

New Physics / Resonances^{/16} in Vector Boson Scattering at the LHC

25th Anniversary of the Rencontres du Vietnam

Aug 5 – 11

ICISE



WINDOWS ON THE UNIVERSE

2018



Jürgen R. Reuter, DESY



based on work with

A. Alboteanu, S. Brass, C. Fleper, W. Kilian, T. Ohl, M. Sekulla

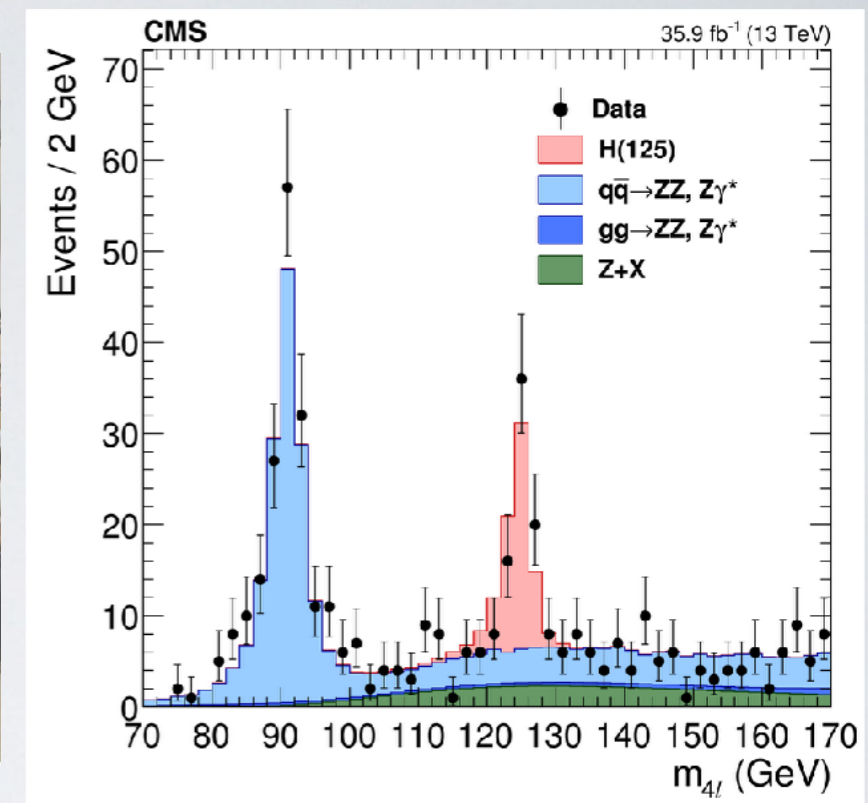
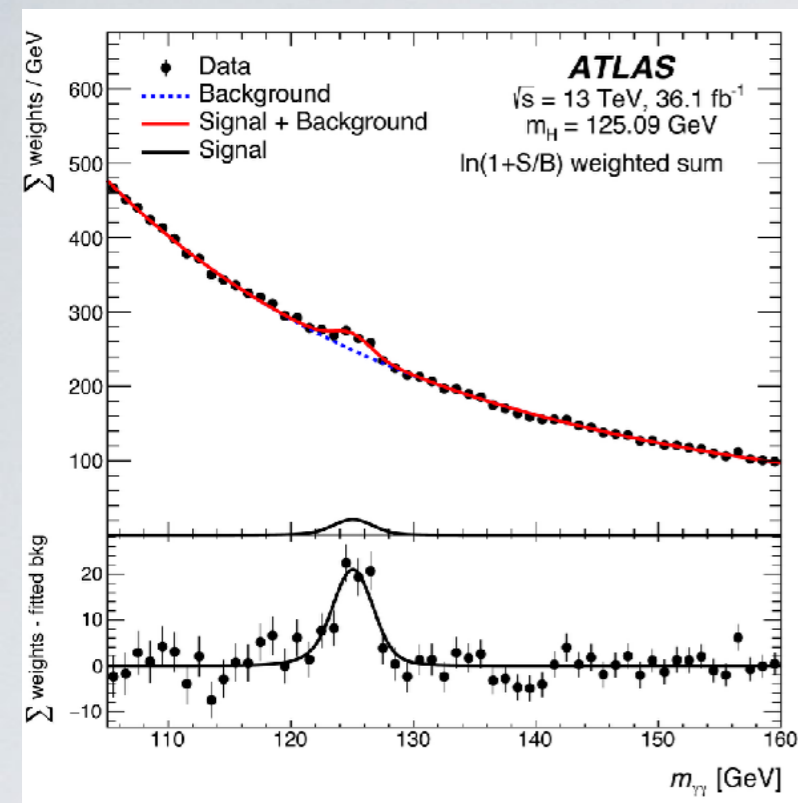
w. EPJC [1807.02512]

PRD93(16),3. 036004 [1511.00022],

PRD91(15) 096007 [1408.6207]

JHEP 0811.010 [0806.4145]

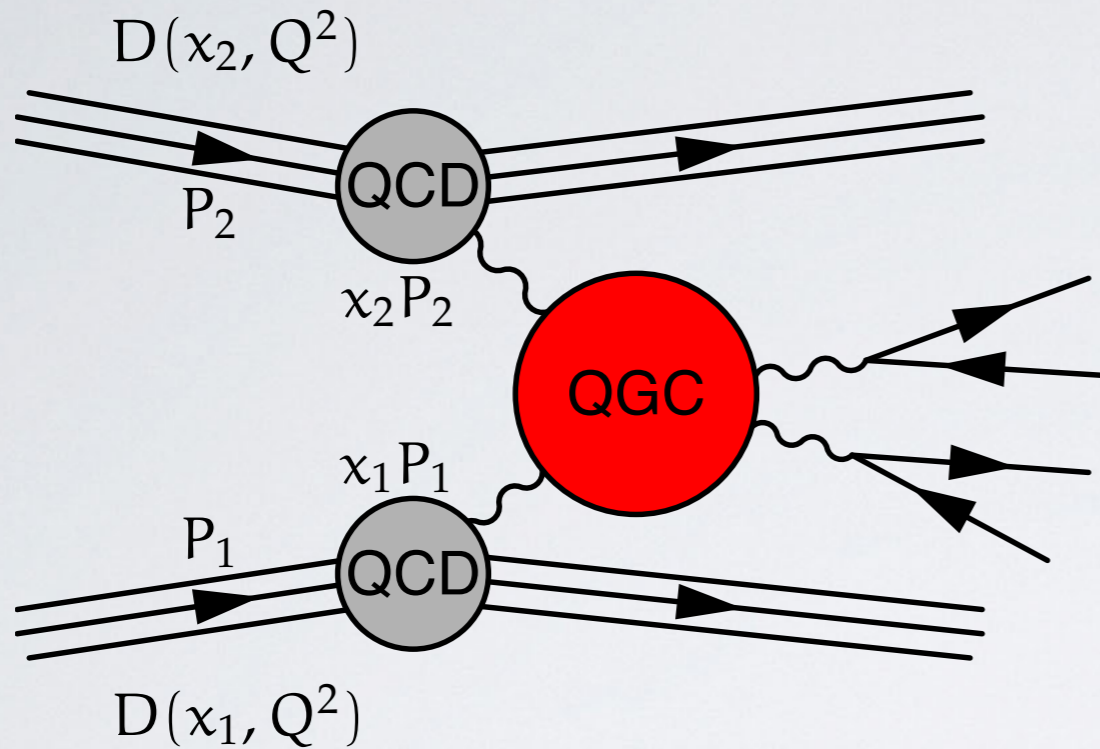




- Discovery of a light Higgs boson leaves still open questions:
 1. Nature of Electroweak Symmetry Breaking
 2. Higgs boson potential, all the way like the Standard Model!?
 3. Does “the Higgs” fulfill the US-fermion/Europe-boson rule?
 4. Is the 125 GeV state the only resonance in the system of EW vector bosons?
 5. How do EW vector bosons scatter? (true heart of weak interactions)
 6. Is there something related to the Little Hierarchy problem (strong or weak)
 7. Look for deviations in intricate cancellations of VBS amplitudes

Anatomy of Vector Boson Scattering (VBS)

$$pp \rightarrow WWjj \rightarrow \ell\nu\nu jj$$

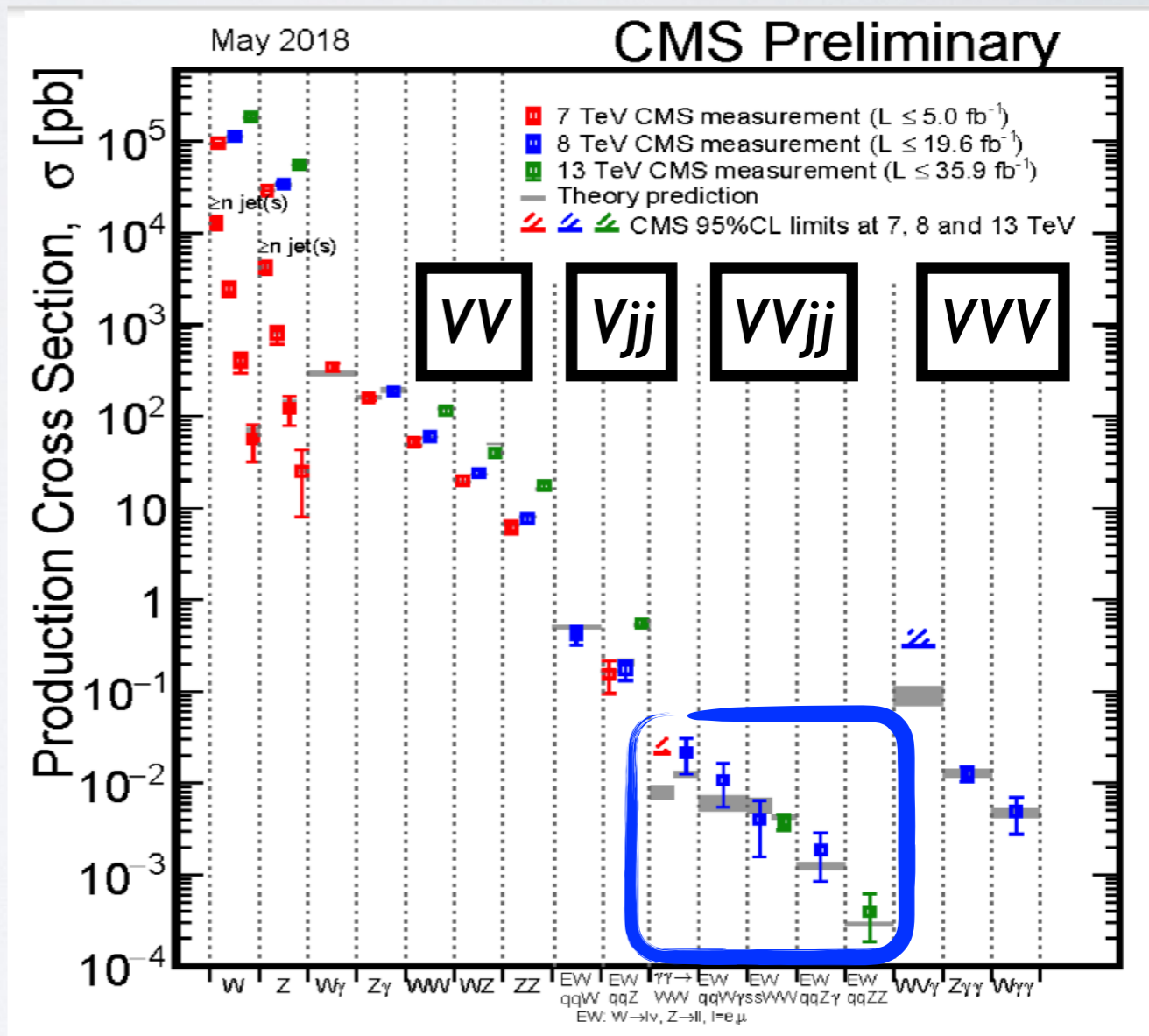


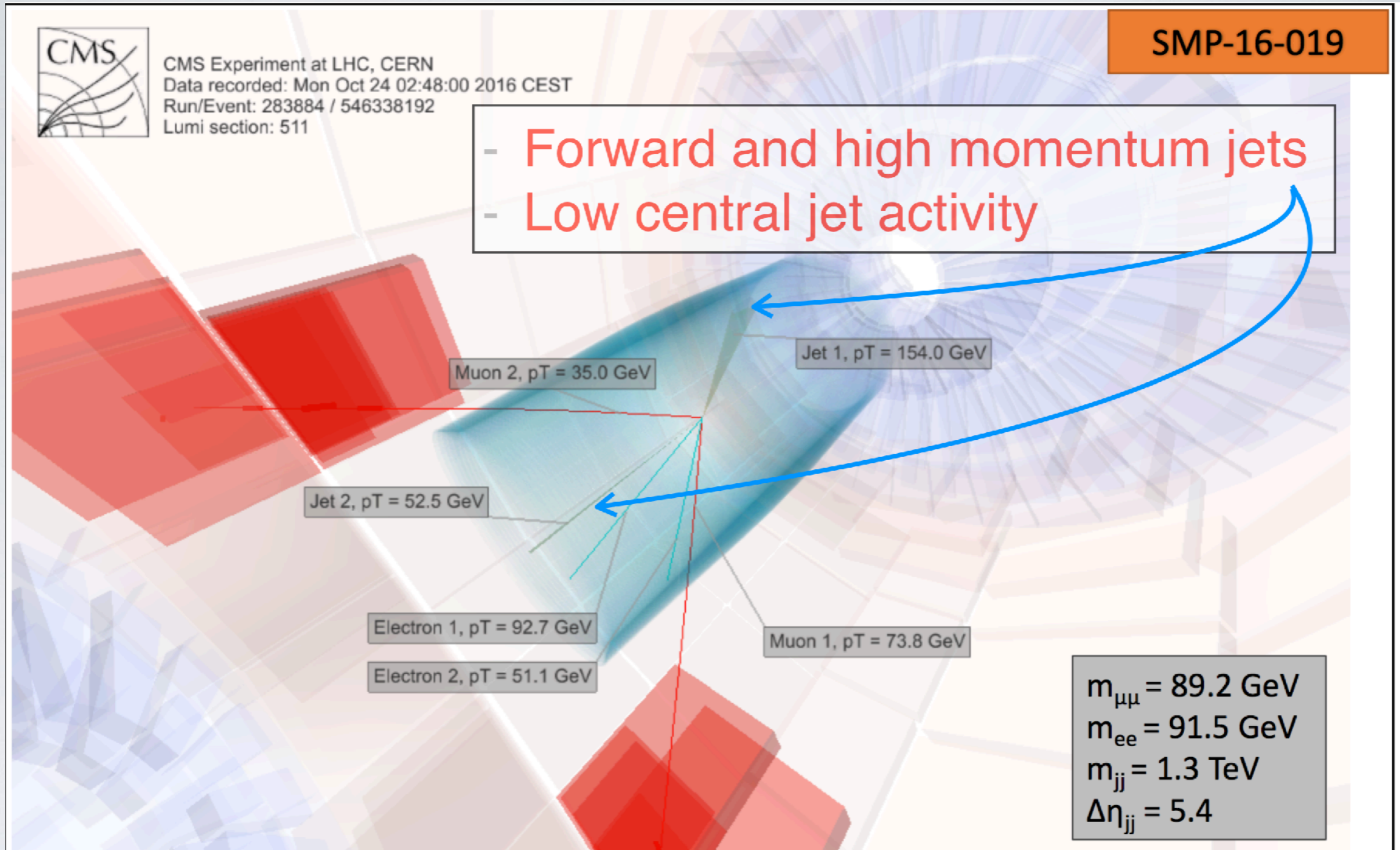
Backgrounds:

- $tt \rightarrow WbWb$
- $W + jets$
- single top, misreconstructed jet
- $WWjj$ QCD production
- $ll + X + Emiss$ (“prompt”)

Fiducial phase space volume:

- $lljj$ tag
- $m_{jj} > 500$ GeV (“jet recoil”)
- $|\Delta y_{jj}| > 2.4$ (“rapidity distance”)
- Cuts on E_j, p_T^j
- No mini jet vetoes





VBS ZZjj Candidate Event from PLB 774 (2017) 682

shown by Kenneth Long, Seoul, ICHEP 2018

The Holy Grail of Vector Boson Scattering

5 / 16

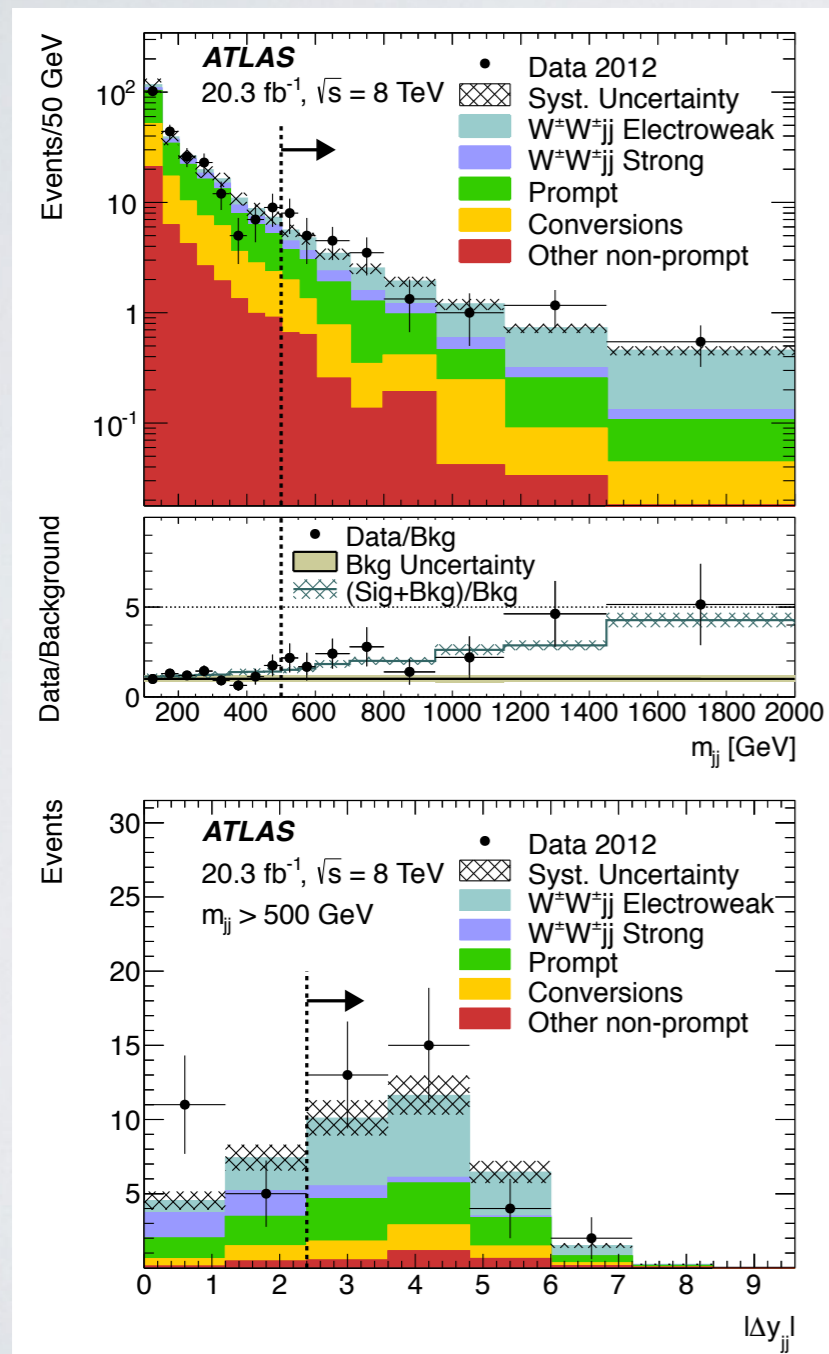
- 📌 Discovery for W^+W^+jj (electroweak production)
ATLAS PRL 113(2014)14, 141803 [1405.6241] & 1611.02428; CMS PRL 114(2015), 051801 [1410.6315]
- 📌 First limits on New Physics in pure electroweak gauge/Goldstone sector

