

# Mass and Mixing, Global Analysis

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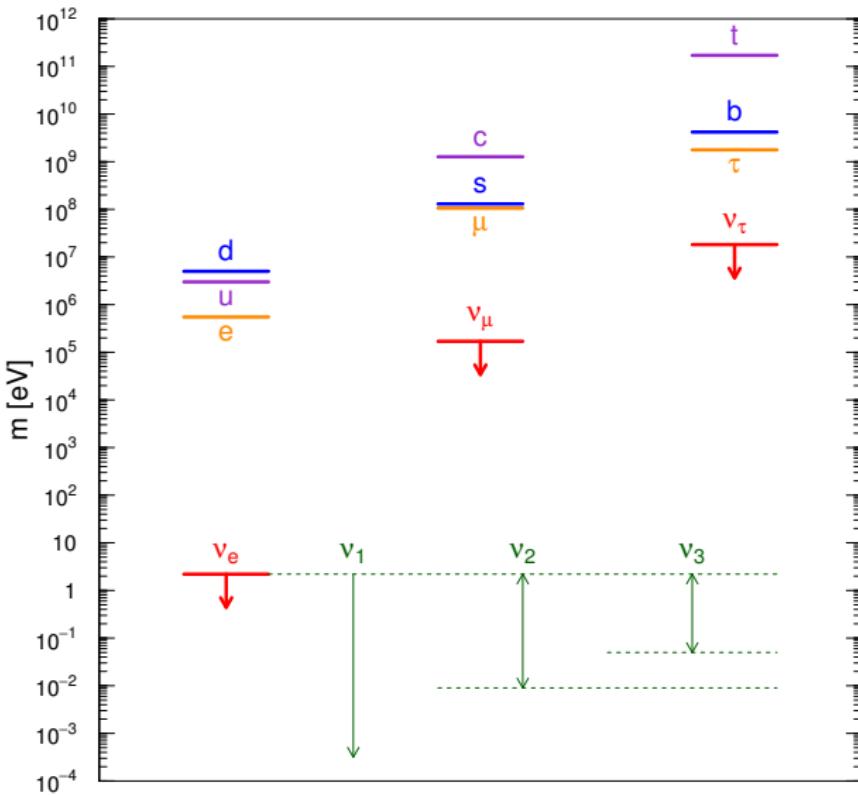
INFN, Torino, Italy

Rencontres du Vietnam 2017: Neutrinos

Qui Nhon, Vietnam, 16-22 July 2017



# Fermion Mass Spectrum



# Neutrino Mixing

Left-handed Flavor Neutrinos produced in Weak Interactions

$$|\nu_e, -\rangle \quad |\nu_\mu, -\rangle \quad |\nu_\tau, -\rangle$$

$$\mathcal{H}_{CC} = \frac{g}{\sqrt{2}} W_\rho (\overline{\nu_{eL}} \gamma^\rho e_L + \overline{\nu_{\mu L}} \gamma^\rho \mu_L + \overline{\nu_{\tau L}} \gamma^\rho \tau_L) + \text{H.c.}$$

Fields  $\nu_{\alpha L} = \sum_k U_{\alpha k} \nu_{kL} \implies |\nu_\alpha, -\rangle = \sum_k U_{\alpha k}^* |\nu_k, -\rangle$  States

$$|\nu_1, -\rangle \quad |\nu_2, -\rangle \quad |\nu_3, -\rangle$$

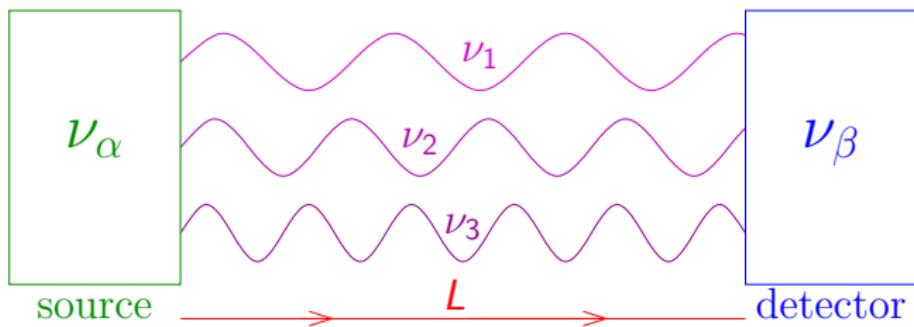
Left-handed Massive Neutrinos propagate from Source to Detector

3 × 3 Unitary Mixing Matrix:

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix}$$

# Neutrino Oscillations

$$|\nu(t=0)\rangle = |\nu_\alpha\rangle = U_{\alpha 1}^* |\nu_1\rangle + U_{\alpha 2}^* |\nu_2\rangle + U_{\alpha 3}^* |\nu_3\rangle$$



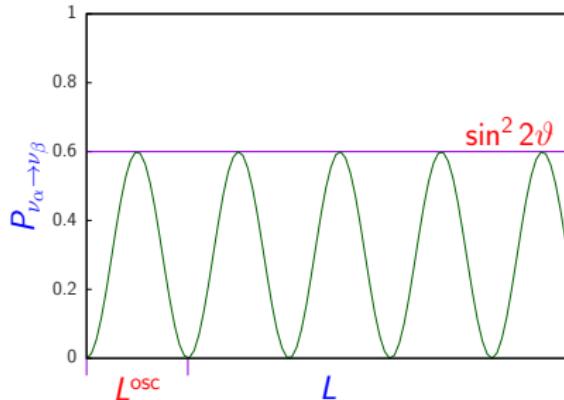
$$|\nu(t > 0)\rangle = U_{\alpha 1}^* e^{-iE_1 t} |\nu_1\rangle + U_{\alpha 2}^* e^{-iE_2 t} |\nu_2\rangle + U_{\alpha 3}^* e^{-iE_3 t} |\nu_3\rangle \neq |\nu_\alpha\rangle$$

$$E_k^2 = p^2 + m_k^2 \quad t = L$$

$$P_{\nu_\alpha \rightarrow \nu_\beta}(L) = |\langle \nu_\beta | \nu(L) \rangle|^2 = \sum_{k,j} U_{\beta k} U_{\alpha k}^* U_{\beta j}^* U_{\alpha j} \exp\left(-i \frac{\Delta m_{kj}^2 L}{2E}\right)$$

the oscillation probabilities depend on  $U$  and  $\Delta m_{kj}^2 \equiv m_k^2 - m_j^2$

$$2\nu\text{-mixing: } P_{\nu_\alpha \rightarrow \nu_\beta} = \sin^2 2\vartheta \sin^2 \left( \frac{\Delta m^2 L}{4E} \right) \implies L^{\text{osc}} = \frac{4\pi E}{\Delta m^2}$$



