕ_s^{SM} from global fit ignoring penguin contribution, CKM fitter arXiv:1501.05013 $\phi_s^{SM} = -0.0364 \pm 0.0016 \ rad$

LHCb -
$$\phi_s$$
, $\Delta \Gamma_s$ in $B_s \rightarrow J/\psi \phi$

Need an angular analysis to statistically separate CP eigenstates $\phi \rightarrow K^+ K^-$ in P wave \rightarrow CP-even, CP-odd helicity angles $\Omega = (\theta_{\mu}, \theta_{K}, \phi_{h})$ $\phi \to K^+ K^-$ in S wave $\to \mathsf{CP}\text{-}\mathsf{odd}$ depending on rel. orbital momentum of J/ψ and ϕ $\underline{d^4\Gamma(B^0_s\to J\psi KK)}$ > LHCb determines in a max. likelihood fit to $dt d\Omega$

- the physics quantities ϕ_s , $\Delta\Gamma_s$, Γ_s ,
 - Key ingredients to t dependent flavour tagged angular analysis
 - probability of getting the initial B flavour wrong
 - decay time measurement
 - event by event decay time resolution
 - knowledge of Δm_s





LHCb - ϕ_s , $\Delta \Gamma_s$ in $B_s \rightarrow J/\psi \phi$

Need an angular analysis to statistically separate CP eigenstates $\phi \to K^+K^-$ in P wave → CP-even, CP-odd $\phi \to K^+K^-$ in S wave → CP-odd
depending on rel. orbital
momentum of J/ψ and ϕ LHCb determines in a max. likelihood fit to $d^4\Gamma(B_s^0 \to J\psi KK))$

the physics quantities ϕ_s , $\Delta\Gamma_s$, Γ_s ,



CP-even CP-odd S-wave



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Summary - ϕ_s , $\Delta \Gamma_s$ in $B_s \rightarrow J/\psi \phi$

► LHCb results for $B_s \to J/\psi K^+ K^-, J/\psi \pi^+ \pi^-, J/\psi \phi (K^+ K^-)$ JHEP 08 (2017) 037 $\phi_s = -25 \pm 45 \pm 8 \, mrad$ $\Delta \Gamma_s \equiv \Gamma_L - \Gamma_H = 0.0813 \pm 0.0073 \pm 0.0036 \, ps^{-1}$ $\Gamma_s = 0.6588 \pm 0.0022 \pm 0.0015 \, ps^{-1} \, |\lambda| = 0.978 \pm 0.013 \pm 0.003$ no evidence for CPV

- Measurements include the mass range above $\phi(1020)$
- $\nabla \Gamma_{sd}^{s} g^{0.14}$ HFLAV D0 8 fb⁻¹ PDG 2018 68% CL contours $(\Delta \log \mathcal{L} = 1.15)$ CMS 19.7 fb⁻¹ 0.10 CDF 9.6 fb⁻¹ 0.08 LHCb 3 fb $^{-1}$ ATLAS 19.2 fb⁻¹ 0.06 0.2 -0.2 0.4 -0.4 -0.0 $\phi_{s}^{c\bar{c}s}[rad]$
- > HFLAV average for $\phi_s^{c\overline{c}s}$ vs $\Delta\Gamma_s$
- Reached a precision of 33 mrad with LHC run 1 data, which includes also ATLAS PLB 757 (2016) 97 and CMS JHEP 08 (2016) 147 data
- Good agreement with Standard Model expectations



Selected results

Measurement of the CKM angle γ



Introduction - γ in $B \rightarrow DK$ Decays

 $b \rightarrow c$ and $b \rightarrow u$ transitions are sensitive to $\gamma \equiv arg[-V_{ud}V_{ub}^*/V_{cd}V_{cb}^*]$





Reconstruct D^0/\bar{D}^0 in a final state f_D accessible to both to achieve interference

- Interference allows to determine
 - weak phase γ
 - relative strong phase δ_B
 - magnitude ratio of the amplitudes r_B
- Hadronic parameters (r_D, δ_D) of f_D (either determined or external input)



Considered as SM benchmark since only tree level processes are involved



γ Measurement in $B \rightarrow DK$ Decays

Depending on the final state $D^0 \rightarrow f_D$ distinguish methods of γ extraction