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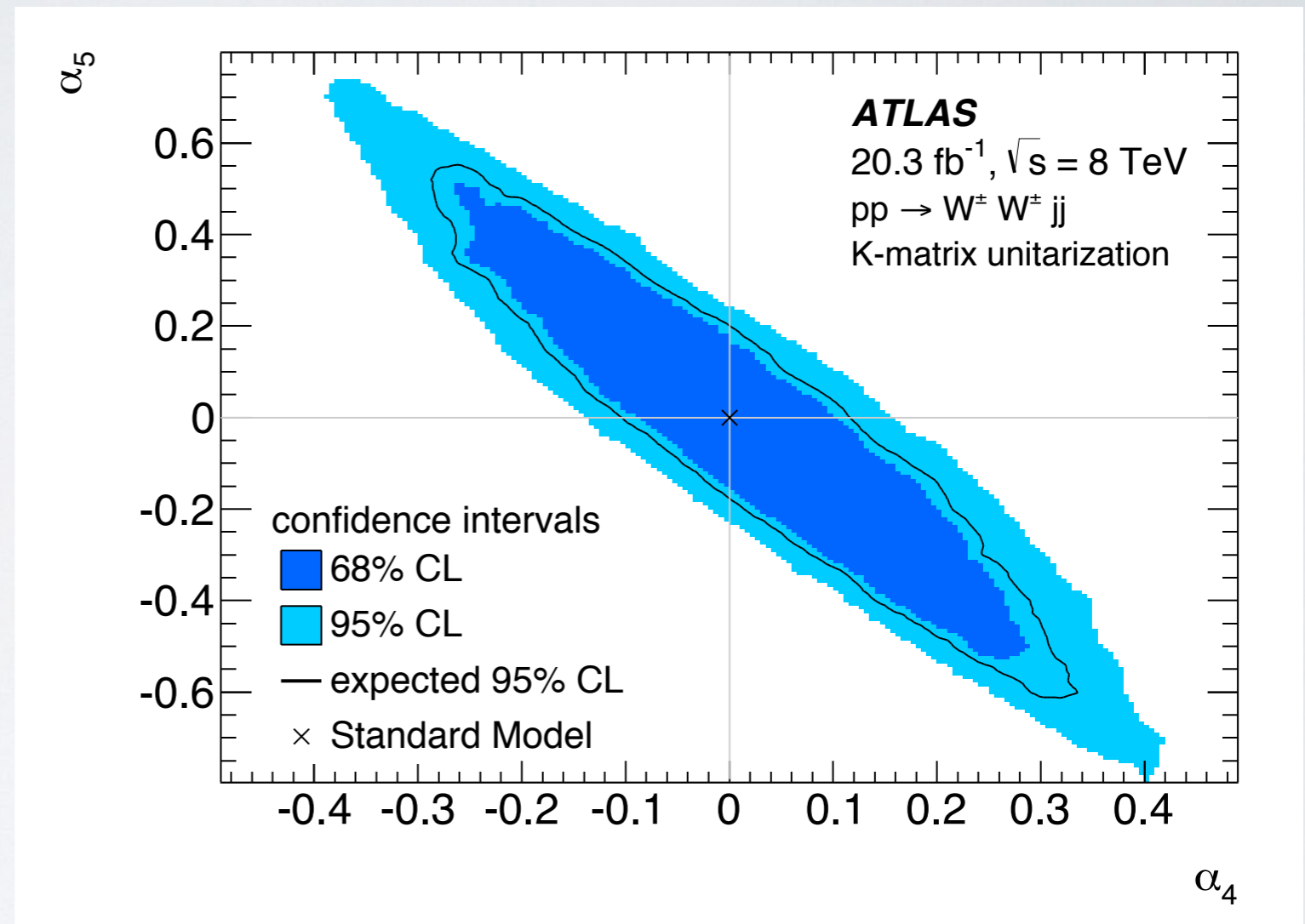
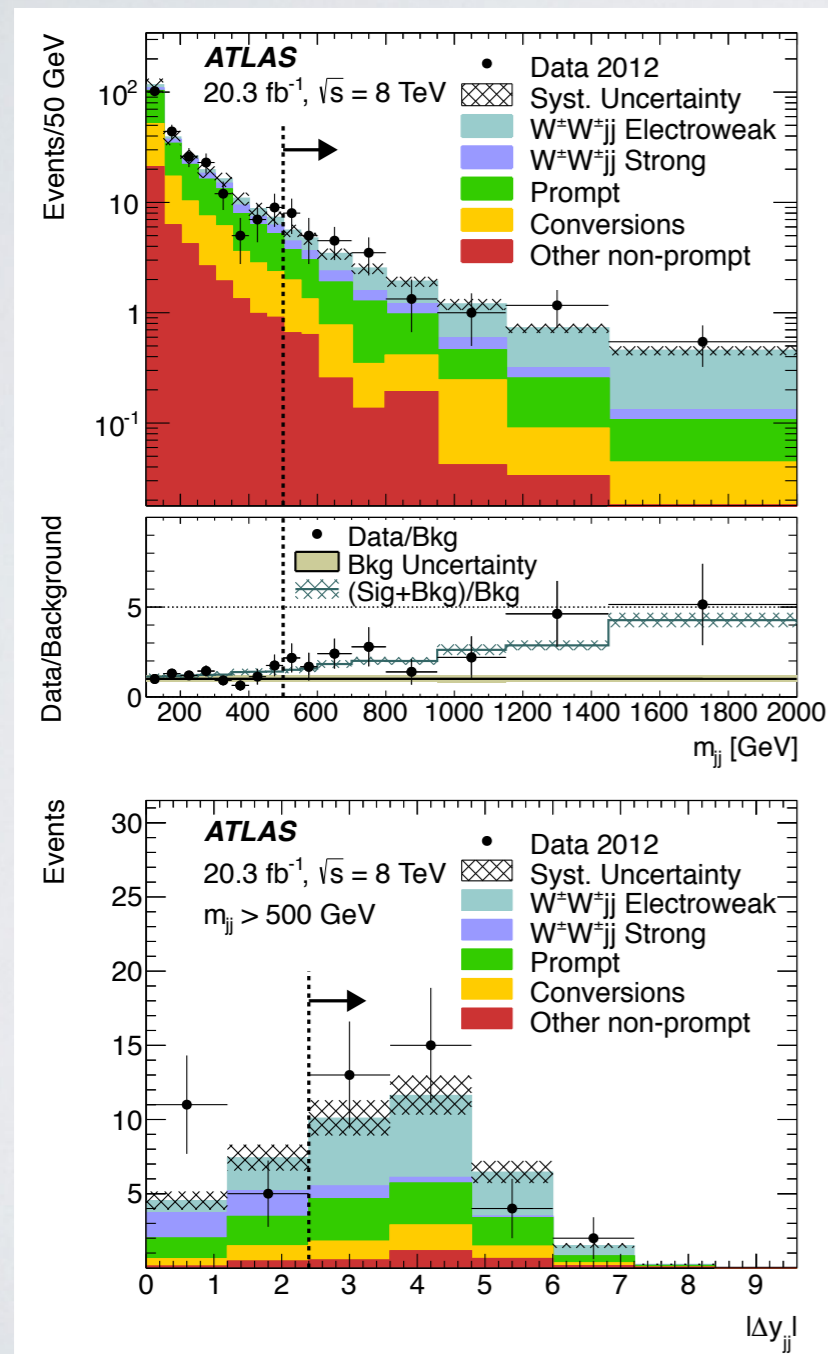


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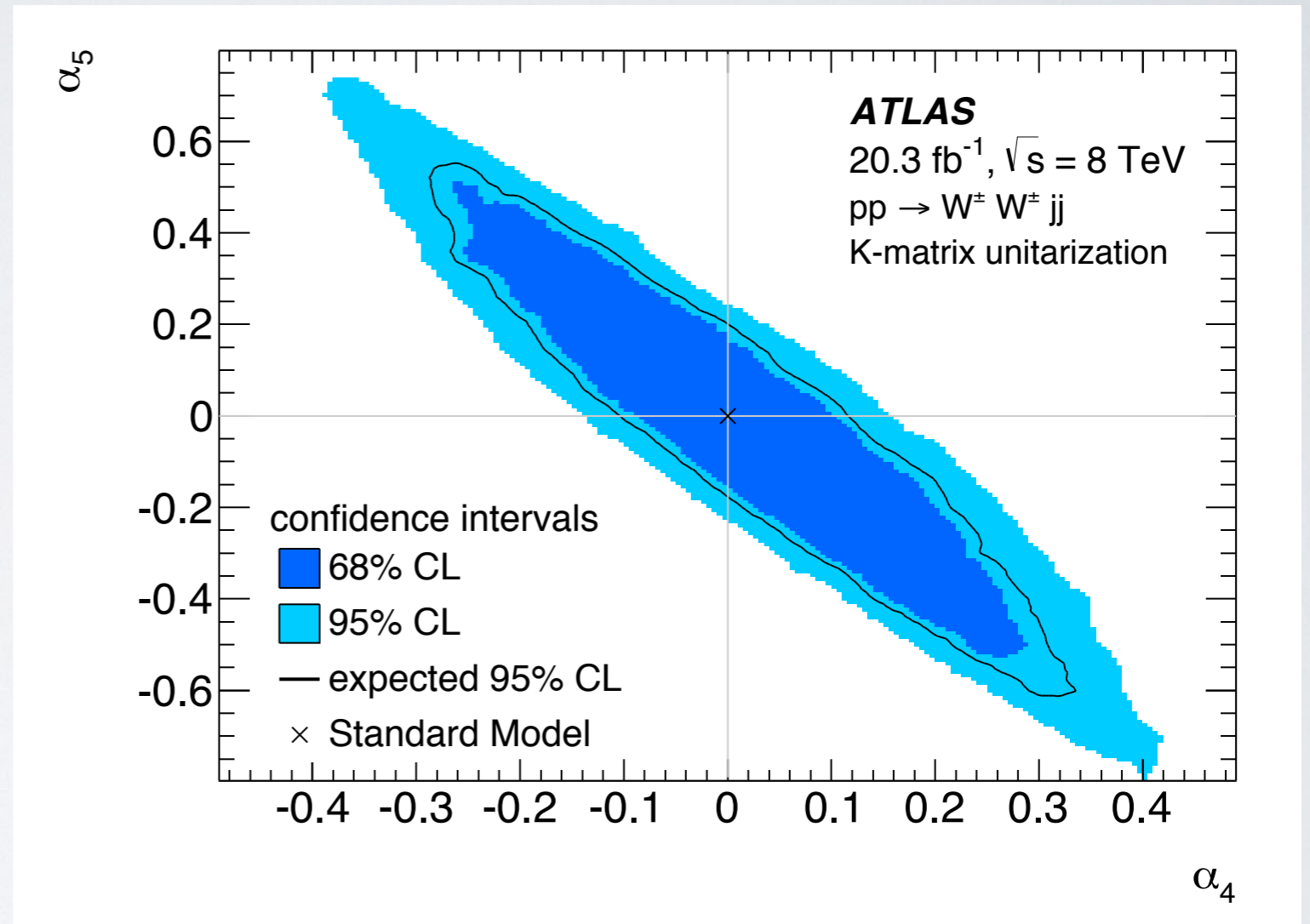
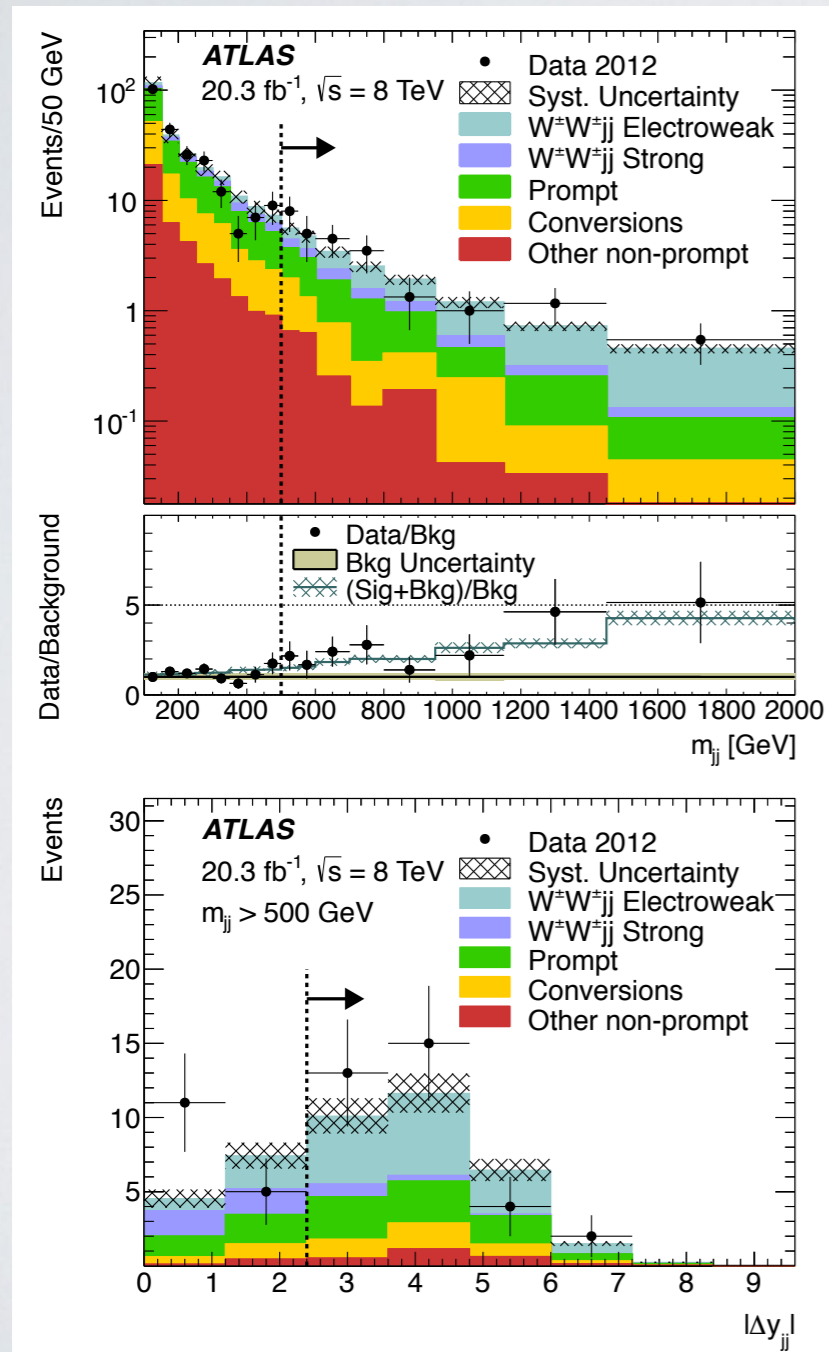
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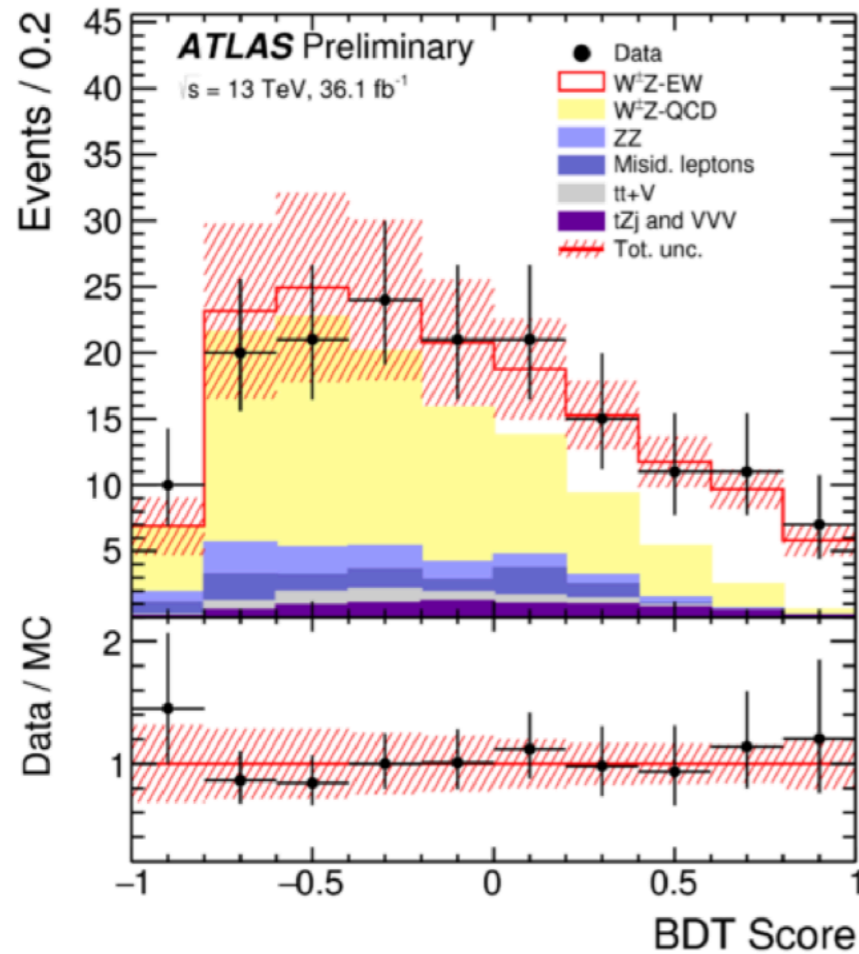
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Exploration of E-frontier → look for heavy objects, including high-mass $V_L V_L$ scattering:
 □ requires as much integrated luminosity as possible (cross-section goes like 1/s)

F. Gianotti, 01/2014





Post-fit background normalisations

$$\mu_{\text{WZ-QCD}} = 0.60 \pm 0.25$$

$$\mu_{\text{ttV}} = 1.18 \pm 0.19$$

$$\mu_{\text{ZZ}} = 1.34 \pm 0.29$$

$$pp \rightarrow WZjj \rightarrow l\nu lljj$$

ATLAS-CONF-2018-033
 M.-A. Pleier, Seoul, ICHEP 2018

WZjj-EW measured signal strength:

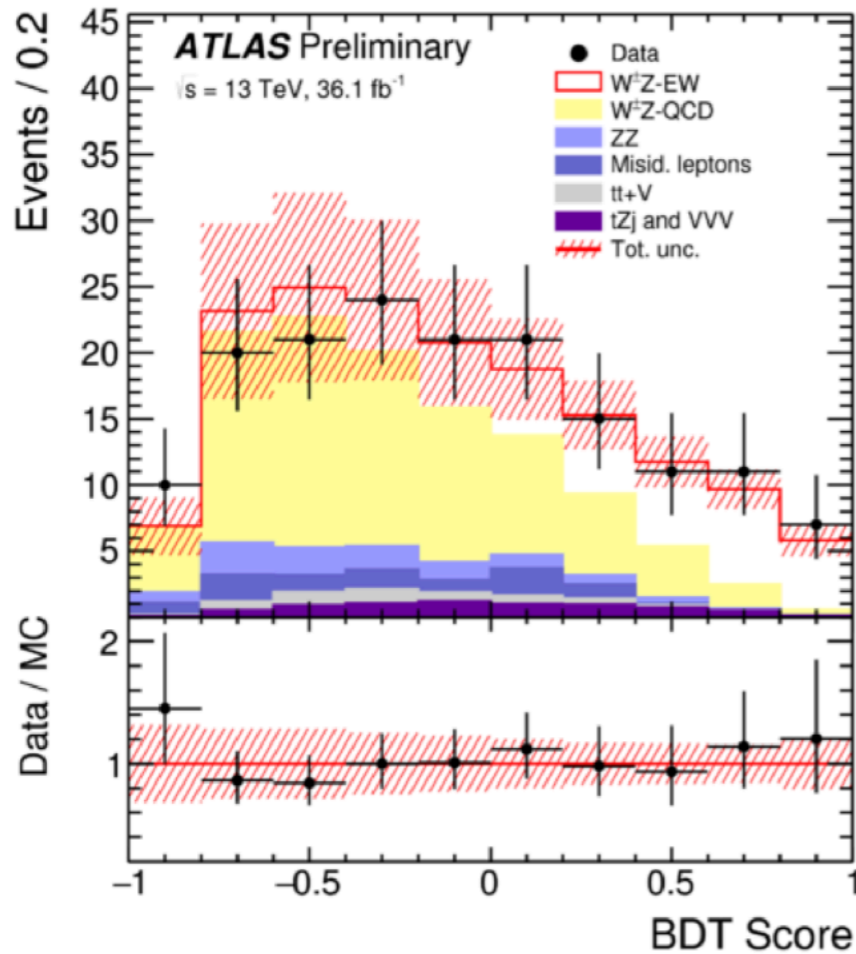
$$\mu_{\text{EW}} = 1.77 \pm 0.41(\text{stat.}) \pm 0.17(\text{syst.}) = 1.77 \pm 0.45$$

Observed sign.: 5.6σ (3.3σ expected)

Corresponding fid. cross section:

$$\begin{aligned} \sigma_{\text{WZ}^\pm jj \rightarrow l\nu lljj}^{\text{fid., EW}} &= 0.57^{+0.15}_{-0.14} \text{ fb} \\ &= 0.57^{+0.14}_{-0.13} (\text{stat.})^{+0.05}_{-0.04} (\text{syst.})^{+0.04}_{-0.03} (\text{th.}) \text{ fb} \end{aligned}$$

More channels coming up ...