Tiny neutrino masses lead to observable macroscopic oscillation distances!

$$L \gtrsim \begin{cases} 10 \frac{m}{MeV} \left( \frac{km}{GeV} \right) & \text{short-baseline experiments} & \Delta m^2 \gtrsim 10^{-1} \text{ eV}^2 \\ 10^3 \frac{m}{MeV} \left( \frac{km}{GeV} \right) & \text{long-baseline experiments} & \Delta m^2 \gtrsim 10^{-3} \text{ eV}^2 \\ 10^4 \frac{km}{GeV} & \text{atmospheric neutrino experiments} & \Delta m^2 \gtrsim 10^{-4} \text{ eV}^2 \\ 10^{11} \frac{m}{MeV} & \text{solar neutrino experiments} & \Delta m^2 \gtrsim 10^{-11} \text{ eV}^2 \end{cases}$$

Neutrino oscillations are the optimal tool to reveal tiny neutrino masses!

# Three-Neutrino Mixing Paradigm

Standard Parameterization of Mixing Matrix (as CKM)

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{13}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix}$$
$$= \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{13}} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta_{13}} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta_{13}} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta_{13}} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta_{13}} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix}$$

$$c_{ab} \equiv \cos \vartheta_{ab} \quad s_{ab} \equiv \sin \vartheta_{ab} \quad 0 \leq \vartheta_{ab} \leq \frac{\pi}{2} \quad 0 \leq \delta_{13}, \lambda_{21}, \lambda_{31} < 2\pi$$

OSCILLATION  
PARAMETERS

{ 3 Mixing Angles:  $\vartheta_{12}$ ,  $\vartheta_{23}$ ,  $\vartheta_{13}$   
1 CPV Dirac Phase:  $\delta_{13}$   
2 independent  $\Delta m_{kj}^2 \equiv m_k^2 - m_j^2$ :  $\Delta m_{21}^2$ ,  $\Delta m_{31}^2$

2 CPV Majorana Phases:  $\lambda_{21}$ ,  $\lambda_{31} \iff |\Delta L| = 2$  processes

# Three-Neutrino Mixing Ingredients

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{13}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix}$$

Solar  
 $\nu_e \rightarrow \nu_\mu, \nu_\tau$   
VLBL Reactor  
 $\bar{\nu}_e$  disappearance

$$\left. \begin{array}{c} \text{SNO, Borexino} \\ \text{Super-Kamiokande} \\ \text{GALLEX/GNO, SAGE} \\ \text{Homestake, Kamiokande} \\ \text{(KamLAND)} \end{array} \right\} \rightarrow \left\{ \begin{array}{l} \Delta m_S^2 = \Delta m_{21}^2 \simeq 7.4 \times 10^{-5} \text{ eV}^2 \\ \sin^2 \vartheta_S = \sin^2 \vartheta_{12} \simeq 0.30 \end{array} \right.$$

# Three-Neutrino Mixing Ingredients

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{13}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix}$$

Atmospheric  
 $\nu_\mu \rightarrow \nu_\tau$

Super-Kamiokande  
Kamiokande, IMB  
MACRO, Soudan-2

LBL Accelerator  
 $\nu_\mu$  disappearance

K2K, MINOS  
T2K, NO $\nu$ A

LBL Accelerator  
 $\nu_\mu \rightarrow \nu_\tau$

(OPERA)



$$\left. \begin{array}{l} \Delta m_A^2 \simeq |\Delta m_{31}^2| \simeq 2.5 \times 10^{-3} \text{ eV}^2 \\ \sin^2 \vartheta_A = \sin^2 \vartheta_{23} \simeq 0.50 \end{array} \right\}$$

# Three-Neutrino Mixing Ingredients

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{13}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{13}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\lambda_{21}} & 0 \\ 0 & 0 & e^{i\lambda_{31}} \end{pmatrix}$$

LBL Accelerator

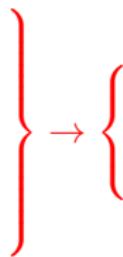
$$\nu_\mu \rightarrow \nu_e$$

(T2K, MINOS, NO $\nu$ A)

LBL Reactor

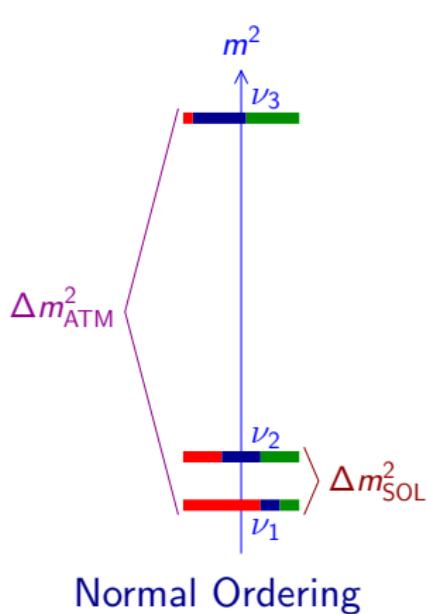
$\bar{\nu}_e$  disappearance

$\begin{pmatrix} \text{Daya Bay, RENO} \\ \text{Double Chooz} \end{pmatrix}$



$$\left\{ \begin{array}{l} \Delta m_A^2 \simeq |\Delta m_{31}^2| \simeq 2.5 \times 10^{-3} \text{ eV}^2 \\ \sin^2 \vartheta_{13} \simeq 0.022 \end{array} \right.$$