- Measure CP asymmetries and charge averaged partial decay rate ratios \rightarrow fit for (γ, r_B, δ_B)
- > GLW (Gronau, London, Wyler) [1991]): $f_D = [K^+K^-], [\pi^+\pi^-], [K_s^0\pi^0]$
 - f_D is a suppressed CP eigenstate
- ≻ ADS (Atwood, Dunietz, Soni [1997,2001]): $f_D = [K^- \pi^+], [K^+ \pi^-]$
 - Interference of suppressed *B* / favoured *D* decays and favoured *B* / doubly suppressed *D* decays



γ Measurement in $B \rightarrow DK$ Decays

> GGSZ(Giri, Grossman, Sofer, Zupan) [2003]): $f_D = [K_s^0 \pi^+ \pi^-, K_s^0 K^+ K^-]$

- f_D is a multibody selfconjugate state, which contains different quasi 2-body final states ($D^0 \to K^{*-}\pi^+$, $D^0 \to K^{*+}\pi^-$, $D^0 \to K^0_s \rho^0$) arXiv:1806.01202
- Interference of suppressed and favoured *B* decays followed by *D* decays provide amplitude and phase variations across the Dalitz plane
- Measure yields in bins across the Dalitz plane and compare B⁺ and B⁻

Use strong *D* phase measurement, CLEO-C PRD 82 (2010) 112006





LHCb - Results of γ Measurements

> Combine 16 γ measurements to obtain global LHCb average

- \bullet Determine γ and hadronic amplitude ratio and phase difference
- To achieve optimal sensitivity, use 94 observables with a likelihood ansatz for the combination



$B^+ \to DK^+$	$D \to h^+ h^-$	GLW	Run 1 & 2
$B^+ \to DK^+$	$D \rightarrow h^+ h^-$	ADS	Run 1
$B^+ \to DK^+$	$D \to h^+ \pi^- \pi^+ \pi^-$	GLW/ADS	Run 1
$B^+ \to DK^+$	$D \to h^+ h^- \pi^0$	GLW/ADS	Run 1
$B^+ \to DK^+$	$D \to K^0_{ m s} h^+ h^-$	GGSZ	Run 1
$B^+ \to DK^+$	$D \to K_{\rm s}^0 h^+ h^-$	GGSZ	Run 2
$B^+ \to DK^+$	$D \to K^0_{\rm s} K^+ \pi^-$	GLS	Run 1
$B^+ \to D^* K^+$	$D \rightarrow h^+ h^-$	GLW	Run 1 & 2
$B^+ \to DK^{*+}$	$D \rightarrow h^+ h^-$	GLW/ADS	Run 1 & 2
$B^+ \to DK^{*+}$	$D \to h^+ \pi^- \pi^+ \pi^-$	$\mathrm{GLW}/\mathrm{ADS}$	Run 1 & 2
$B^+ \to D K^+ \pi^+ \pi^-$	$D \rightarrow h^+ h^-$	GLW/ADS	Run 1
$B^0 \to DK^{*0}$	$D \to K^+ \pi^-$	ADS	Run 1
$B^0 \rightarrow DK^+\pi^-$	$D \rightarrow h^+ h^-$	GLW-Dalitz	Run 1
$B^0 \to DK^{*0}$	$D \to K^0_{\rm s} \pi^+ \pi^-$	GGSZ	Run 1
$B_s^0 \to D_s^{\mp} K^{\pm}$	$D_s^+ \rightarrow h^+ h^- \pi^+$	TD	Run 1
$B^0 \rightarrow D^{\mp} \pi^{\pm}$	$D^+\!\to K^+\pi^-\pi^+$	TD	Run 1



LHCb - Results of γ Measurements

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- \bullet Determine γ and hadronic amplitude ratio and phase difference
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• Uncertainty of γ decreased by a factor

World Average of γ Measurements

> All experiments performed measurements with the various methods

- \bullet Determine γ and hadronic amplitude ratio and phase difference
- Uncertainties of the averages are dominated by the LHCb results





Selected results

CP violation in the Baryon Sector



CPV in Charmless 2-body Λ_b^0 **Decays**

Phys. Rev. Lett. 113, 242001 (2014)