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$$pp \rightarrow WZjj \rightarrow lljj + X$$

$$\sigma_{WZjj}^{\text{fid}} = 2.91^{+0.53}_{-0.49} (\text{stat})^{+0.41}_{-0.34} (\text{syst})$$

CMS-SMP-18-001 Observed (expected) of EW WZ 1.9σ (2.7σ)

$$pp \rightarrow W^+W^+jj \rightarrow l\nu l\nu jj$$

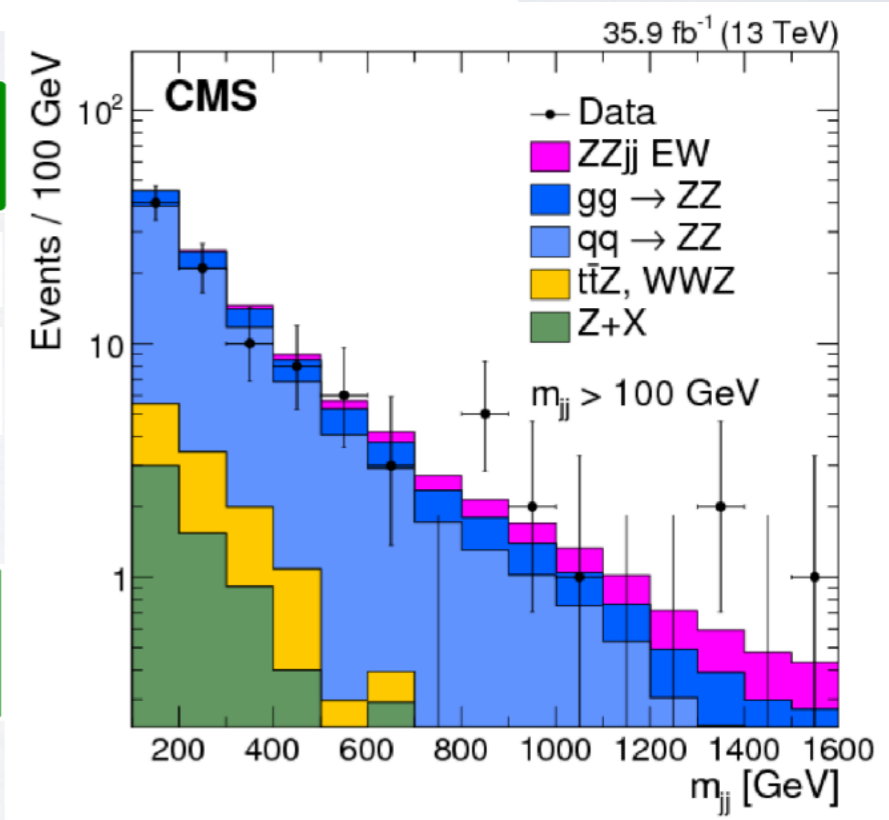
$$\sigma_{\text{fid}} = 3.83 \pm 0.66 (\text{stat}) \pm 0.35 (\text{syst}) \text{ fb}$$

PRL 120, 081801 (2018) Observed (expected) of 5.5σ (5.7σ)

$$pp \rightarrow ZZjj \rightarrow ll lljj$$

$$\mu = \sigma_{\text{obs}}/\sigma_{\text{th.}} = 1.39^{+0.72}_{-0.57} (\text{stat})^{+0.46}_{-0.31} (\text{syst.})$$

PLB 774(2017) 682 Observed (expected) of 2.7σ (1.6σ)



Motivated by SMEFT:

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_i \left[\frac{c_i^{(5)}}{\Lambda} \mathcal{O}_i^{(5)} + \frac{c_i^{(6)}}{\Lambda^2} \mathcal{O}_i^{(6)} + \frac{c_i^{(7)}}{\Lambda^3} \mathcal{O}_i^{(7)} + \frac{c_i^{(8)}}{\Lambda^4} \mathcal{O}_i^{(8)} + \dots \right]$$

Longitudinal operators

$$\mathcal{L}_{S,0} = F_{S,0} \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger (\mathbf{D}_\nu \mathbf{H}) \right] \text{tr} \left[(\mathbf{D}^\mu \mathbf{H})^\dagger (\mathbf{D}^\nu \mathbf{H}) \right]$$

$$\mathcal{L}_{S,1} = F_{S,1} \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger (\mathbf{D}^\mu \mathbf{H}) \right] \text{tr} \left[(\mathbf{D}_\nu \mathbf{H})^\dagger (\mathbf{D}^\nu \mathbf{H}) \right]$$

Mixed operators

$$\mathcal{L}_{M,0} = -g^2 F_{M_0} \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger (\mathbf{D}^\mu \mathbf{H}) \right] \text{tr} \left[\mathbf{W}_{\nu\rho} \mathbf{W}^{\nu\rho} \right]$$

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$$\mathcal{L}_{M,7} = -g^2 F_{M_7} \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger \mathbf{W}_{\nu\rho} \mathbf{W}^{\nu\mu} (\mathbf{D}^\rho \mathbf{H}) \right];$$

S. Weinberg, 1979

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Transversal operators

$$\mathcal{L}_{T,0} = g^4 F_{T_0} \text{tr} \left[\mathbf{W}_{\mu\nu} \mathbf{W}^{\mu\nu} \right] \text{tr} \left[\mathbf{W}_{\alpha\beta} \mathbf{W}^{\alpha\beta} \right],$$

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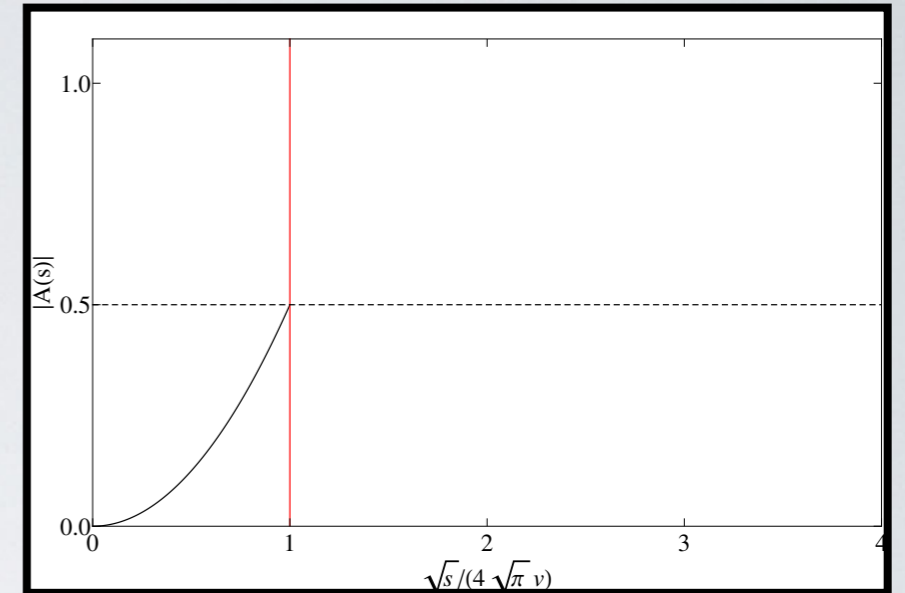
Energy rise of operators lead to unitarity violation

Unitarity violation between operators in UV-complete Theory

Procedures to treat unitarity violations

Cut-off (a.k.a. “Event clipping”) $\theta(\Lambda_C^2 - s)$

unitarity bound (0th partial wave) at Λ_C
no continuous transition beyond
Effect on BDT training not clear



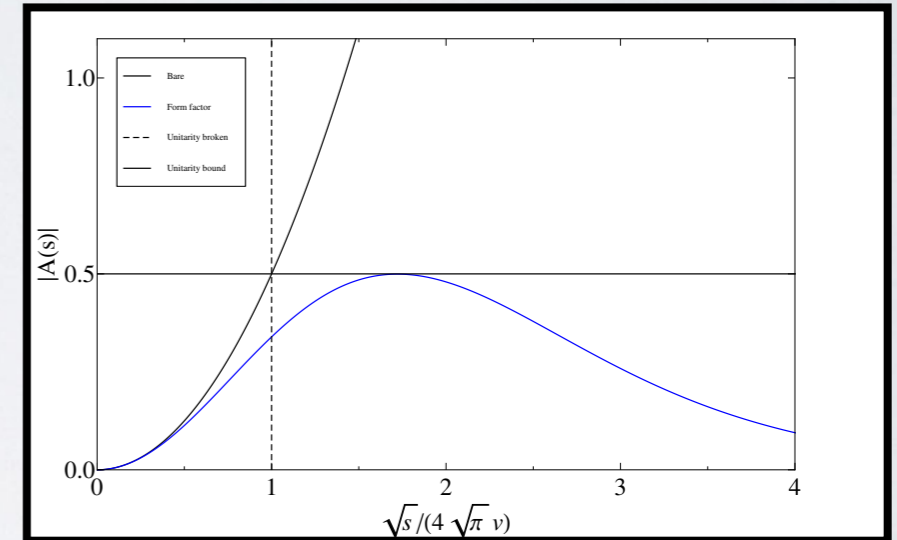
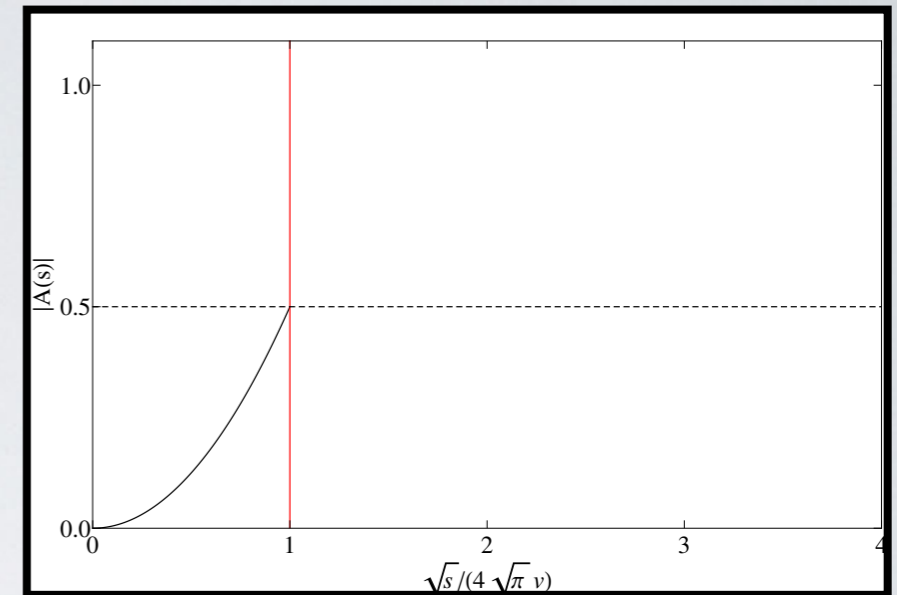
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Form factor

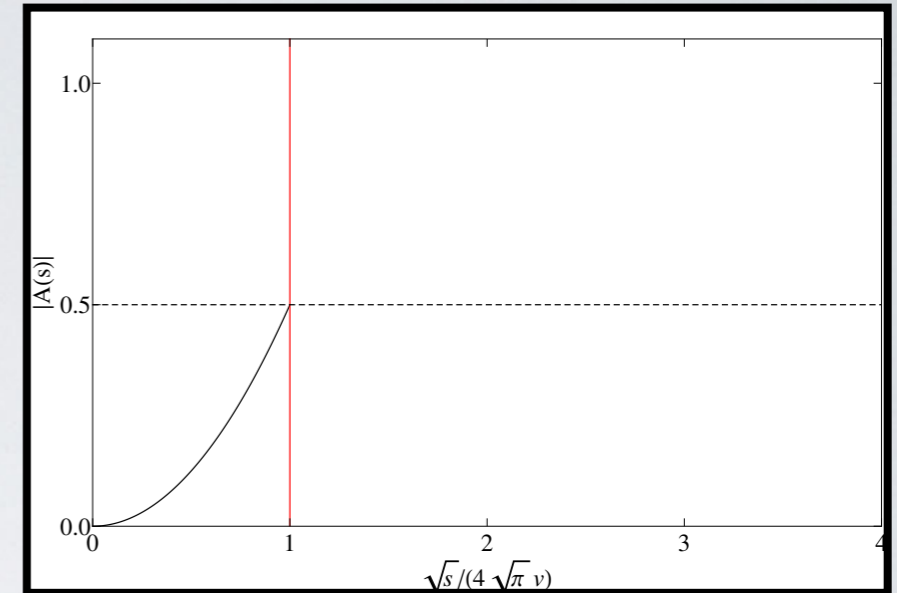
$$\frac{1}{\left(1 + \frac{s}{\Lambda_{FF}^2}\right)^n}$$

Applicable for arbitrary operators, tuning in 2 parameters: n damps unitarity violation, Λ_{FF} highest value to satisfy 0th partial wave



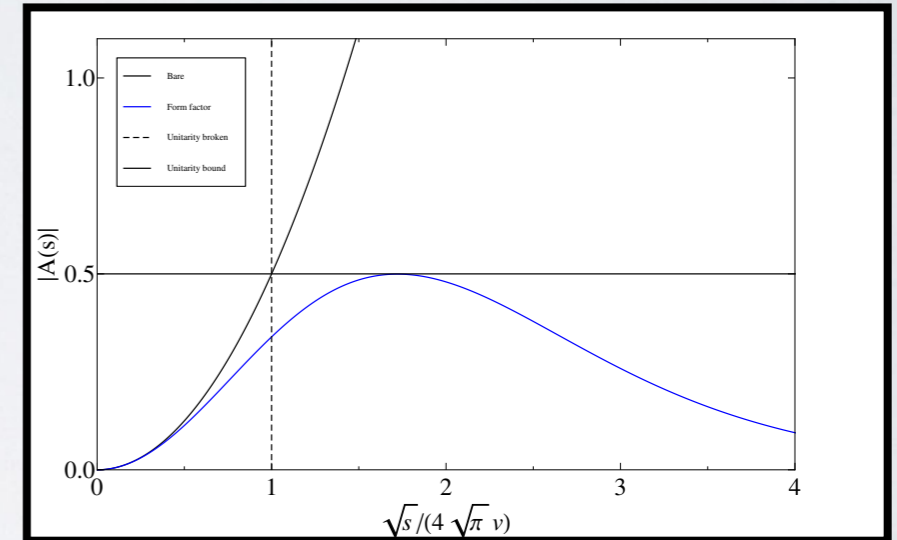
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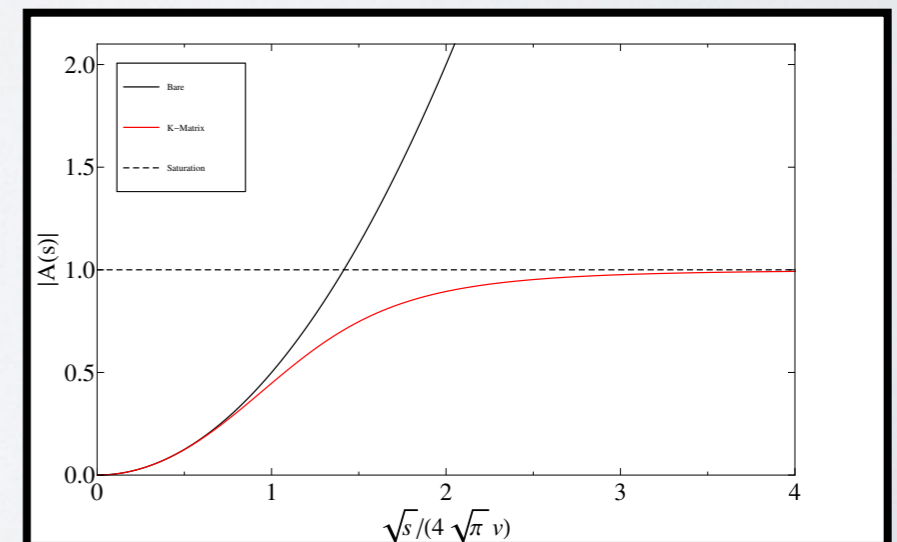
Form factor $\frac{1}{\left(1 + \frac{s}{\Lambda_{FF}^2}\right)^n}$

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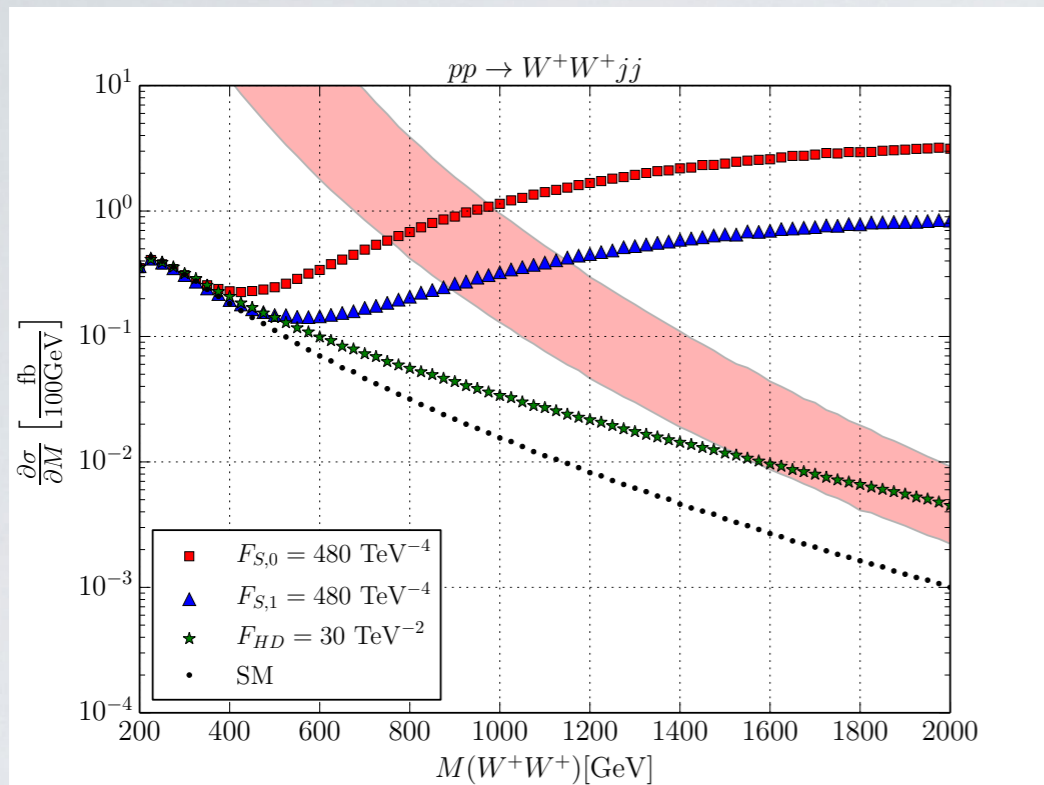
K-/T-matrix saturation $a = \frac{1}{\text{Re}\left(\frac{1}{a_0}\right) - i}$

saturates amplitude [projection to unitarity circle], also for complex ampl., **no additional parameters**