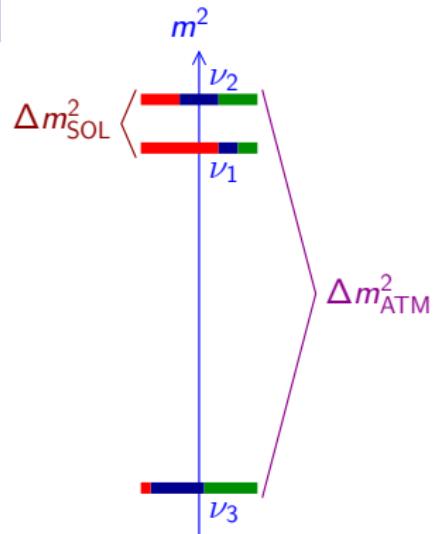
# Mass Ordering



Normal Ordering

$$\Delta m_{31}^2 > \Delta m_{32}^2 > 0$$

absolute scale is not determined by neutrino oscillation data



Inverted Ordering

$$\Delta m_{32}^2 < \Delta m_{31}^2 < 0$$

# CP Transformation

Right-handed antineutrinos are described by CP-conjugated states

$$\text{Particle} \xrightleftharpoons{\text{C}} \text{Antiparticle}$$

$$\text{Left-Handed Helicity} \xrightleftharpoons{\text{P}} \text{Right-Handed Helicity}$$

$$|\nu_\alpha, -\rangle = \sum_k U_{\alpha k}^* |\nu_k, -\rangle \xrightleftharpoons{\text{CP}} |\bar{\nu}_\alpha, +\rangle = \sum_k U_{\alpha k} |\bar{\nu}_k, +\rangle$$

In oscillation probabilities: Neutrino  $U \xrightleftharpoons{\text{CP}} U^*$  Antineutrino

$$\begin{aligned}
 P_{\nu_\alpha \rightarrow \nu_\beta} &= \delta_{\alpha\beta} - 4 \sum_{k>j} \text{Re}[U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^*] \sin^2 \left( \frac{\Delta m_{kj}^2 L}{4E} \right) \quad \leftarrow \text{CP Even} \\
 &\quad + 2 \sum_{k>j} \text{Im}[U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^*] \sin \left( \frac{\Delta m_{kj}^2 L}{2E} \right) \quad \leftarrow \text{CP Odd}
 \end{aligned}$$

Survival probabilities:

$$P_{\nu_\alpha \rightarrow \nu_\alpha} = P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\alpha}$$

CPT

# CP Asymmetries

$$\begin{aligned} A_{\alpha\beta}^{\text{CP}} &= P_{\nu_\alpha \rightarrow \nu_\beta} - P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta} \\ &= 16 J_{\alpha\beta} \sin\left(\frac{\Delta m_{21}^2 L}{4E}\right) \sin\left(\frac{\Delta m_{31}^2 L}{4E}\right) \sin\left(\frac{\Delta m_{32}^2 L}{4E}\right) \end{aligned}$$

$$J_{\alpha\beta} = \text{Im}[U_{\alpha 1}^* U_{\beta 1} U_{\alpha 2} U_{\beta 2}^*] = \pm J_{\text{CP}} \quad \text{Jarlskog Invariant}$$

$$J_{\text{CP}} = \text{Im}[U_{\mu 1}^* U_{e 1} U_{\mu 2} U_{e 2}^*] = c_{12} s_{12} c_{23} s_{23} c_{13}^2 s_{13} \sin \delta_{13}$$

$$J_{\text{CP}} \neq 0 \iff \vartheta_{12}, \vartheta_{23}, \vartheta_{13} \neq 0, \pi/2 \quad \text{and} \quad \delta_{13} \neq 0, \pi$$

Necessary conditions for observation of CP violation:

- ▶ Sensitivity to all mixing angles, including small  $\vartheta_{13}$ .
- ▶ Sensitivity to oscillations due to  $\Delta m_{21}^2$  and  $\Delta m_{31}^2$ .

## LBL $\nu_\mu \rightarrow \nu_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E} \quad A = \frac{2EV}{\Delta m_{31}^2} \quad V = \sqrt{2} G_F N_e$$

$$\sin \theta_{13} \ll 1 \quad \Delta m_{21}^2 / \Delta m_{31}^2 \ll 1$$

$$P_{\nu_\mu \rightarrow \nu_e}^{\text{LBL}} \simeq \sin^2 2\vartheta_{13} \overset{\text{octant}}{\downarrow} \sin^2 \vartheta_{23} \frac{\sin^2[(1-A)\Delta]}{(1-A)^2}$$

$$+ \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \sin 2\vartheta_{13} \sin 2\vartheta_{12} \sin 2\vartheta_{23} \cos(\Delta + \delta_{13}) \frac{\sin(A\Delta)}{A} \frac{\sin[(1-A)\Delta]}{1-A}$$

$$+ \left( \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \right)^2 \sin^2 2\vartheta_{12} \cos^2 \vartheta_{23} \frac{\sin^2(A\Delta)}{A^2} \overset{\text{CPV}}{\uparrow}$$

NO:  $\Delta m_{31}^2 > 0$

IO:  $\Delta m_{31}^2 < 0$

for antineutrinos:  $\delta_{13} \rightarrow -\delta_{13}$  (CPV) and  $A \rightarrow -A$  (Fake CPV!)

[see: Mezzetto, Schwetz, JPG 37 (2010) 103001]

# Global Fits

NuFIT 3.0 (2016)

[arXiv:1611.01514]

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[arXiv:1703.04471]



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## Updated fit to three neutrino mixing: exploring the accelerator-reactor complementarity

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JHEP01(2017)087

PHYSICAL REVIEW D 95, 096014 (2017)