CDF measured the first time in 2014 CP asymmetries in baryonic decays, $\Lambda_b^0 \to p\pi^-$ and $\Lambda_b^0 \to pK^-$, using the full dataset (9.3 fb^{-1})

$$A_{CP}^{dir} = \frac{\Gamma(\Lambda_b^0 \to f) - \Gamma(\bar{\Lambda}_b^0 \to \bar{f})}{\Gamma(\Lambda_b^0 \to f) + \Gamma(\bar{\Lambda}_b^0 \to \bar{f})} = \frac{N_{b \to f} - c_f N_{\bar{b} \to \bar{f}}}{N_{b \to f} + c_f N_{\bar{b} \to \bar{f}}}$$

Phys. Rev. D 80, 034011 (2009)

While theory predicted 30 % asymmetries the measurements are compatible with no direct CP violation

$$A_{CP}^{dir}(\Lambda_b^0 \to p\pi^-) = +0.06 \pm 0.07 \pm 0.03$$

$$A_{CP}^{dir}(\Lambda_b^0 \to pK^-) = -0.10 \pm 0.08 \pm 0.04$$





CPV in Charmless 2-body Λ_b^0 Decays

arXiv: 1807.06544 (2018)

> LHCb measurement of direct CP asymmetries in $\Lambda_b^0 \to p\pi^$ and $\Lambda_b^0 \to pK^-$ using all run 1 data



Most precise measurement to date, compatible with no direct CP violation

 $A_{CP}^{p\pi^{-}} = -0.035 \pm 0.017 \pm 0.020$ $A_{CP}^{pK^{-}} = -0.020 \pm 0.013 \pm 0.019$

 $A_{CP}^{pK^{-}} - A_{CP}^{p\pi^{-}} = 0.014 \pm 0.022 \pm 0.013$

cancels most systematic uncertainties



CPV in $\Lambda_b^0 \to p\pi^-\pi^+\pi^-$ **Decays**

Nature Physics 13,391-396 (2017)

- ➢ Measure P and CP violation observables $a_{P/CP}^{T-odd} = 1/2(A_{\hat{T}} \pm \bar{A}_{\hat{T}})$ using triple products of final state particle momenta in Λ_b⁰ → pπ⁻π⁺π⁻ decays in the Λ_b⁰ c.m.s. Non-zero values indicate P / CP violation.

$$\Lambda_b^0: \ C_{\hat{T}} \equiv \vec{p}_p \cdot (\vec{p}_{\pi^-} \times \vec{p}_{\pi^+})$$

$$\bar{\Lambda}^0_b: \ \bar{C}_{\hat{T}} \equiv \vec{p}_{\bar{p}} \cdot (\vec{p}_{\pi^+} \times \vec{p}_{\pi^-})$$





CPV in $\Lambda_b^0 \to p \pi^- \pi^+ \pi^-$ Decays

- Two amplitudes of comparable magnitude interfere and give access to CP violation $\Lambda_b^0 \begin{bmatrix} u & & & & & \\ 0 &$
- ➤ Measure P and CP violation observables $a_{P/CP}^{T-odd} = 1/2(A_{\hat{T}} \pm \bar{A}_{\hat{T}})$ using triple products of final state particle momenta in Λ_b⁰ → pπ⁻π⁺π⁻ decays in the Λ_b⁰ c.m.s. Non-zero values indicate P / CP violation.

$$\begin{split} \Lambda_b^0: \ C_{\hat{T}} &\equiv \vec{p}_p \cdot (\vec{p}_{\pi^-} \times \vec{p}_{\pi^+}) \qquad A_{\hat{T}} \equiv \frac{N(C_{\hat{T}} > 0) - N(C_T < 0)}{N(C_{\hat{T}} > 0) + N(C_{\hat{T}} < 0)} \\ \bar{\Lambda}_b^0: \ \bar{C}_{\hat{T}} \equiv \vec{p}_{\bar{p}} \cdot (\vec{p}_{\pi^+} \times \vec{p}_{\pi^-}) \qquad \bar{A}_{\hat{T}} \equiv \frac{\bar{N}(-\bar{C}_{\hat{T}} > 0) - \bar{N}(-\bar{C}_T < 0)}{\bar{N}(-\bar{C}_{\hat{T}} > 0) + \bar{N}(-\bar{C}_{\hat{T}} < 0)} \end{split}$$



CPV in $\Lambda_b^0 \rightarrow p\pi^-\pi^+\pi^-$ **Decays**

Signal Distributions



> The Λ_b^0 four body decays have a rich resonance structure dominated in $\Lambda_b^0 \to p \pi^- \pi^+ \pi^-$ by $\Delta(1232)^{++} \to p \pi^+$ and $\rho^0(770) \to \pi^+ \pi^-$ decays

CPV in $\Lambda_b^0 \to p \pi^- \pi^+ \pi^-$ Decays

CP asymmetries vary strongly over the phase space, therefore determine $a_{P/CP}^{T-odd}$ in a phase space and Φ dependent binning.