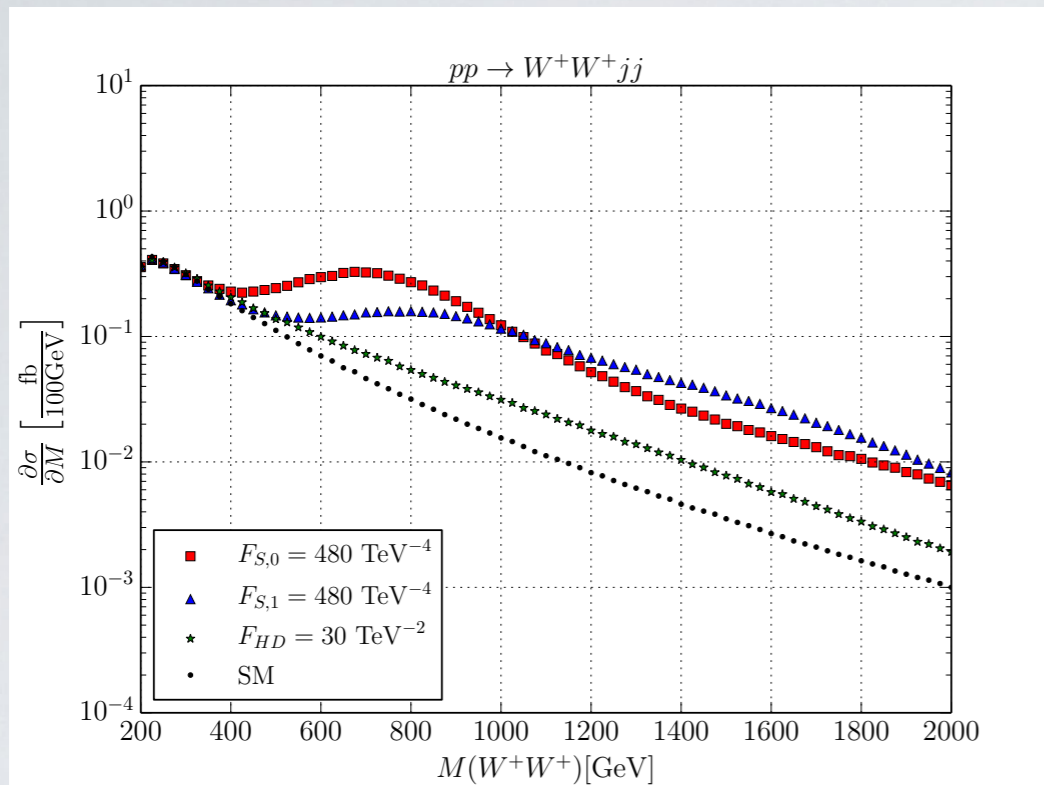


Alboteanu/Kilian/JRR, 2008

Kilian/Ohl/JRR/Sekulla, 2014

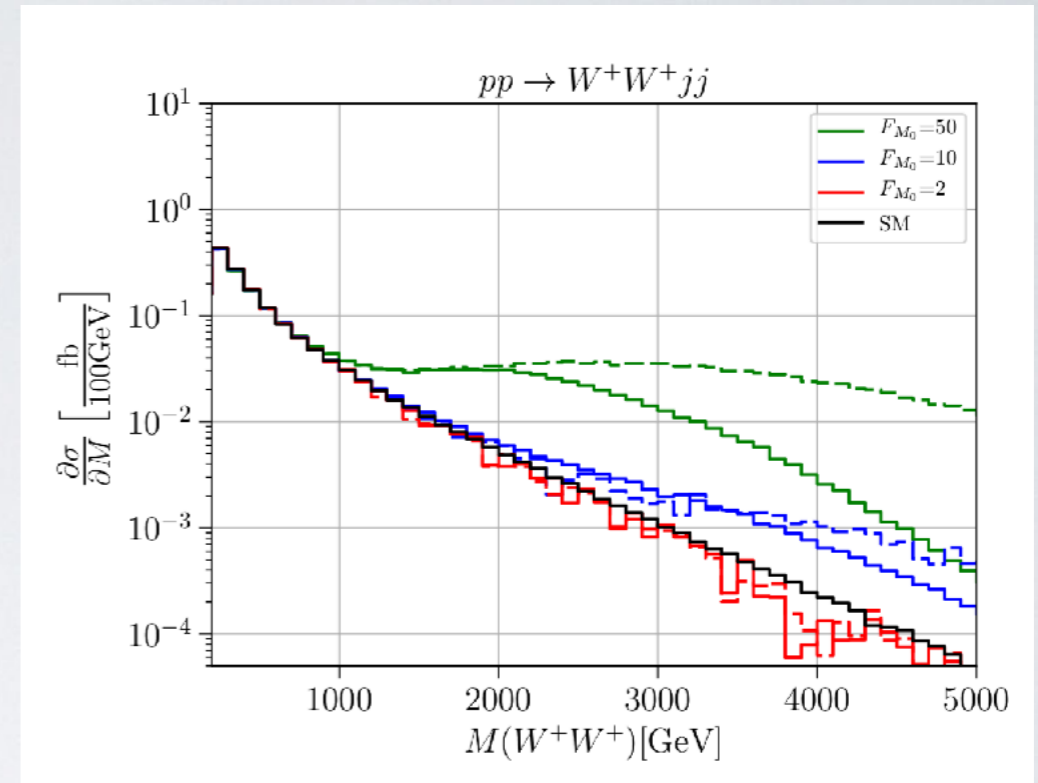
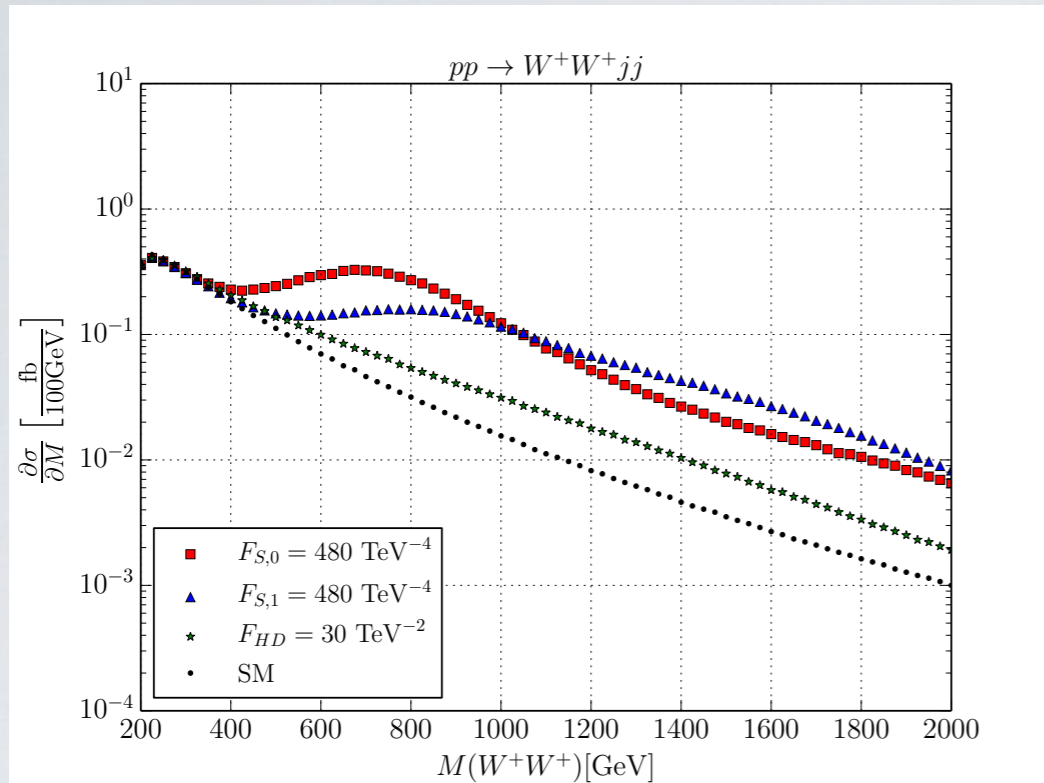


General cuts: $M_{jj} > 500 \text{ GeV}$; $\Delta\eta_{jj} > 2.4$; $p_T^j > 20 \text{ GeV}$; $|\Delta\eta_j| < 4.5$



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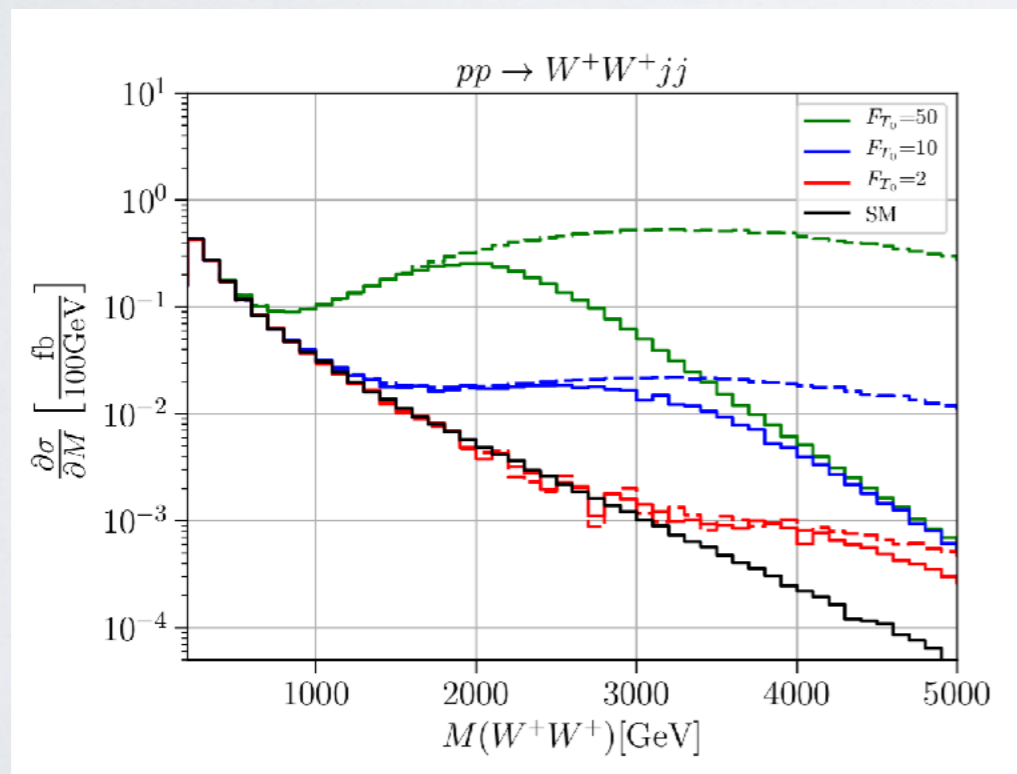
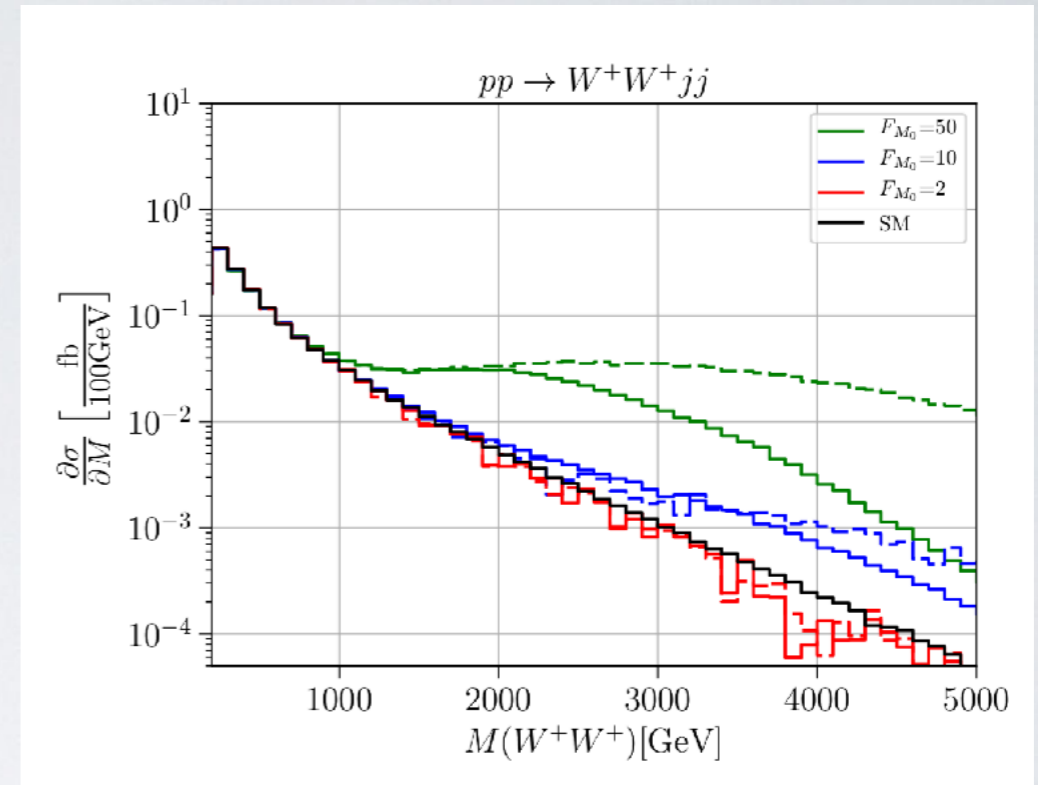
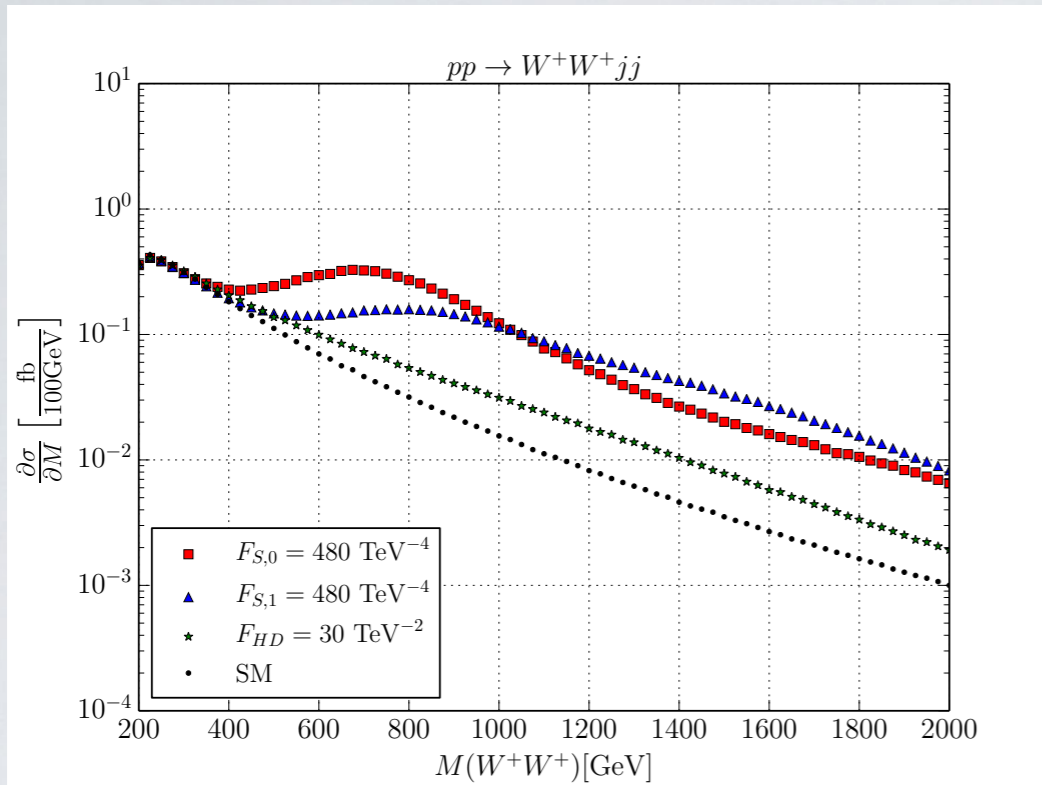
VBS diboson spectra



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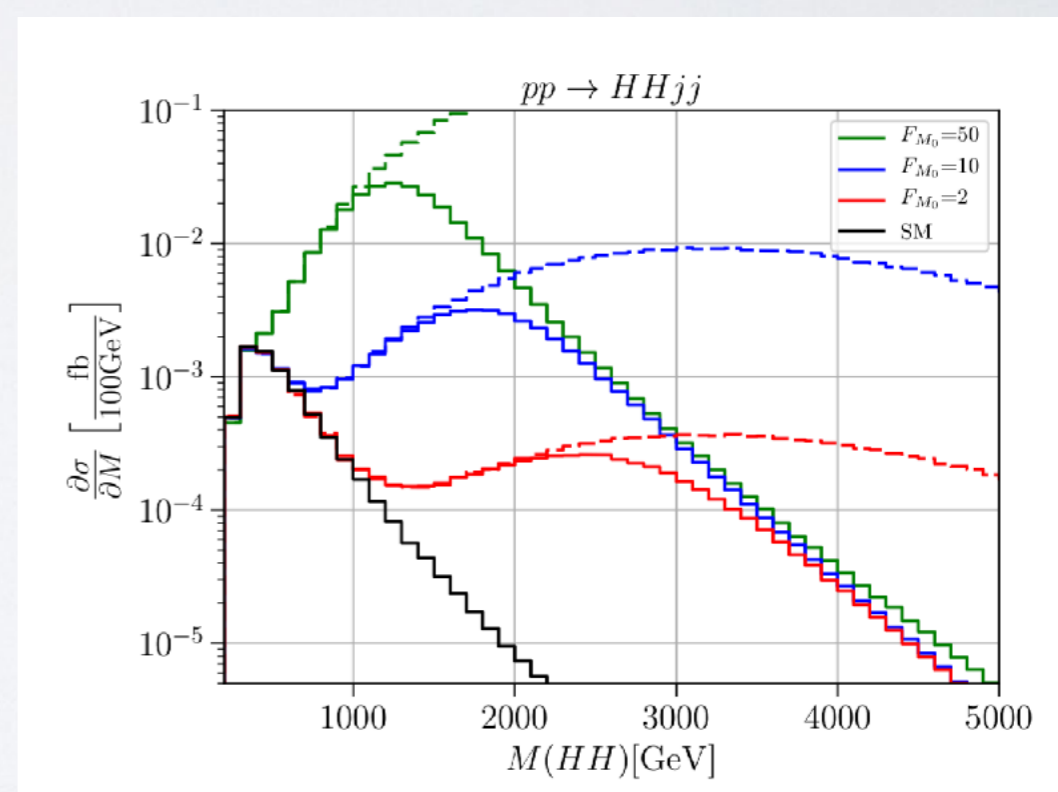
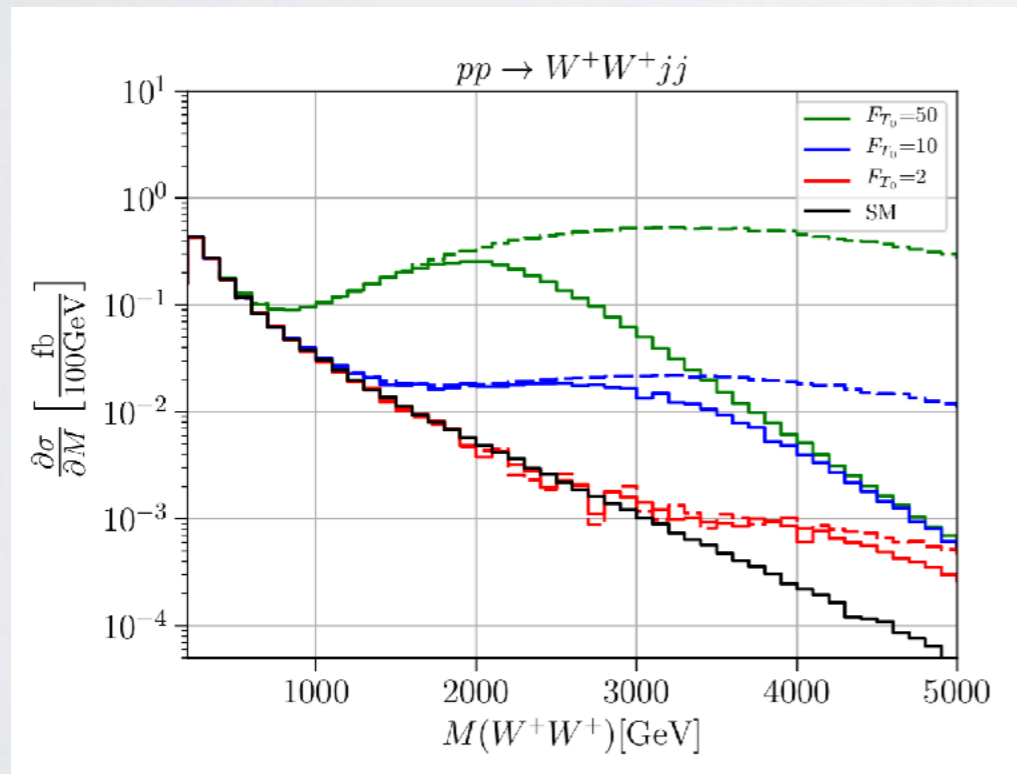
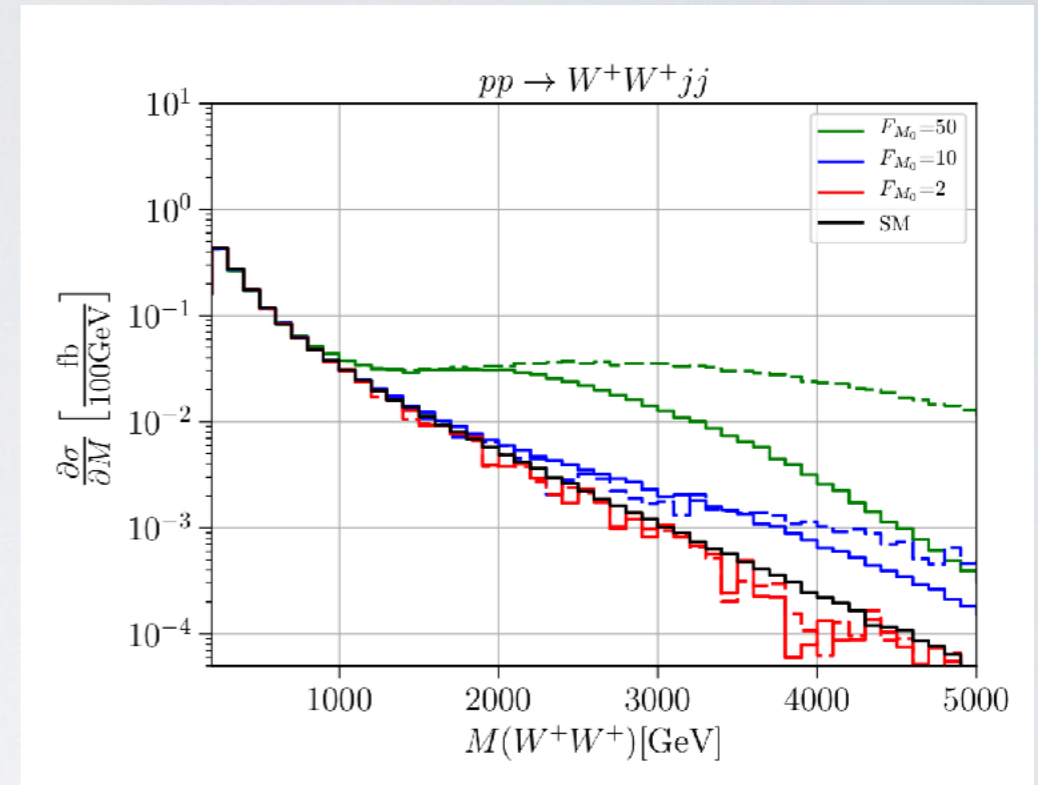
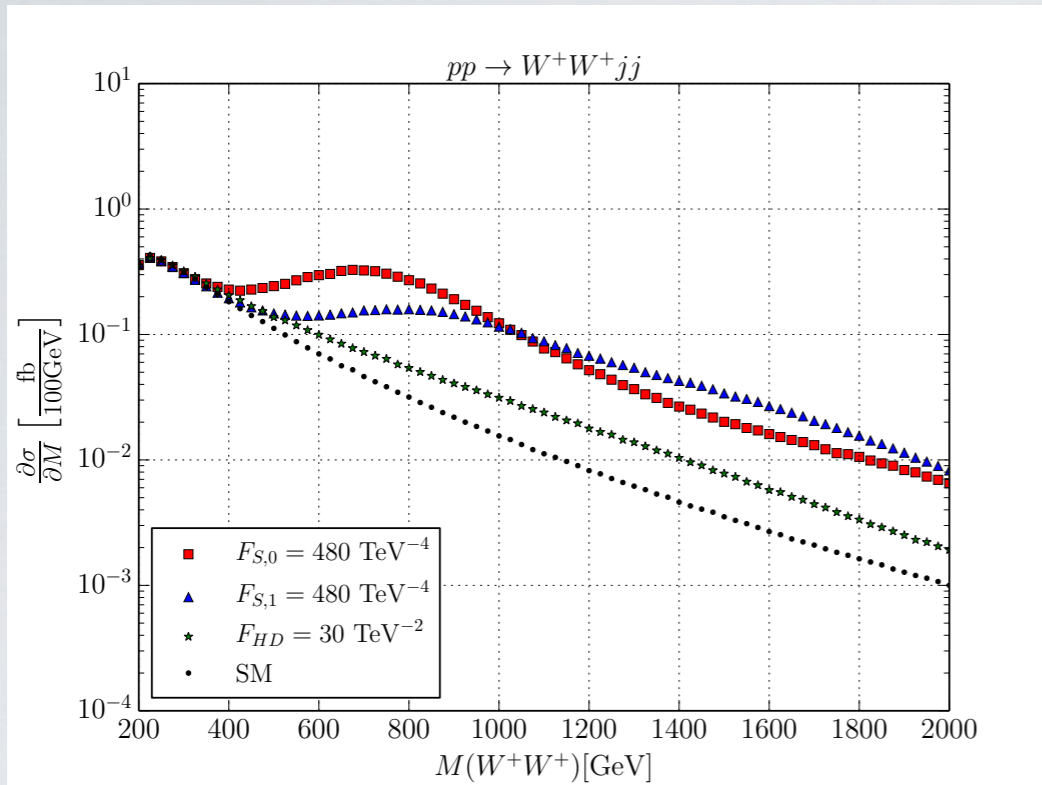
VBS diboson spectra