## Global constraints on absolute neutrino masses and their ordering

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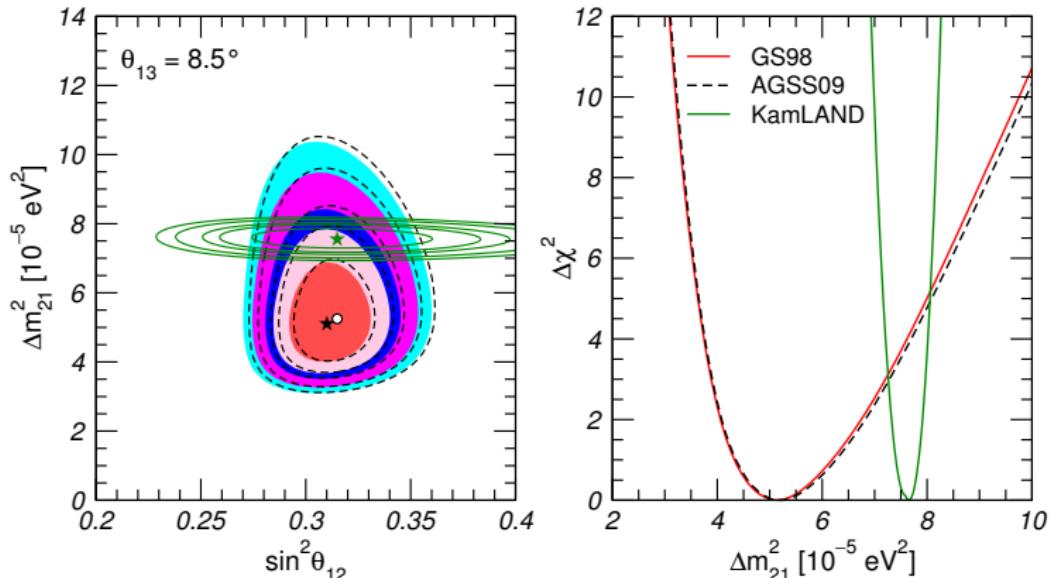
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# Solar Neutrinos

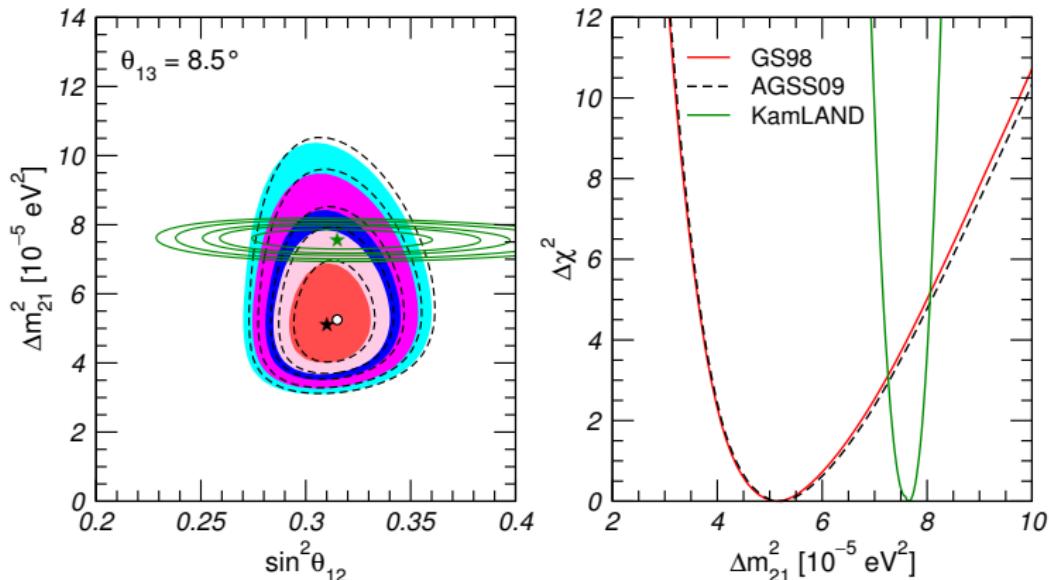
NuFIT 3.0 (2016)



- Metallicity (C, N, O, Ne, Mg, Si, S, Ar, Fe):  
high (GS98)      low (AGSS09)
- A new Generation of Standard Solar Models:  
Vinyoles et al., arXiv:1611.09867

# Solar Neutrinos

NuFIT 3.0 (2016)



## Solar – KamLAND Tension

- ▶ No SK+SNO low-energy spectrum up-turn expected for  $(\Delta m_{21}^2)_{KL}$
- ▶ Larger SK day-night asymmetry than expected for  $(\Delta m_{21}^2)_{KL}$

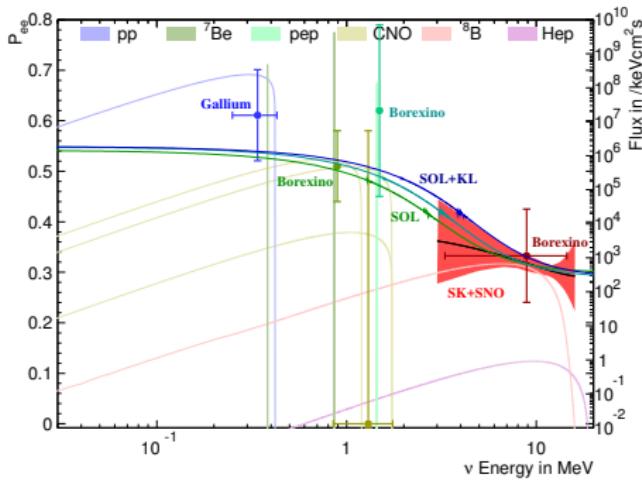
# Solar Neutrino Spectrum

$$\begin{aligned}\overline{P}_{ee}^{\text{SOL}} &= \sum_{k=1}^3 |U_{ek}|^2 |U_{ek}^0|^2 = \left( \cos^2 \vartheta_{12} \cos^2 \vartheta_{12}^0 + \sin^2 \vartheta_{12} \sin^2 \vartheta_{12}^0 \right) \cos^4 \vartheta_{13} + \sin^4 \vartheta_{13} \\ &= \left( \frac{1}{2} + \frac{1}{2} \cos 2\vartheta_{12}^0 \cos 2\vartheta_{12} \right) \cos^4 \vartheta_{13} + \sin^4 \vartheta_{13}\end{aligned}$$

Averaged  
Vacuum  
Oscillations

$$\theta_{12}^0 \simeq \theta_{12}$$

$$\begin{aligned}\overline{P}_{ee}^{\text{SOL}} &\simeq \left( 1 - \frac{1}{2} \sin^2 \vartheta_{12} \right) \\ &\times \left( 1 - \sin^2 \vartheta_{13} \right)\end{aligned}$$