Asymmetry measurements over the entire phase space do not show P / CP violation



Selected results

Mixing and CP violation in Charm



Mixing and CPV in the D System

- \blacktriangleright SM predicts small effects $~(\sim 10^{-3})$ in mixing and CP violation
 - Complement K and B systems, access to new physics coupling to up-type sector
 - D mixing first observed at hadron colliders by LHCb in 2013
 - Established earlier by combining many B factory results



no mixing excluded at $~9.1~\sigma$

• Not a precision testing ground, since calculations are difficult due to the intermediate mass range



LHCb - CPV in $D^0 \to K^+ \pi^-$

A determination of mixing parameters for D^0 and \bar{D}^0 gives access to CPV

► Using data of 2011 – 2016 LHCb determined the time dependent $D^0 \rightarrow K^+ \pi^- / D^0 \rightarrow K^- \pi^+$ rate $R^{\pm}(t)$



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Results

- CPV in mixing 1.00 < |q/p| < 1.35 @ 68.3% CL
- direct CPV

$$A_D = \frac{R^+ - R^-}{R^+ + R^-} = (-0.1 \pm 9.1) \cdot 10^{-3}$$

No indication for direct or indirect CPV





Rencontres du Vietnam: CPV in Heavy-Flavoured Hadrons

t-dependent CPV in
$$D^0 \rightarrow K^+ K^- / \pi^+ \pi^-$$

 $\begin{array}{l} \blacktriangleright \text{ Access CPV through decay time dependent asymmetry measurements} \\ A_{CP}(f;t) \equiv \frac{\Gamma(D^0(t) \to f) - \Gamma(\bar{D}^0(t) \to f)}{\Gamma(D^0(t) \to f) + \Gamma(\bar{D}^0(t) \to f)} \approx \underbrace{a_{CP}^{dir}(f)}_{\mathsf{CPV in decay}} - \underbrace{\frac{t}{\tau}}_{\mathsf{CPV in mixing + interfer.}}^{t} \end{aligned}$



• World best measurement using $\mathcal{L}_{int} \approx 3 \ fb^{-1}$

$$A_{\Gamma}(\pi^{+}\pi^{-}) = (0.46 \pm 0.58 \pm 0.12) \cdot 10^{-3}$$

 $A_{\Gamma}(K^+K^-) = (-0.30 \pm 0.32 \pm 0.10) \cdot 10^{-3}$

No indirect CPV, level of precision 10^{-3}



Summary - CPV in Charm

LHCb takes the fine results of the B factories and Tevatron in charm physics to an even higher precision



• Experimental data indicates mixing

- No indication for CP violation
- > All results are compatible with Standard Model expectations



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Summary

Jörg Marks

- Measurements of the CKM triangle made an enormous progress in the LHC era and are a very important tool in the search for new physics
- All measurements of CP violation are consistent with the predictions of the CKM mechanism of the Standard Model
 - \rightarrow no evidence for new physics contributions
- ➢ BELLE II and LHCb Upgrade will further improve the precision in CPV measurements → stay tuned





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Backup



Contributors Beauty / Charm Physics

B physics experiments are also well suited for charm physics





Experimental Aspects

Different experimental requirements at pp and e^+e^- facilities



Asymmetry Measurements

Measure CP asymmetries at $10^{-3} \rightarrow$ control systematic uncertainties at a similar level

- ➢ Determine raw asymmetries A_{raw} by comparing partial decay rates of a process and its CP conjugate $A_{raw} = \frac{N(x) N(\bar{x})}{N(x) + N(\bar{x})}$
 - Systematic uncertainties cancel, 2 contributions in case of LHCb
 - production rate asymmetries, A_P (measurements for $B^0, B_s^0, B^+, \Lambda_b^0$) Phys. Lett. B774 (2017) 139
 - asymmetries due to particle / anti-particle detection, A_D
 - For LHCb A_P and A_D are determined by
 - use of control samples
 - regularly reverse magnet polarity
 - remove fiducial volumes contributing to asymmetries

 \succ CP asymmetry: $A_{CP} = A_{raw} - A_P - A_D$