General cuts: $M_{jj} > 500 \text{ GeV}$; $\Delta\eta_{jj} > 2.4$; $p_T^j > 20 \text{ GeV}$; $|\Delta\eta_j| < 4.5$



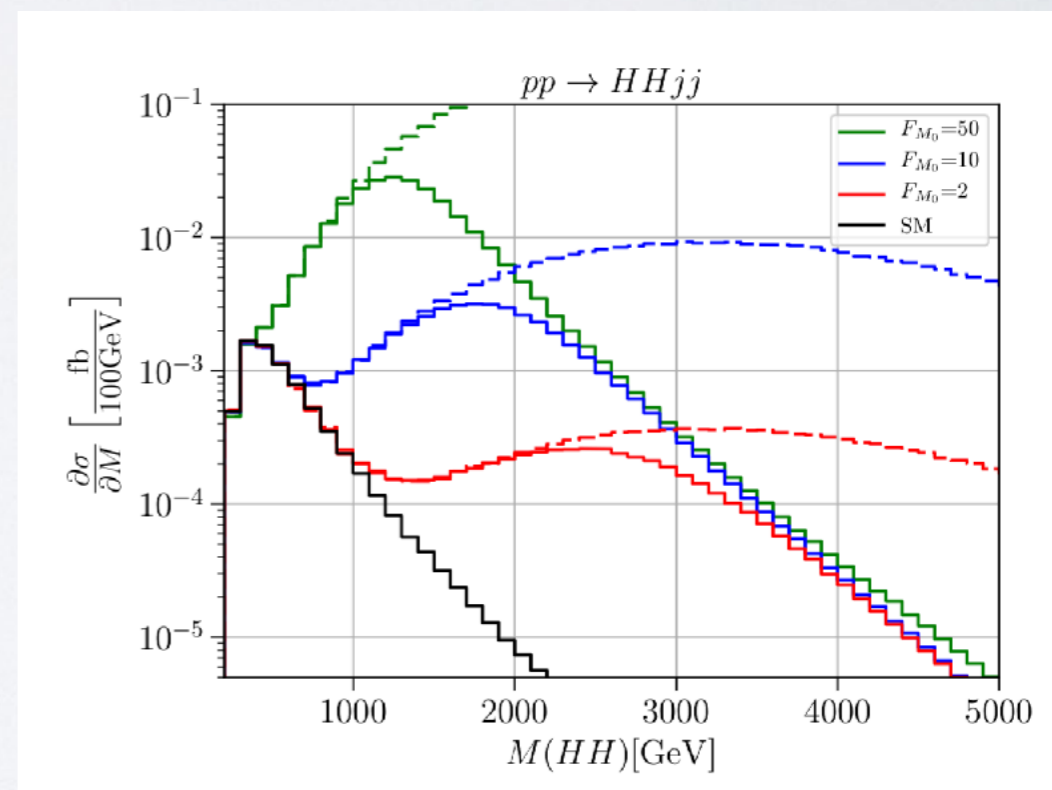
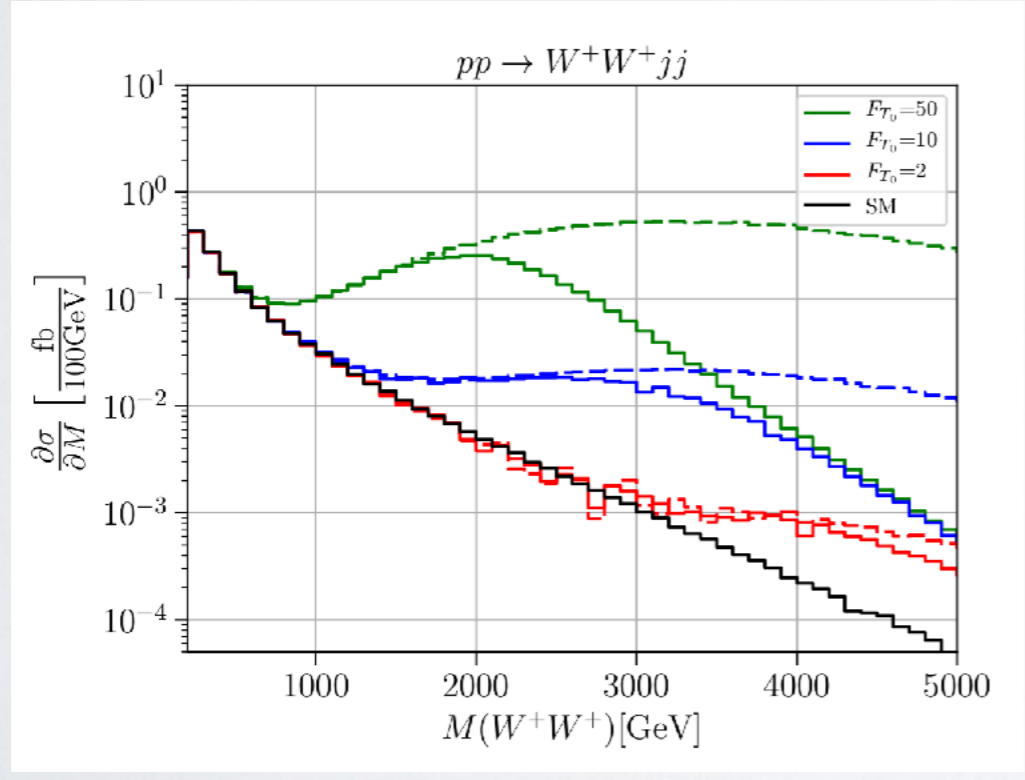
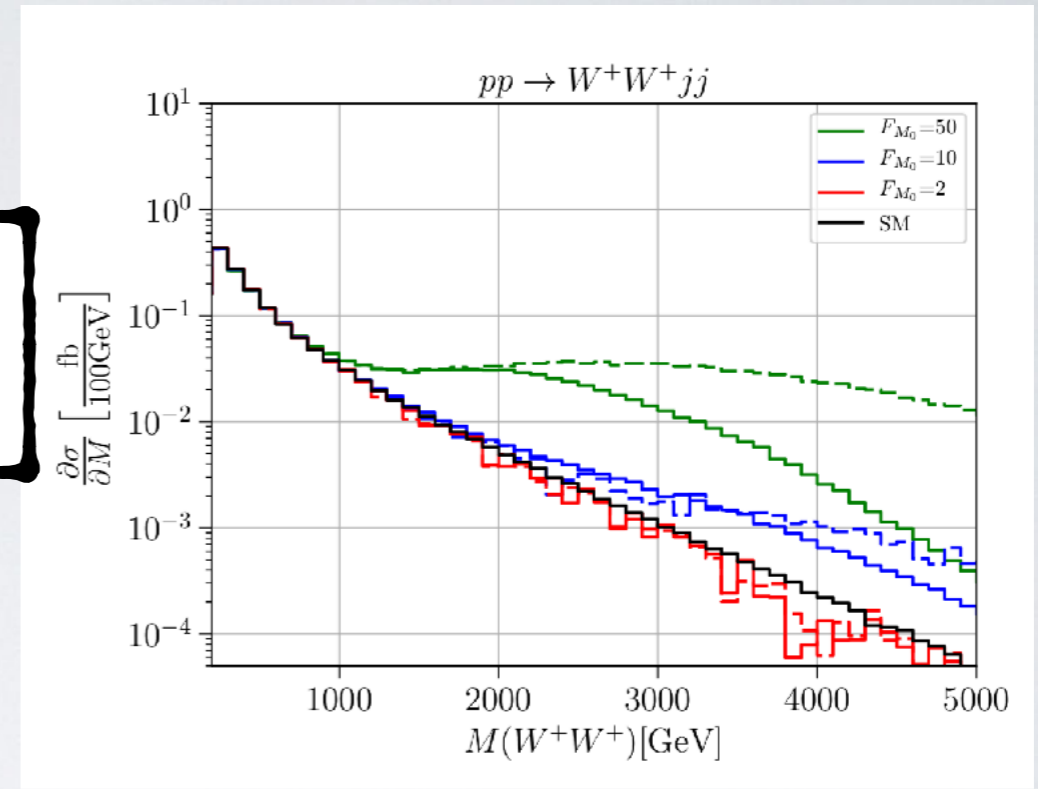
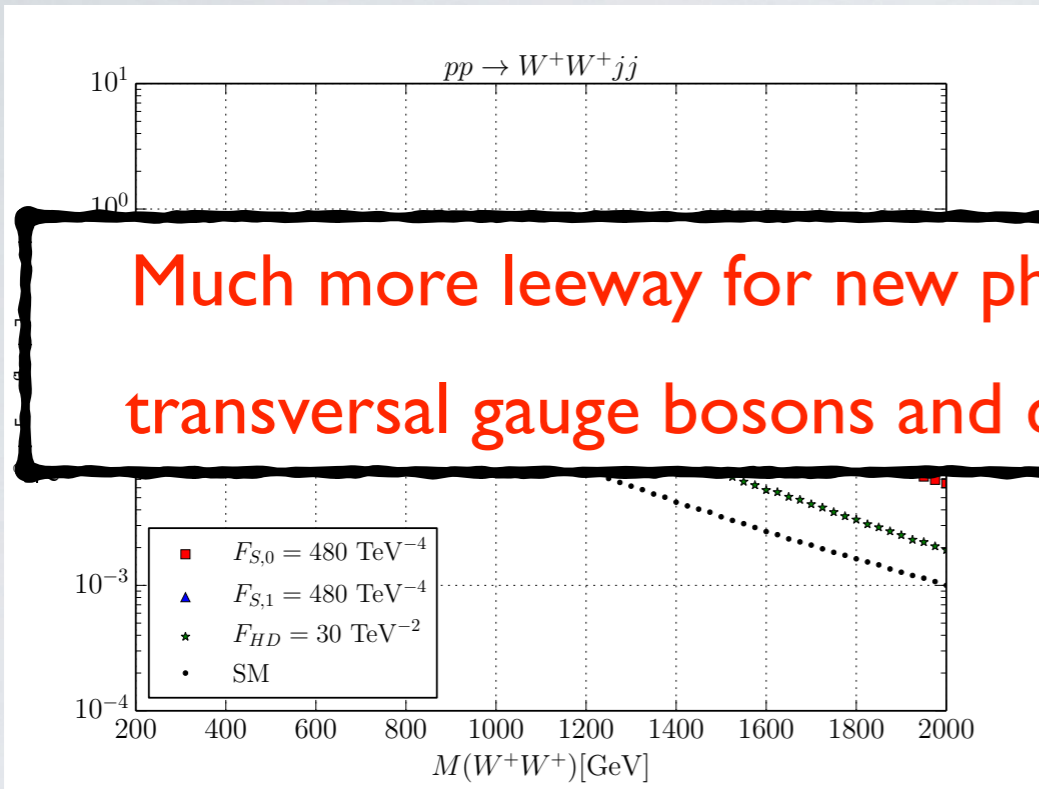
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Differential spectra in VBS

$$pp \rightarrow e^+ \mu^+ \nu_e \nu_\mu jj \quad \sqrt{s} = 14 \text{ TeV} \quad \mathcal{L} = 1 \text{ ab}^{-1}$$

Simulations with WHIZARD [<http://whizard.hepforge.org>, Kilian/Ohl/JRR]

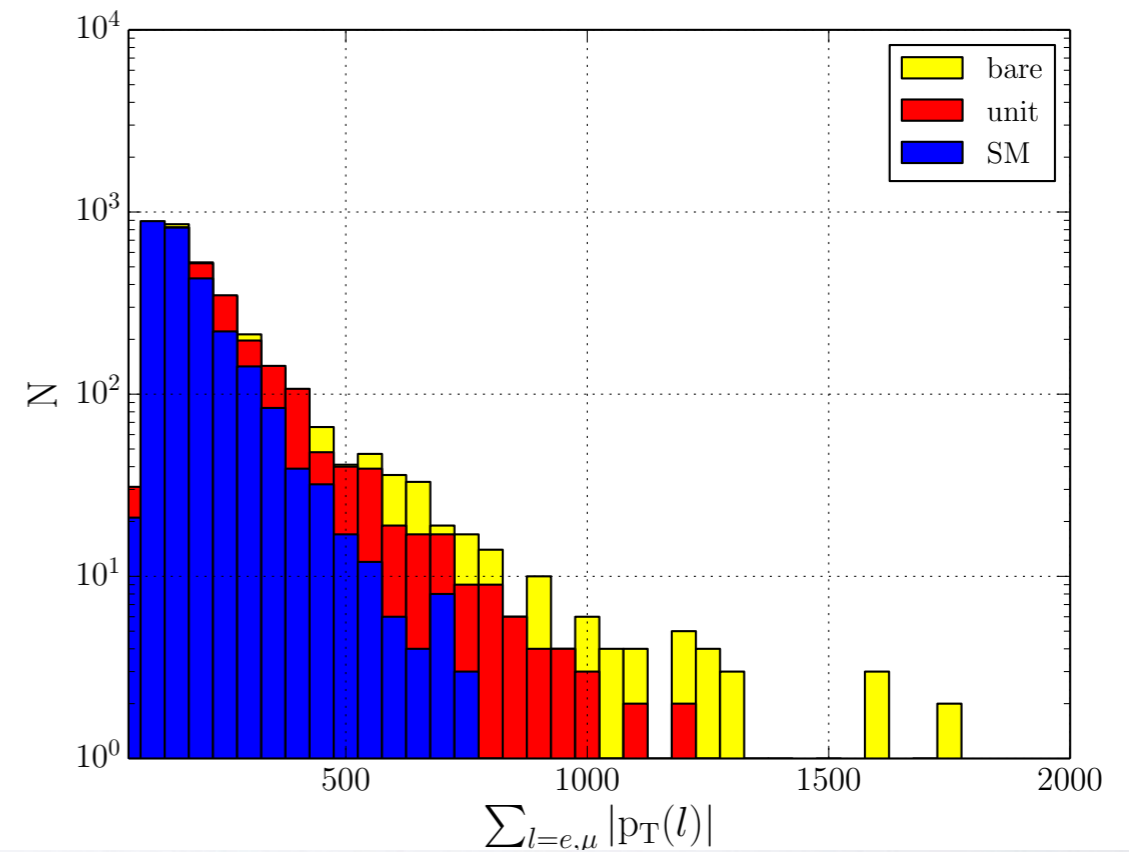
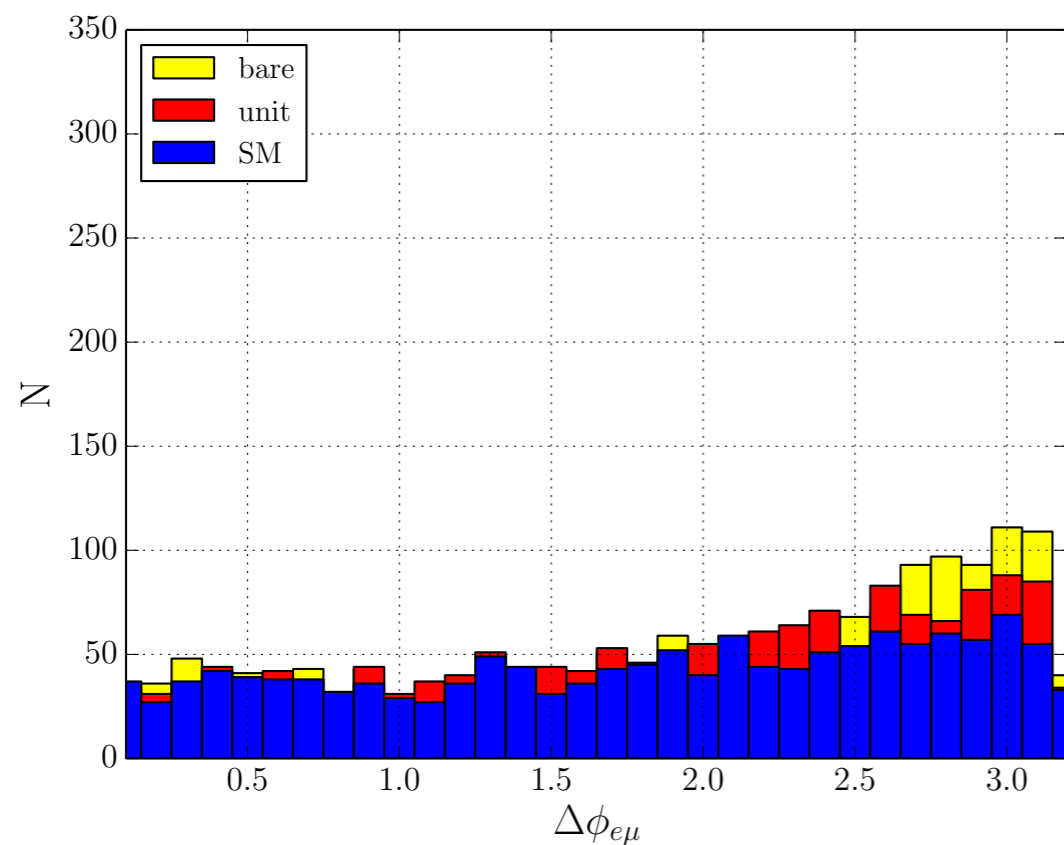


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$$\mathcal{L}_{HD} = F_{HD} \text{tr} \left[\mathbf{H}^\dagger \mathbf{H} - \frac{v^2}{4} \right] \cdot \text{tr} \left[(\mathbf{D}_\mu \mathbf{H})^\dagger \mathbf{D}_\mu \mathbf{H} \right] \quad F_{HD} = 30 \text{ TeV}^{-2}$$