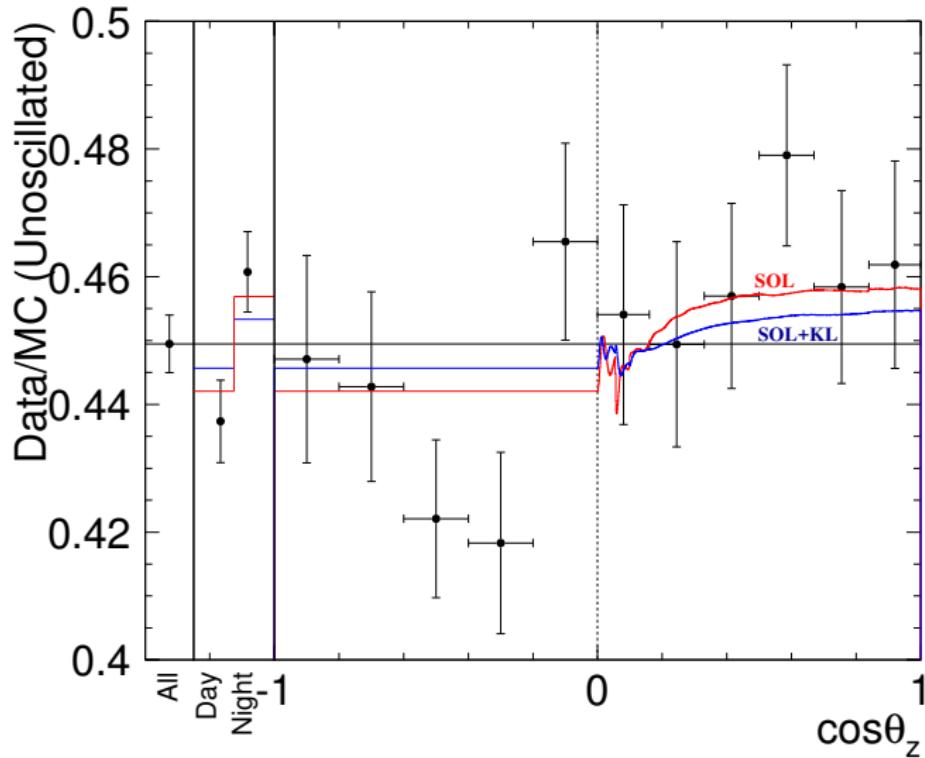
[SK, arXiv:1606.07538]

Adiabatic  
MSW  
Transitions

$$\theta_{12}^0 \simeq \pi/2$$

$$\begin{aligned}\overline{P}_{ee}^{\text{SOL}} &\simeq \sin^2 \vartheta_{12} \\ &\times \left( 1 - \sin^2 \vartheta_{13} \right)\end{aligned}$$

# SK Day-Night Asymmetry

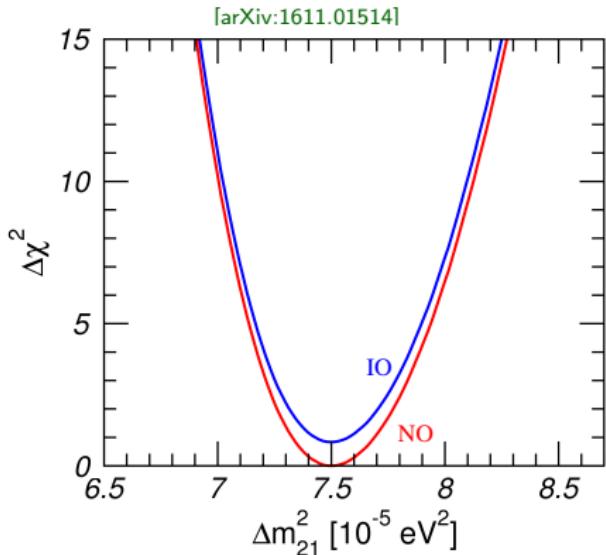


[SK, arXiv:1606.07538]

$$\Delta m_{21}^2$$

NuFIT 3.0 (2016)

[arXiv:1611.01514]



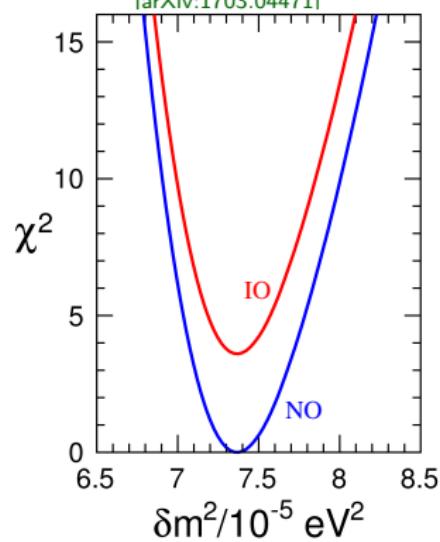
$$\Delta m_{21}^2 = 7.50^{+0.19}_{-0.17} \times 10^{-5} \text{ eV}^2$$

$\sim 2.5\%$  precision

$\sim 1.7\%$  difference

Bari 2017

[arXiv:1703.04471]



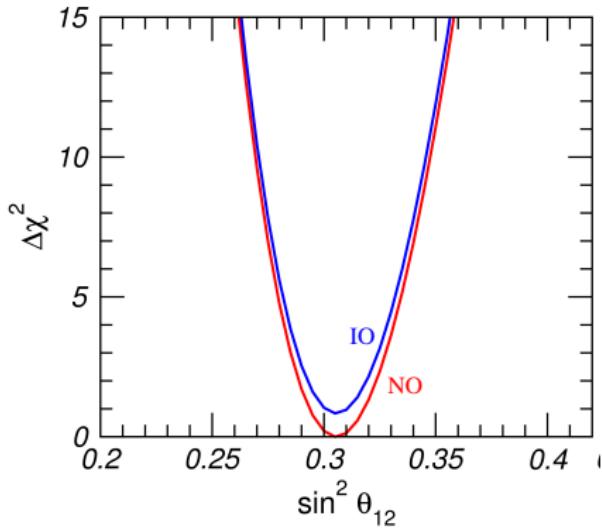
$$\Delta m_{21}^2 = 7.37^{+0.17}_{-0.16} \times 10^{-5} \text{ eV}^2$$