Differential spectra in VBS

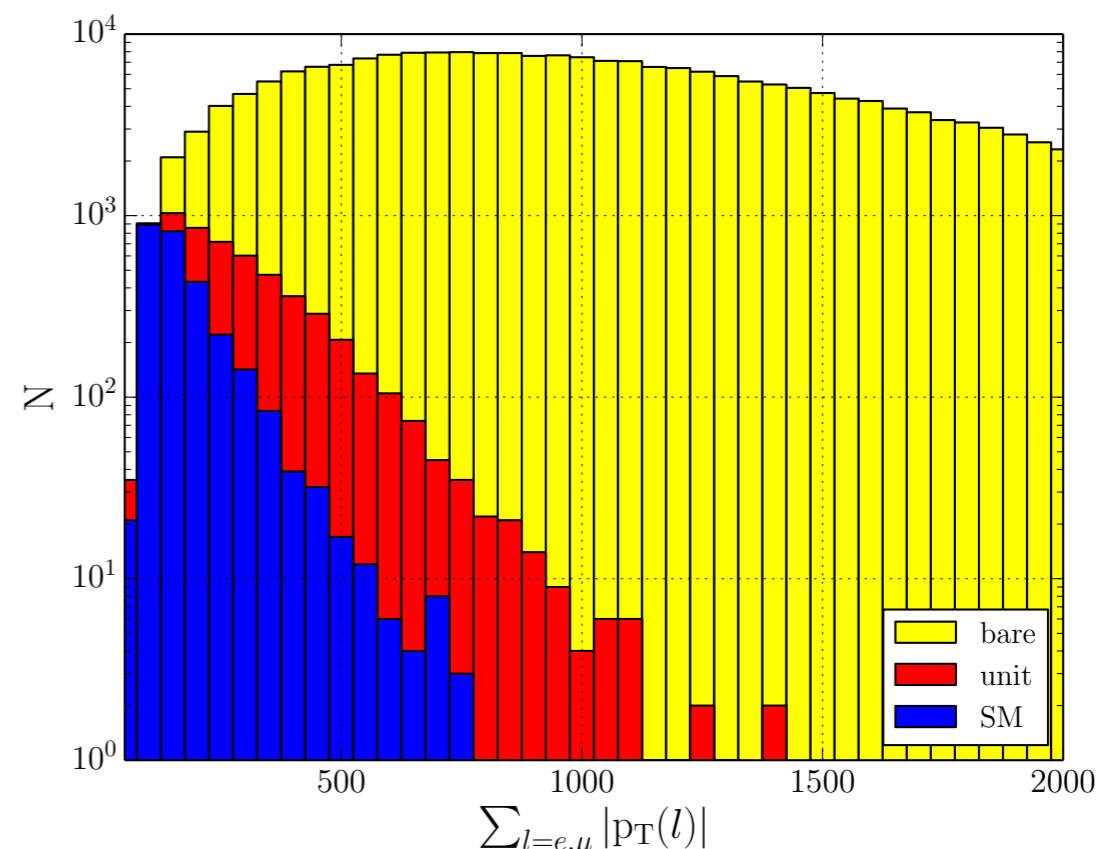
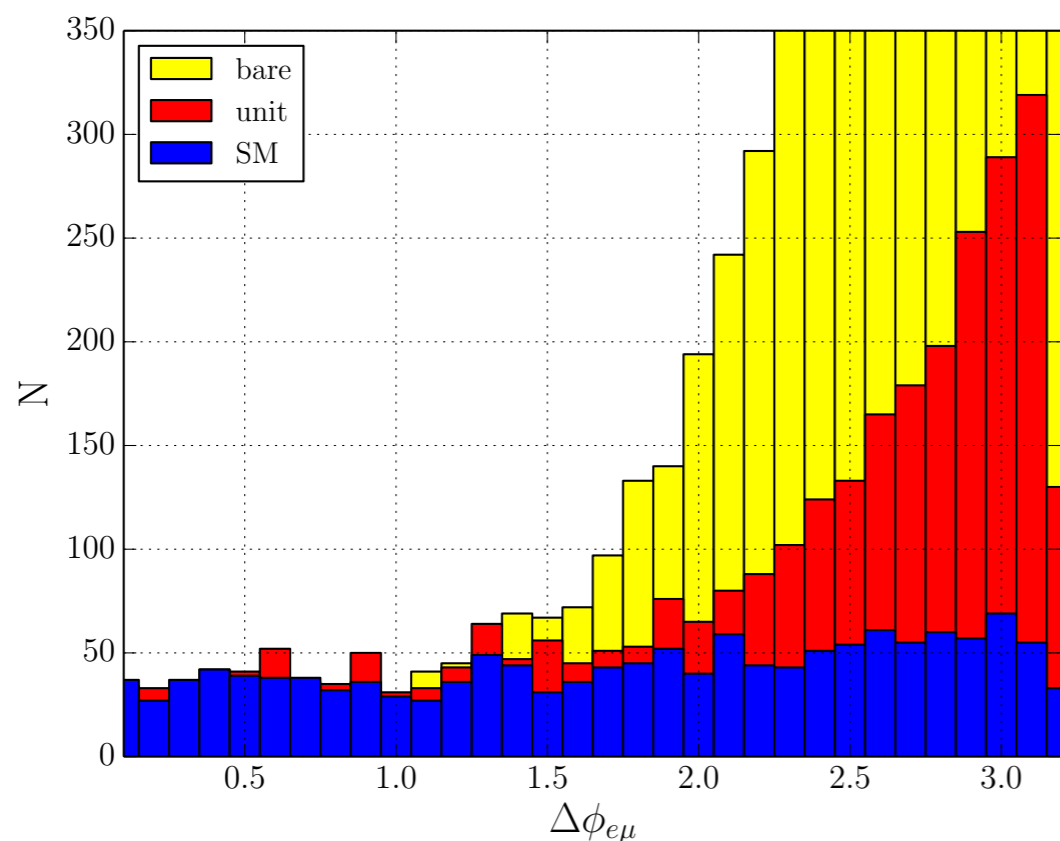
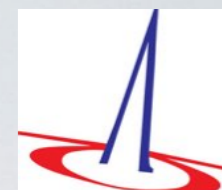
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$$F_{S,0} = 480 \text{ TeV}^{-4}$$



$$M_{jj} > 500 \text{ GeV}; \quad \Delta\eta_{jj} > 2.4; \quad p_T^j > 20 \text{ GeV}; \quad |\Delta\eta_j| < 4.5; \quad p_T^\ell > 20 \text{ GeV}$$

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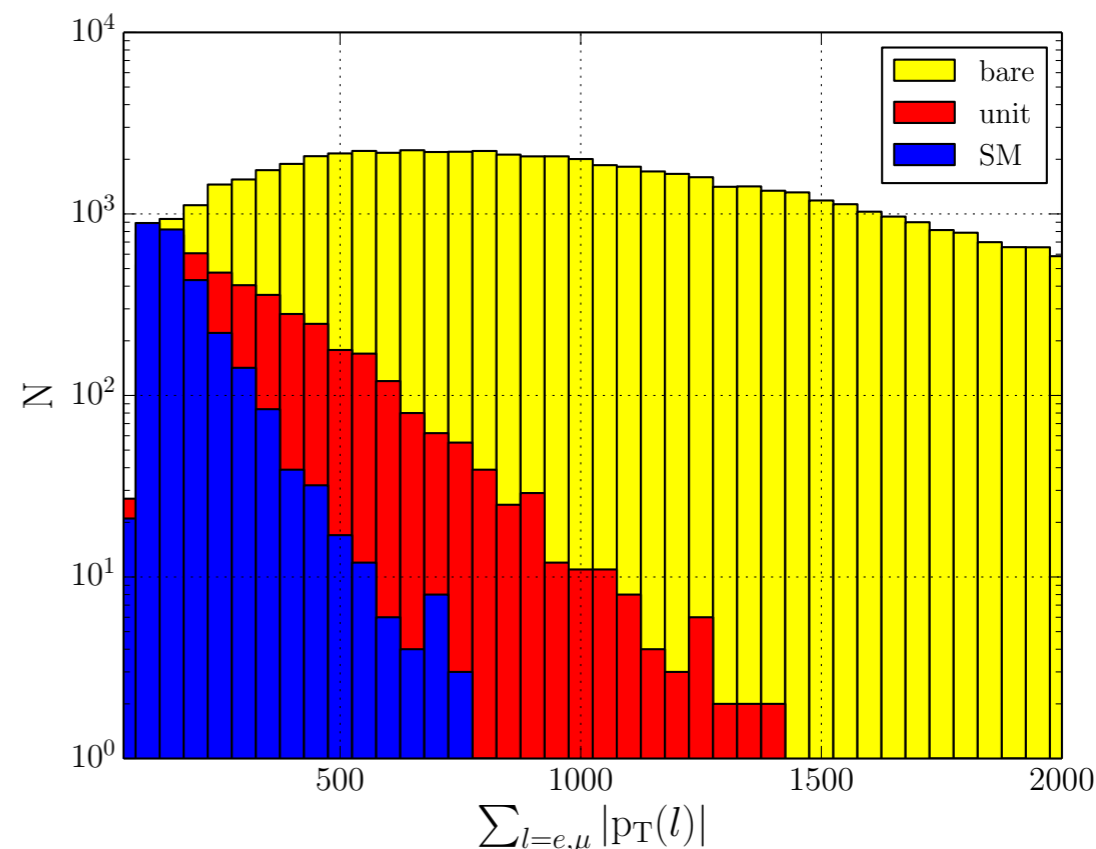
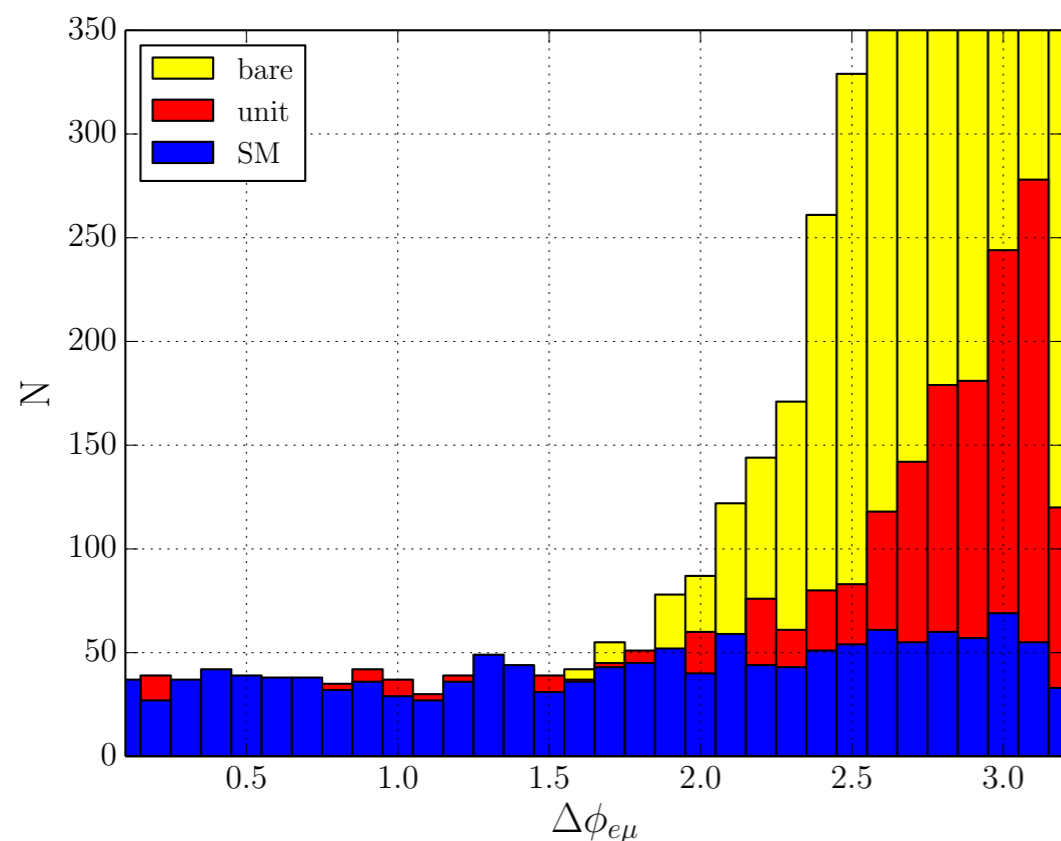
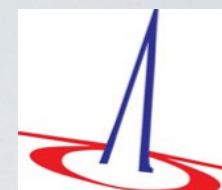
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Resonances: Quantum numbers & simplified models ^{11 / 16}

- Rise of amplitude: is Taylor expansion below a resonance
- Resonances might be in direct reach of LHC
- EFT framework EW-restored regime: $SU(2)_L \times SU(2)_R, SU(2)_L \times U(1)_Y$ gauged
- Include EFT operators in addition (more resonances, continuum contribution)
- Apply T -matrix unitarization beyond resonance (“UV-incomplete” model)

Spins 0, 2 considered, Spin 1 has different physics (mixing with W/Z)

	isoscalar	isotensor
scalar	σ^0	$\phi_t^{--}, \phi_t^-, \phi_t^0, \phi_t^+, \phi_t^{++}$ $\phi_v^-, \phi_v^0, \phi_v^+$ ϕ_s^0
tensor	f^0	$(X_t^{--}, X_t^-, X_t^0, X_t^+, X_t^{++})$ X_v^-, X_v^0, X_v^+ X_s^0
...

$$32\pi\Gamma/M^5$$

	σ	ϕ	f	X
$F_{S,0}$	$\frac{1}{2}$	2	15	5
$F_{S,1}$	–	$-\frac{1}{2}$	-5	-35

Translation into Wilson coefficients below resonance

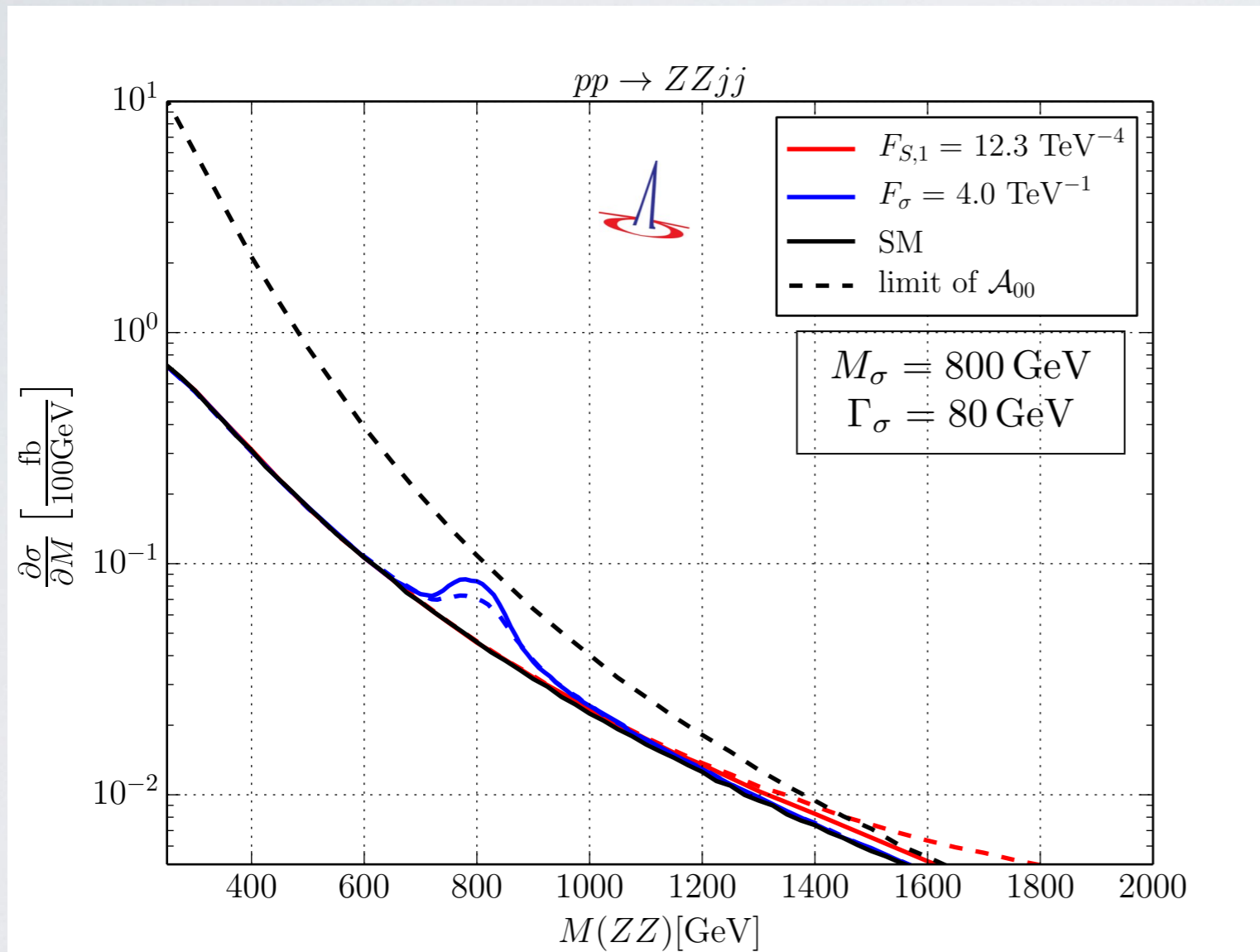
Comparison: Simplified Models & EFT

Kilian/Ohl/JRR/Sekulla: PRD93(16),3. 036004 [1511.00022]

Brass/Fleper/Kilian/JRR/Sekulla: w. EPJC [1807.02512]

Black dashed line:

saturation of $\mathcal{A}_{22}(W^+W^+)/\mathcal{A}_{00}(ZZ)$



- EFT fails at resonance
- aQGC describe rise of resonance
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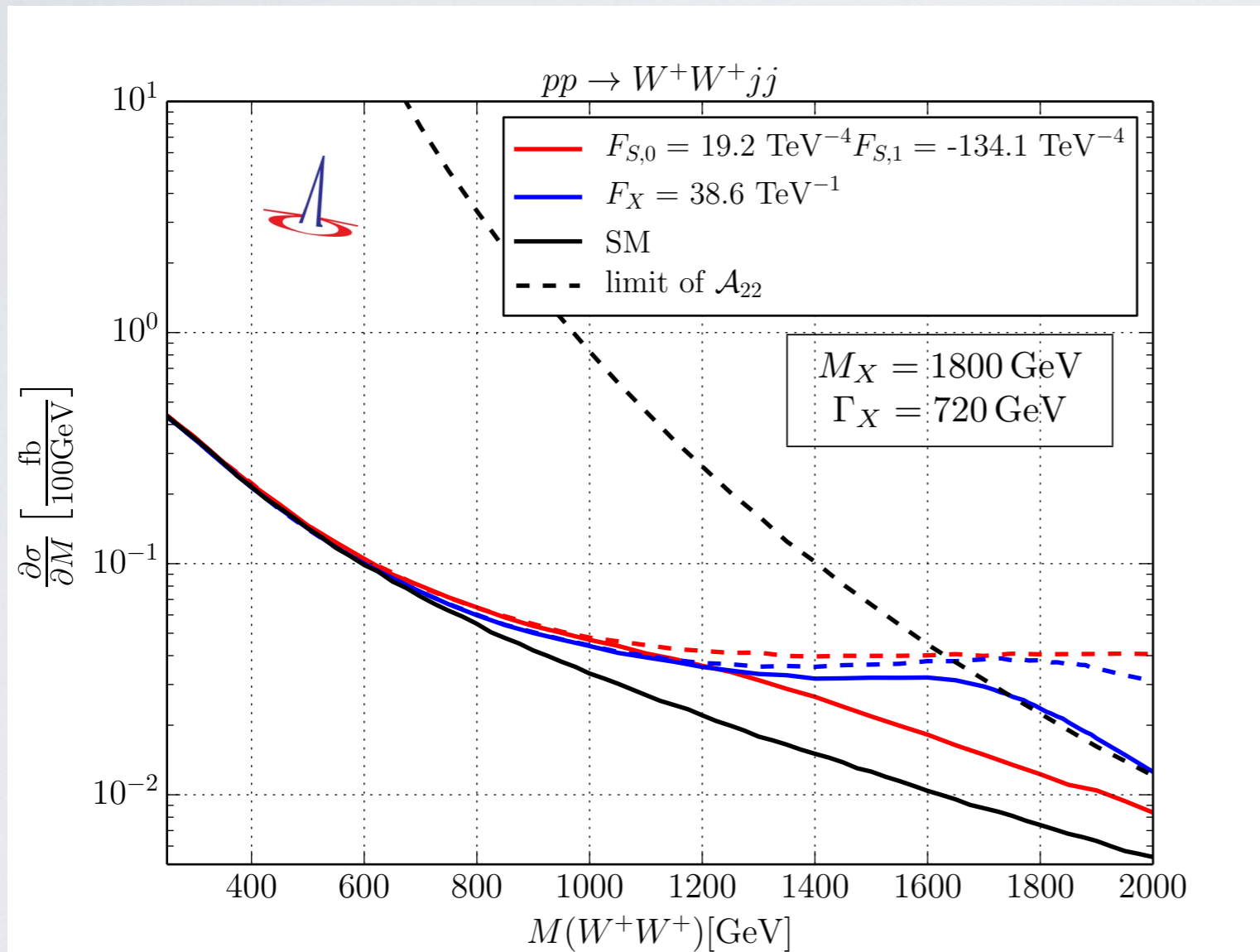
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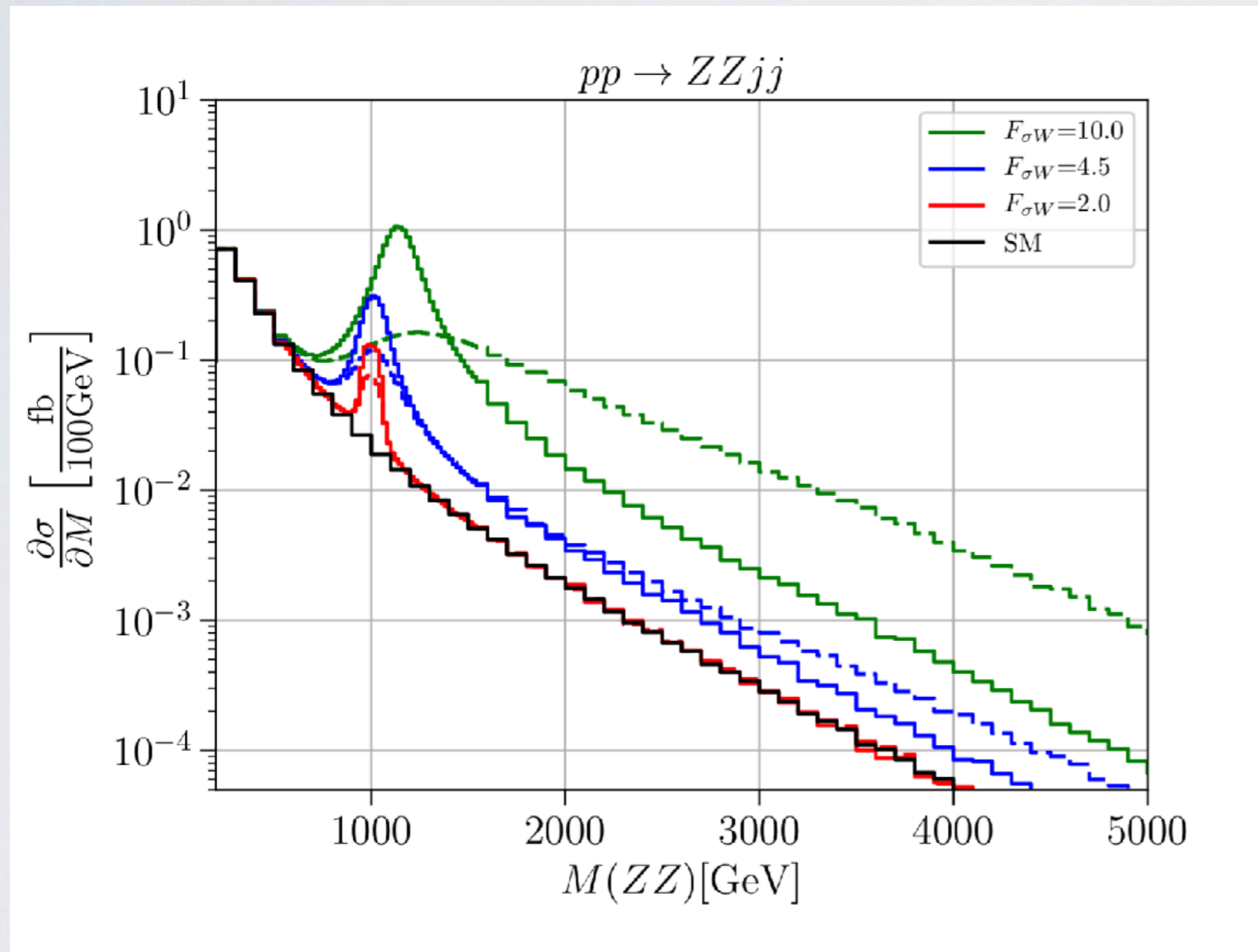
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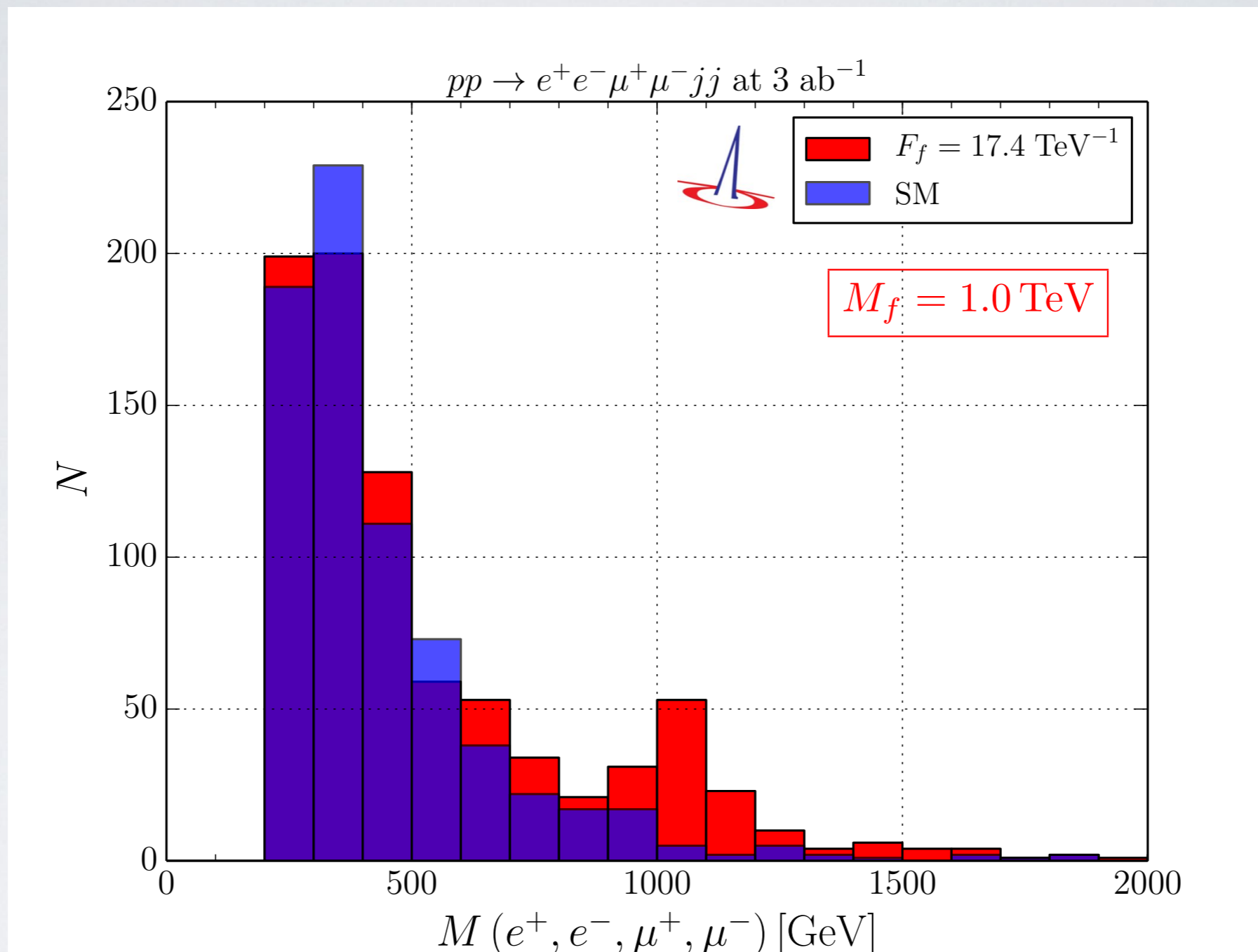
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Complete LHC process at 14 TeV

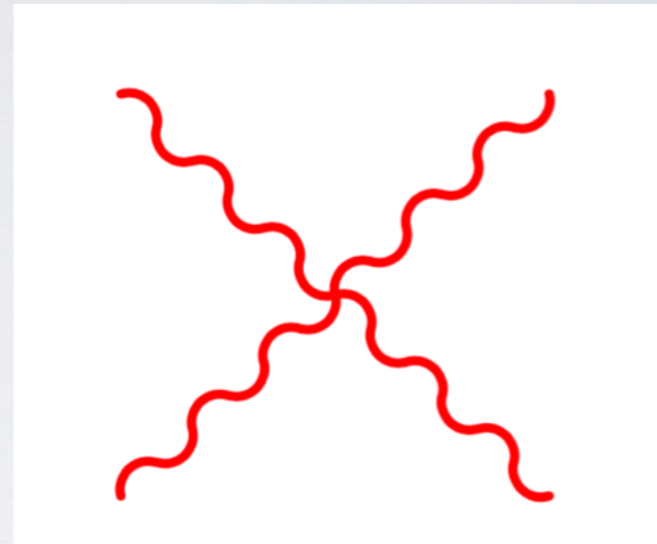


Triple [multiple] Vector Boson Production ?

Relate



to

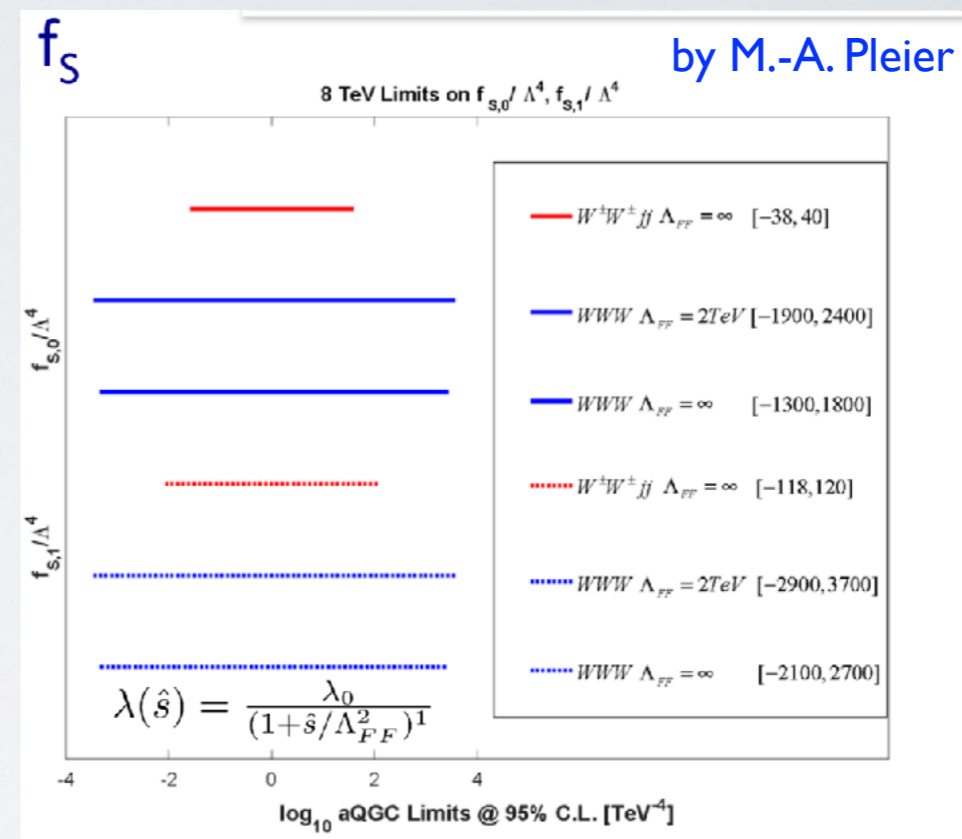
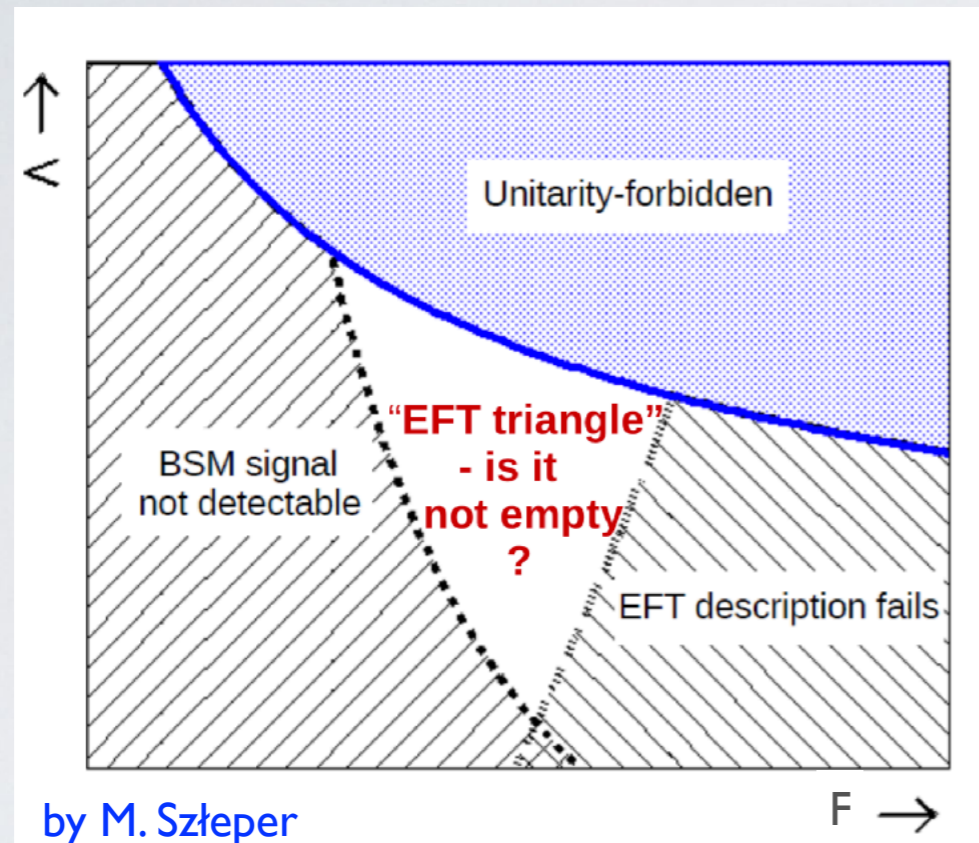


?

- ▶ Yes, same Feynman rule as in VBS, but ...
- ▶ one external $W/Z/\gamma$ always far off-shell
- ▶ Unitarization formalism: work in progress (needs $2 \rightarrow 3$ unitarizations)
- ▶ Different Wilson coefficients dominate (particularly for resonances)
- ▶ Important physics (partially) independent from VBS (“different fiducial vol.”)

Conclusions / Summary

- ◆ Vector boson scattering one of the flagship measurements of Runs II/III
- ◆ EFT provides **well-defined (and very limited) framework for SM deviations**



- ◆ There is not really a true model-independent parameterization!
- ◆ **Unitarization for theoretically sane description (allows reliable BDT analysis)**
- ◆ T -matrix unitarization universal scheme for EFT and resonances
- ◆ **Simplified models: generic electroweak resonances**
- ◆ Limits from LHC still quite limited: $\Lambda_{new\ physics} \sim 500\text{-}600\text{ GeV}$

Multi-Boson Interactions (MBI) 2018

August 28-30, 2018
<http://cern.ch/mbi2018>

University of Michigan
 Ann Arbor, MI, USA

6th Workshop on Multi-Boson Interactions

August 28-30, 2018

U. Michigan, Ann Arbor

1. TU Dresden
2. BNL (Brookhaven Ntl. Lab)
3. DESY
4. U. of Wisconsin — Madison
5. KIT Karlsruhe
6. U. of Michigan — Ann Arbor

Topics:

- Diboson and triboson production
- Higgs production
- Vector boson scattering and vector boson fusion
- Precision calculation and measurement of multiboson production
- New physics in multiboson production
- Monte Carlo generators
- LHC Run 2 results and beyond

Program committee:

John Campbell (FNAL), Sally Dawson (BNL), Lindsey Gray (FNAL), Christophe Grojean (DESY)
 Tao Han (U Pittsburgh), Matthew Herndon (UW Madison), Barbara Jäger (U Tübingen)
 Michael Kobel (TU Dresden), Sabine Lammers (Indiana U), Yurii Maravin (KSU)
 Marc-André Pleier (BNL), Aaron Pierce (Michigan), Jianming Qian (Michigan)
 Jürgen Reuter (DESY), James Wells (Michigan), Bing Zhou (Michigan), Junjie Zhu (Michigan)



BACKUP SLIDES



- ◆ Consider effects from heavy states by using (known) low-energy d.o.f.s

In addition to being a great convenience, effective field theory allows us to ask all the really scientific questions that we want to ask without committing ourselves to a picture of what happens at arbitrarily high energy.

H. Georgi, 1993

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- ◆ Integrating out heavy d.o.f.s marginalizes over details of short-distance interactions
- ◆ Toy Example: two interacting scalar fields φ, Φ

Path integral

$$\mathcal{Z}[j, J] = \int \mathcal{D}[\Phi] \mathcal{D}[\varphi] \exp \left[i \int dx \left(\frac{1}{2} (\partial\varphi)^2 - \frac{1}{2} \Phi (\square + M^2) \Phi - \lambda \varphi^2 \Phi - \dots + J \Phi + j \varphi \right) \right]$$

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Completing the square (Gaussian integration)

$$\Phi' = \Phi + \frac{\lambda}{M^2} \left(1 + \frac{\partial^2}{M^2} \right)^{-1} \varphi^2 \quad \Rightarrow \quad \text{Diagrammatic representation}$$

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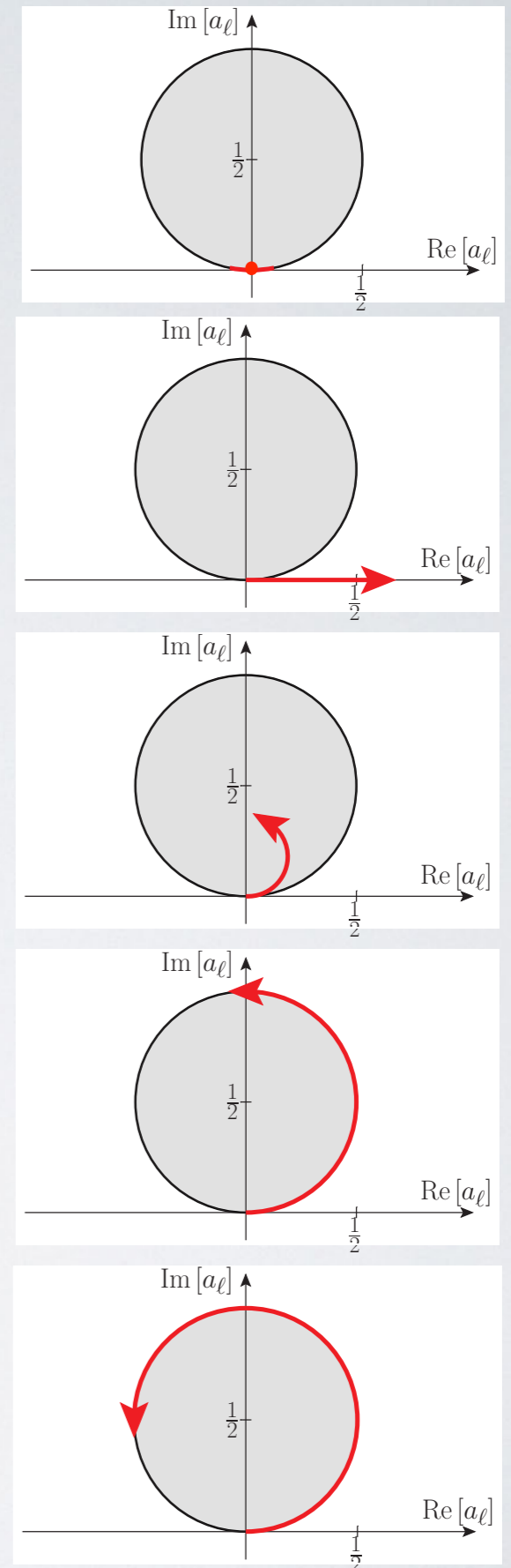
In the Lagrangian remove the high-scale d.o.f.s:

$$\frac{1}{2} (\partial\Phi)^2 - \frac{1}{2} M^2 \Phi^2 - \lambda \varphi^2 \Phi = \underbrace{-\frac{1}{2} \Phi' (M^2 + \partial^2) \Phi'}_{\text{Irrelevant normalization of the path integral}} + \underbrace{\frac{\lambda^2}{2M^2} \varphi^2 \left(1 + \frac{\partial^2}{M^2} \right)^{-1} \varphi^2}_{\text{Tower of higher and higher-dim. operators of light fields}}$$

Irrelevant normalization of the path integral

Tower of higher and higher-dim. operators of light fields

1. **SM or weakly coupled physics (e.g. 2HDM):** amplitude remains close to origin
2. **Rising amplitude (at least one dim-8 operator):** rise beyond unitarity circle [unphys.], strongly interacting regime
3. **Inelastic channel opens (form-factor description):** new channels open out, multi-boson final states
4. **Saturation of amplitude:** maximal amplitude, strongly interacting continuum, K-/T-matrix unitarization
5. **New resonance:** amplitude turns over



Unitarity in vector boson scattering

Optical Theorem (Unitarity of the S(cattering) Matrix):

$$\sigma_{\text{tot}} = \text{Im} [\mathcal{M}_{ii}(t = 0)] / s \quad t = -s(1 - \cos \theta)/2$$