$\sim 2.3\%$  precision

$$\sin^2 \vartheta_{12}$$

NuFIT 3.0 (2016)

[arXiv:1611.01514]



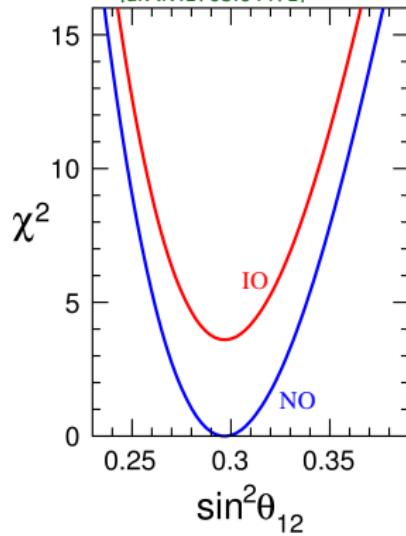
$$\sin^2 \vartheta_{12} = 0.306 \pm 0.012$$

$\sim 3.9\%$  precision

$\sim 3.0\%$  difference

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[arXiv:1703.04471]



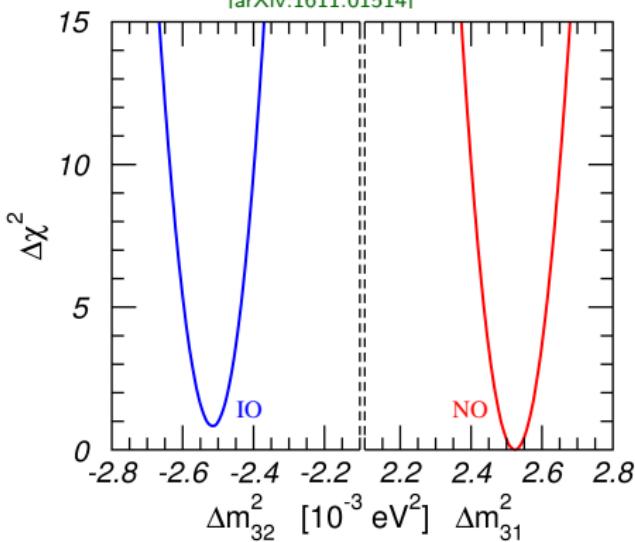
$$\sin^2 \vartheta_{12} = 0.297^{+0.017}_{-0.016}$$

$\sim 5.7\%$  precision

# $\Delta m_{31}^2$ & $\Delta m_{32}^2$

NuFIT 3.0 (2016)

[arXiv:1611.01514]



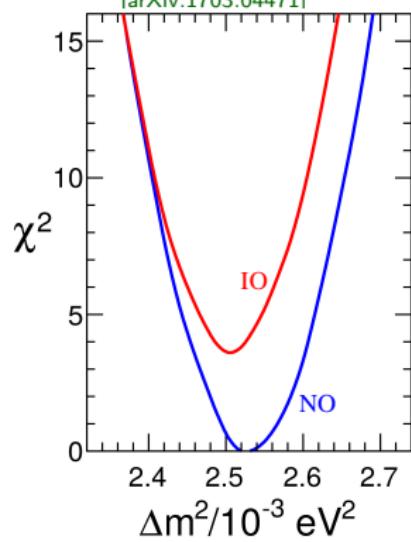
$$\frac{|\Delta m_{31}^2|}{10^{-3} \text{ eV}^2} = \begin{cases} 2.524^{+0.039}_{-0.040} & (\text{NO}) \\ 2.439^{+0.041}_{-0.038} & (\text{IO}) \end{cases}$$

~ 1.7% precision

~ 1.5% difference

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[arXiv:1703.04471]



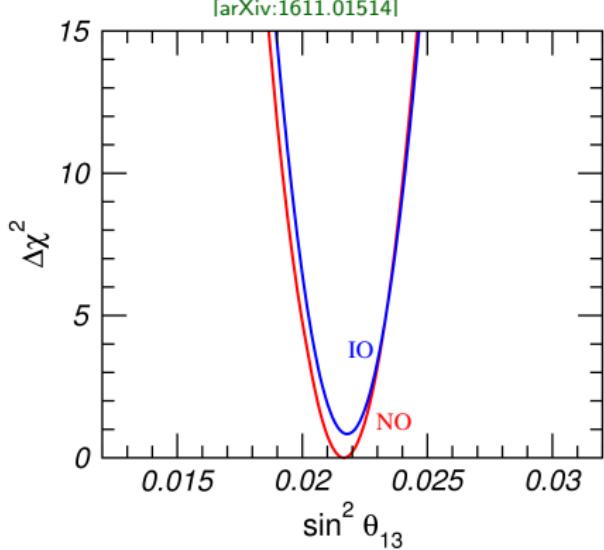
$$\frac{|\Delta m_{31}^2|}{10^{-3} \text{ eV}^2} = \begin{cases} 2.562^{+0.042}_{-0.030} & (\text{NO}) \\ 2.468^{+0.034}_{-0.32} & (\text{IO}) \end{cases}$$

~ 1.6% precision

# $\sin^2 \vartheta_{13}$

NuFIT 3.0 (2016)

[arXiv:1611.01514]

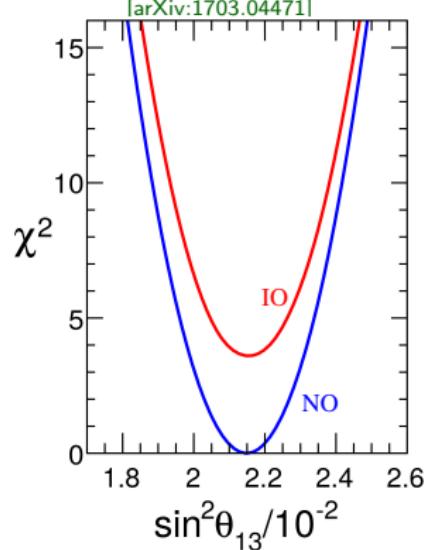


$$\sin^2 \vartheta_{13} = \begin{cases} 0.02166 \pm 0.00075 & (\text{NO}) \\ 0.02179 \pm 0.00076 & (\text{IO}) \end{cases}$$