Partial wave amplitudes:

$$\mathcal{M}(s, t, u) = 32\pi \sum_{\ell} (2\ell + 1) \mathcal{A}_{\ell}(s) P_{\ell}(\cos \theta) \quad (\text{"Power spectrum"})$$

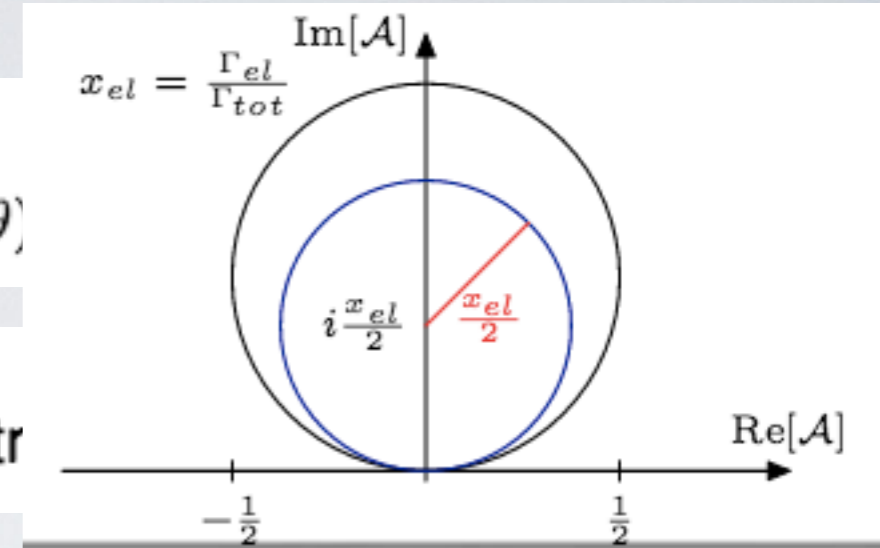
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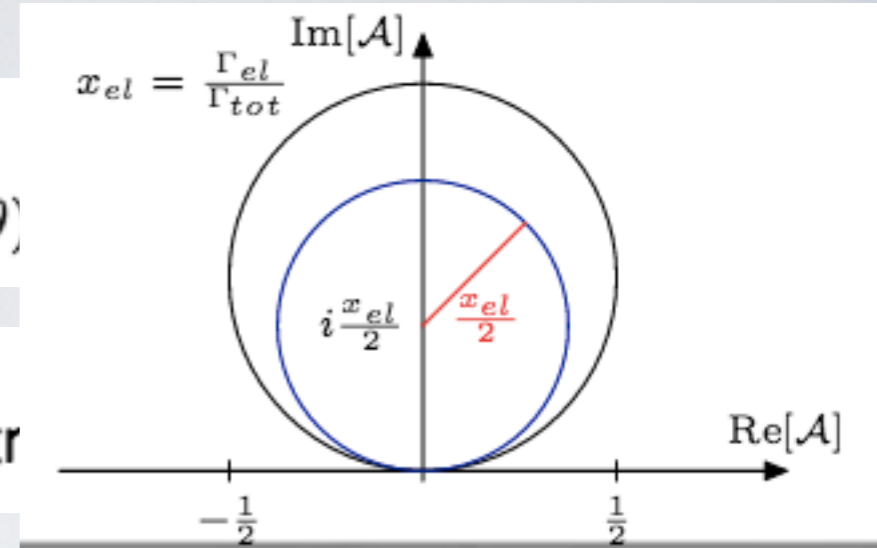
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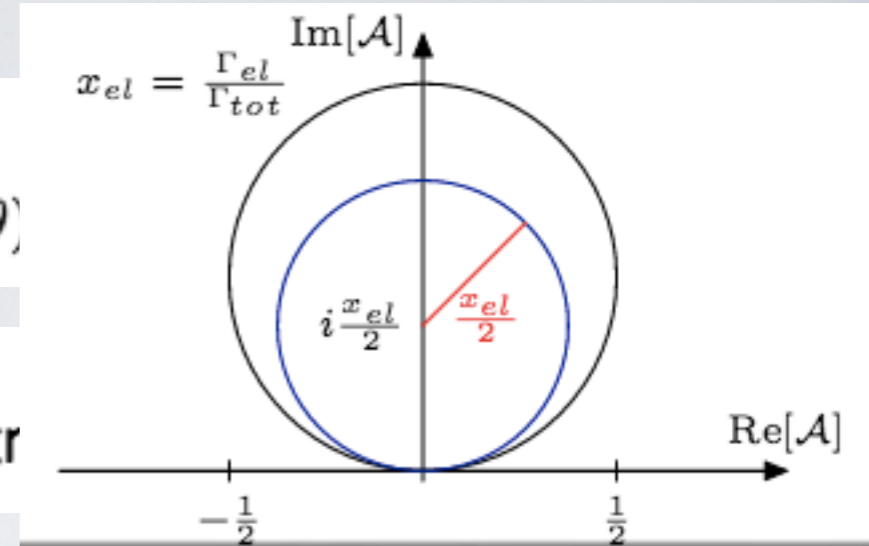
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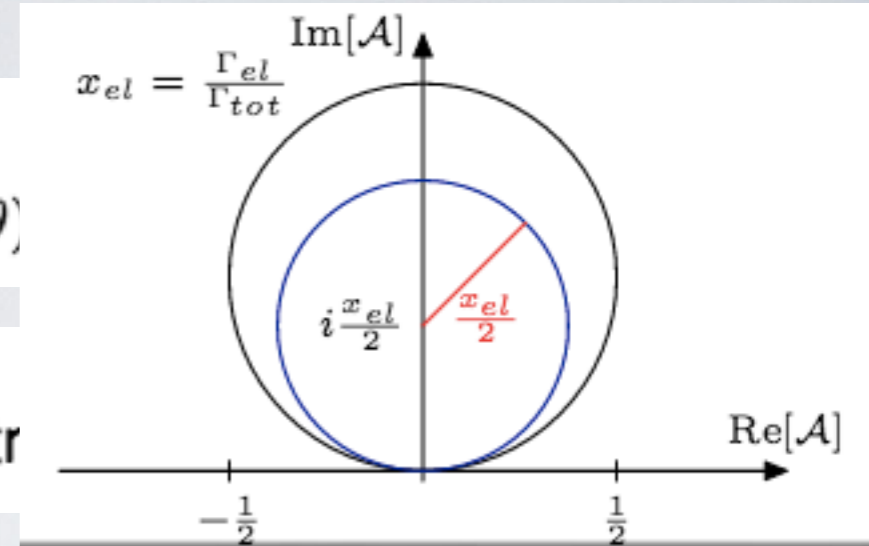
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Lee/Quigg/Thacker, 1973

exceeds unitarity bound $|\mathcal{A}_{IJ}| \lesssim \frac{1}{2}$ at:

$$I = 0 : \quad E \sim \sqrt{8\pi} v = 1.2 \text{ TeV}$$

$$I = 1 : \quad E \sim \sqrt{48\pi} v = 3.5 \text{ TeV}$$

$$I = 2 : \quad E \sim \sqrt{16\pi} v = 1.7 \text{ TeV}$$

Higgs exchange:

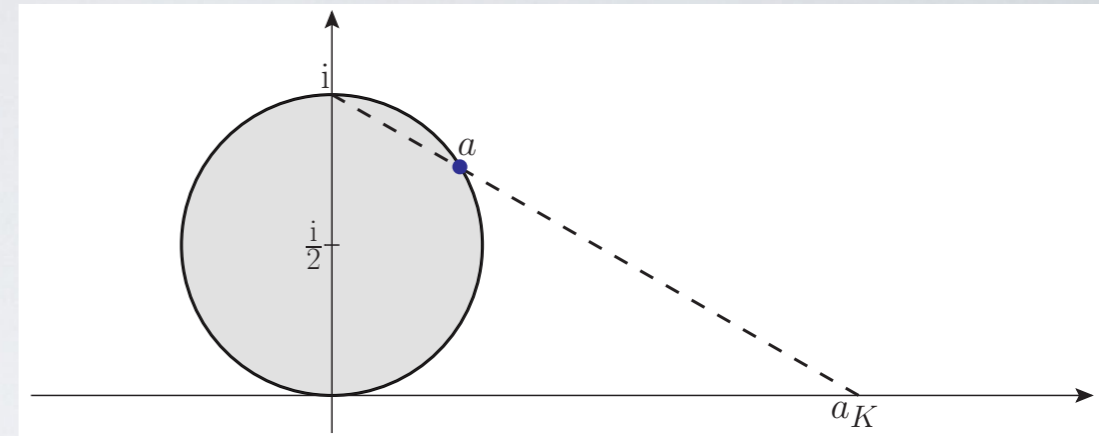
$$\mathcal{A}(s, t, u) = -\frac{M_H^2}{v^2} \frac{s}{s - M_H^2}$$

Unitarity: $M_H \lesssim \sqrt{8\pi} v \sim 1.2 \text{ TeV}$

Different unitarity projections

- **K-matrix:** Cayley transform of S-matrix Heitler, 1941; Schwinger, 1949; Gupta, 1950
- Stereographic projection to Argand circle

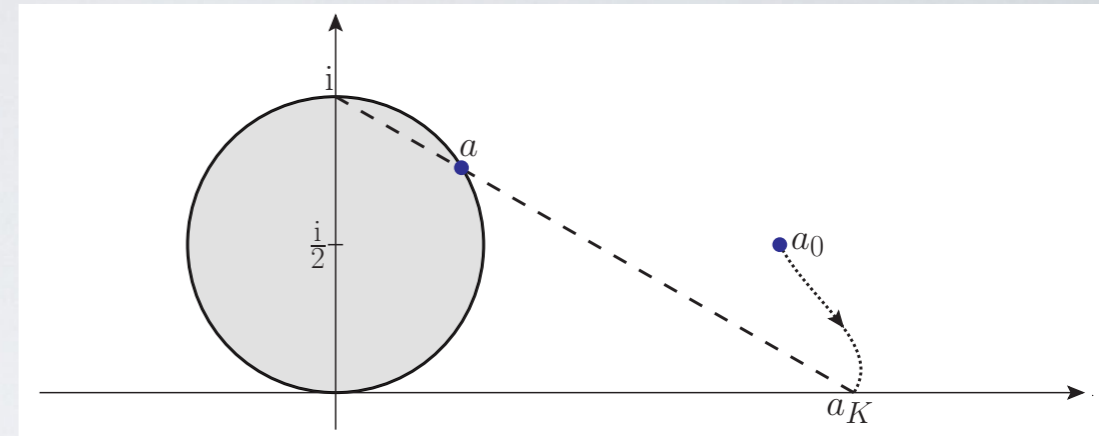
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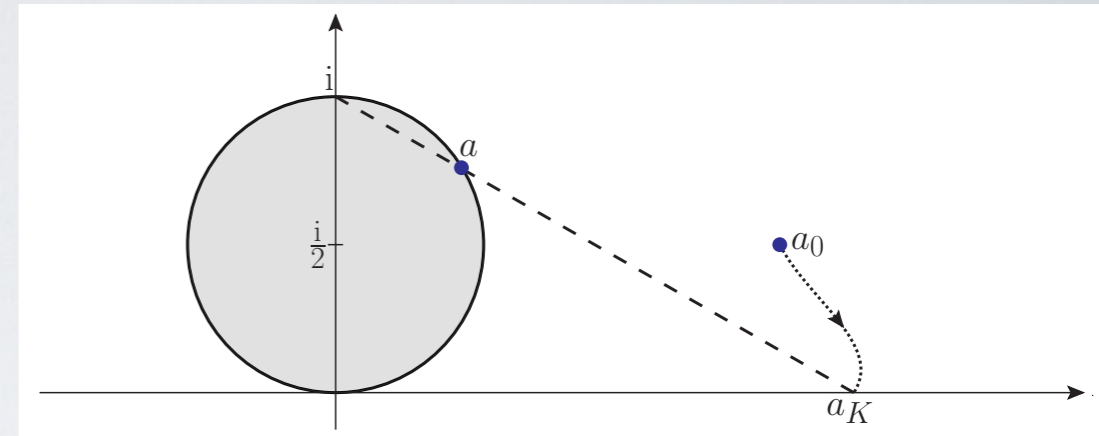
- Stereographic projection to Argand circle
- Formalism does a partial resummation of perturbative series
- **need to construct (orig.) K-matrix as self-adjoint intermediate operator**
Problems, if S-matrix non-diagonal, presence of non-perturbative contrib.

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- Stereographic projection to Argand circle

Formalism does a partial resummation of perturbative series

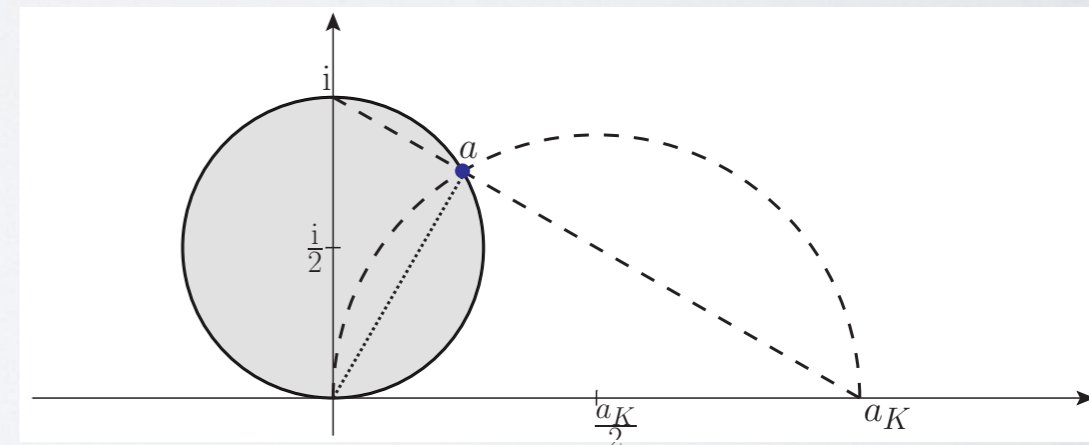
need to construct (orig.) K-matrix as self-adjoint intermediate operator

Problems, if S-matrix non-diagonal, presence of non-perturbative contrib.

- **T-matrix:** Thales circle construction

Kilian/Ohl/JRR/Sekulla, 2014

Defined via $|a - \frac{a_K}{2}| = \frac{a_K}{2} \Rightarrow a = \frac{1}{\text{Re}(\frac{1}{a_0}) - i}$

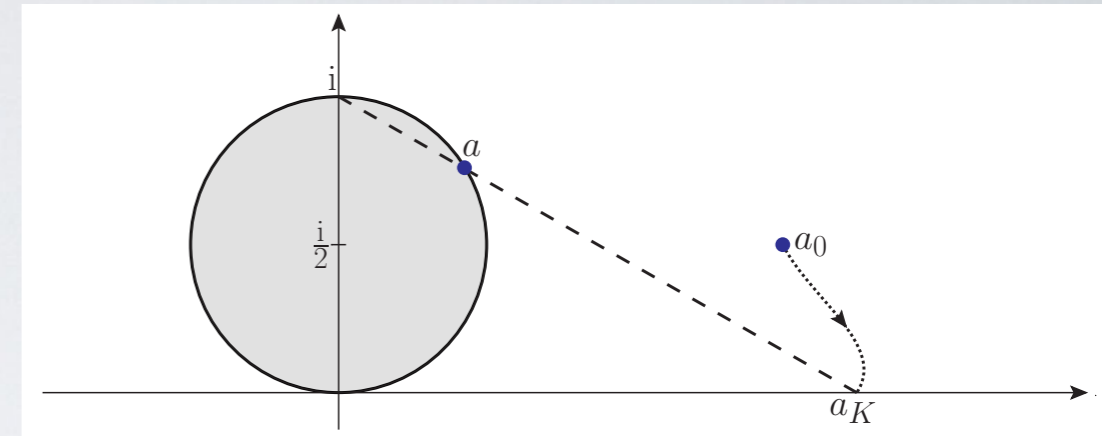


Different unitarity projections

- **K-matrix:** Cayley transform of S-matrix
- Stereographic projection to Argand circle

Heitler, 1941; Schwinger, 1949; Gupta, 1950

$$S = \frac{1+iK/2}{1-iK/2} \quad a_K(s) = \frac{a(s)}{1-ia(s)}$$



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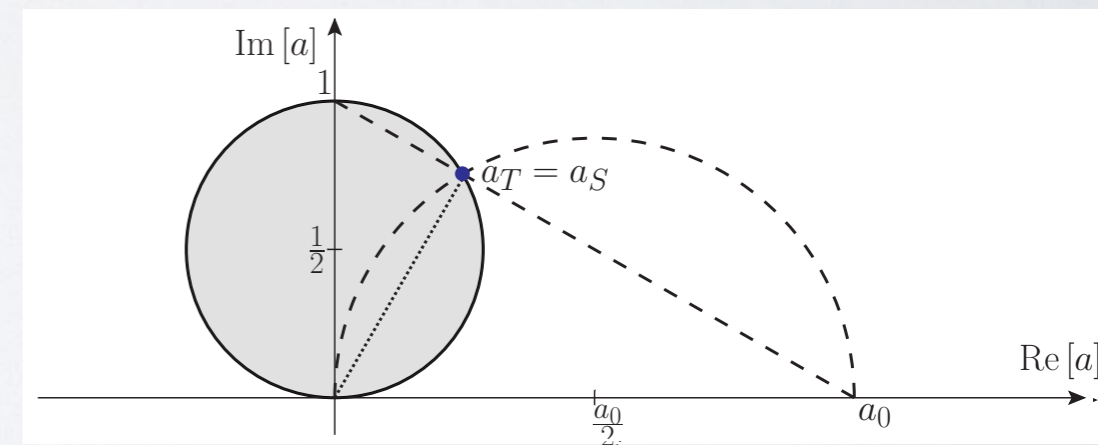
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- Identical to K matrix for real amplitudes

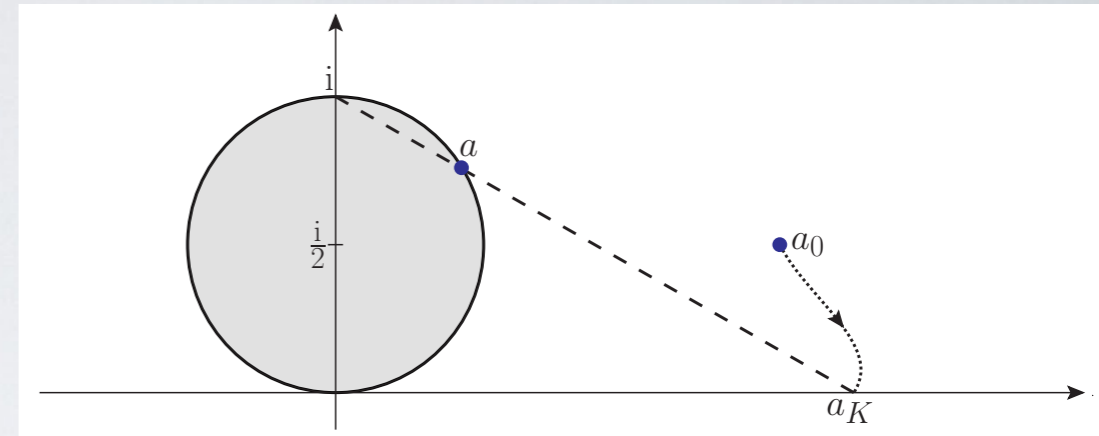
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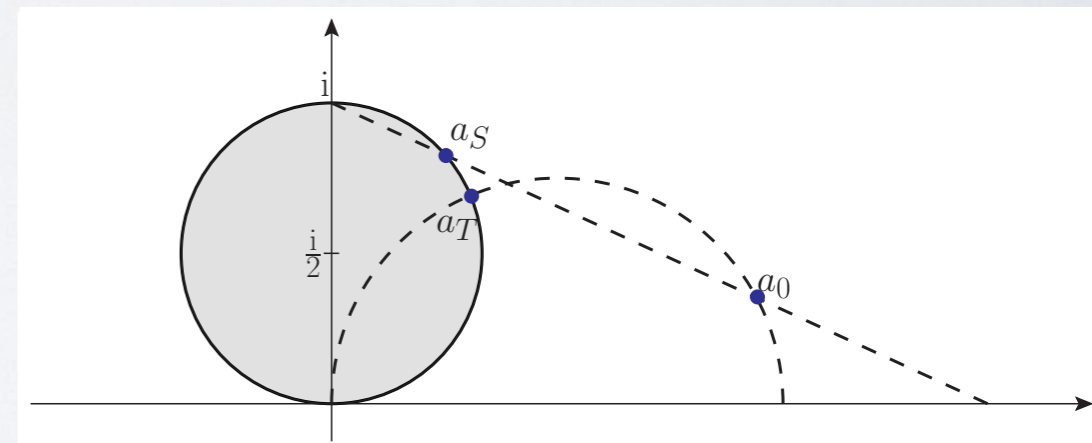
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