$\sim 3.5\%$  precision

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[arXiv:1703.04471]



$$\sin^2 \vartheta_{13} = \begin{cases} 0.0215 \pm 0.0007 & (\text{NO}) \\ 0.0216^{+0.0008}_{-0.0009} & (\text{IO}) \end{cases}$$

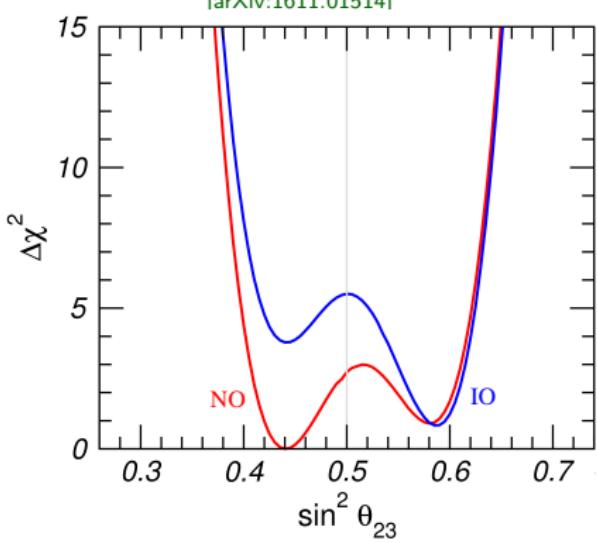
$\sim 3.3\%$  precision

$\sim 0.8\%$  difference

$$\sin^2 \vartheta_{23}$$

NuFIT 3.0 (2016)

[arXiv:1611.01514]



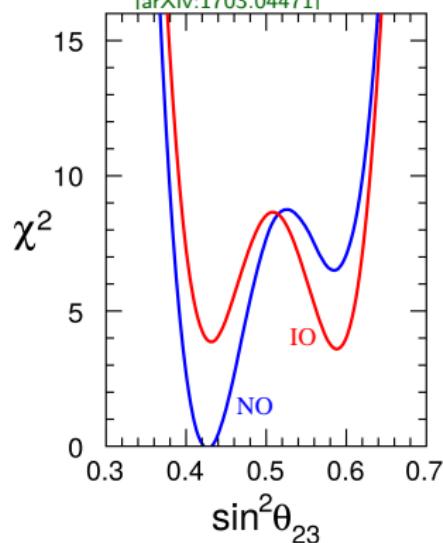
$$\sin^2 \vartheta_{23} = \begin{cases} 0.441^{+0.027}_{-0.021} & (\text{NO}) \\ 0.587^{+0.020}_{-0.024} & (\text{IO}) \end{cases}$$

$\sim 9\%$  precision

Common NO/IO octant flip

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[arXiv:1703.04471]

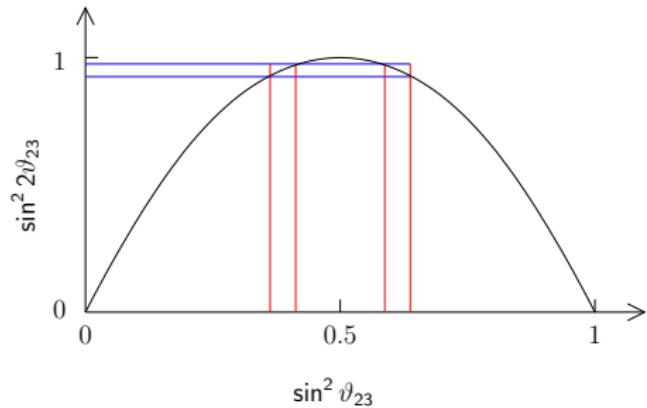
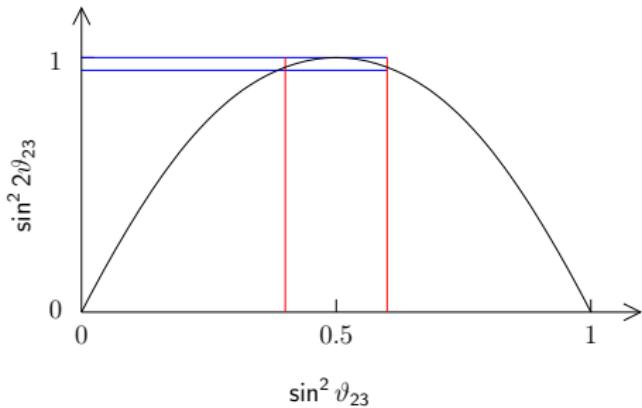


$$\sin^2 \vartheta_{23} = \begin{cases} 0.425^{+0.021}_{-0.015} & (\text{NO}) \\ 0.589^{+0.016}_{-0.022} \oplus [0.417, 0.448] & (\text{IO}) \end{cases}$$

$\sim 9\%$  precision

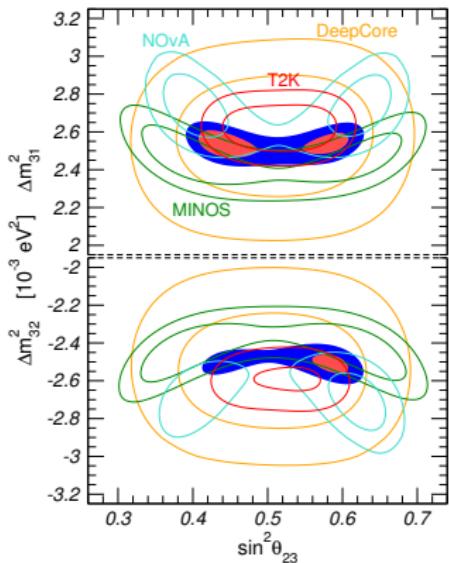
$$P_{\nu_\mu \rightarrow \nu_\mu}^{\text{LBL}} \simeq 1 - \sin^2 2\vartheta_{23} \sin^2 \left( \frac{\Delta m_{31}^2 L}{4E} \right)$$

$$\sin^2 2\vartheta_{23} = 4 \sin^2 \vartheta_{23} (1 - \sin^2 \vartheta_{23})$$



# NuFIT 3.0 (2016)

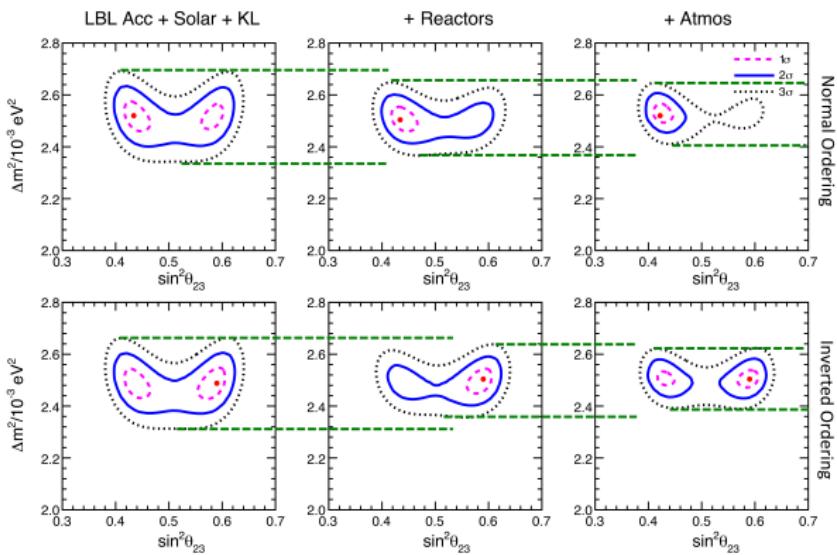
[arXiv:1611.01514]



No SK atmospheric data

# Bari 2017

[Lisi @ EPS-HEP 2017]

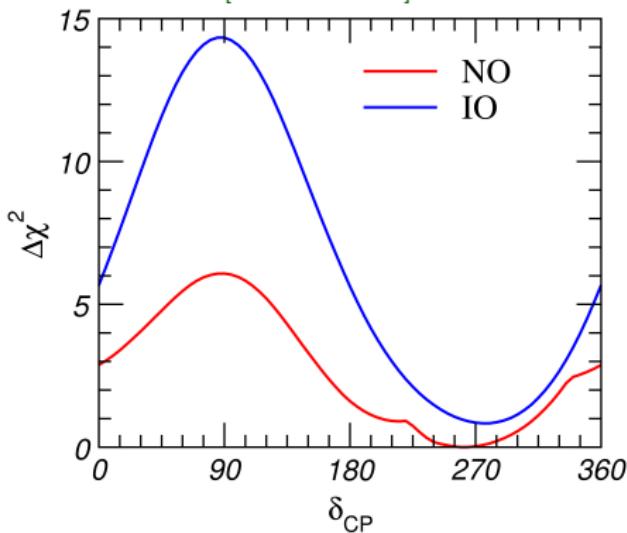


Analyzable subset of SK atmospheric data

$\delta_{13}$

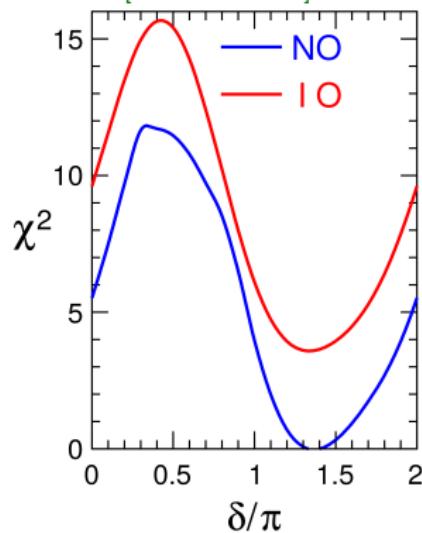
NuFIT 3.0 (2016)

[arXiv:1611.01514]



Bari 2017

[arXiv:1703.04471]

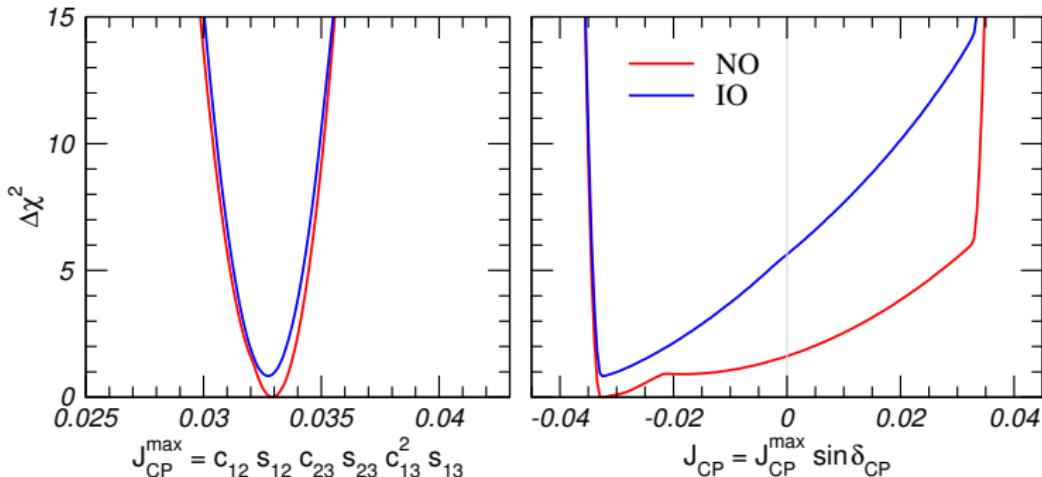


$$\frac{\delta_{13}}{\pi} = \begin{cases} 1.45^{+0.28}_{-0.33} & (\text{NO}) \\ 1.54^{+0.22}_{-0.26} & (\text{IO}) \end{cases}$$

$$\frac{\delta_{13}}{\pi} = \begin{cases} 1.38^{+0.23}_{-0.20} & (\text{NO}) \\ 1.31^{+0.31}_{-0.19} & (\text{IO}) \end{cases}$$

still unknown       $\Delta\chi^2_{CPV} = 1.7$

$\sim 20\%$  precision       $\Delta\chi^2_{CPV} \simeq 4$

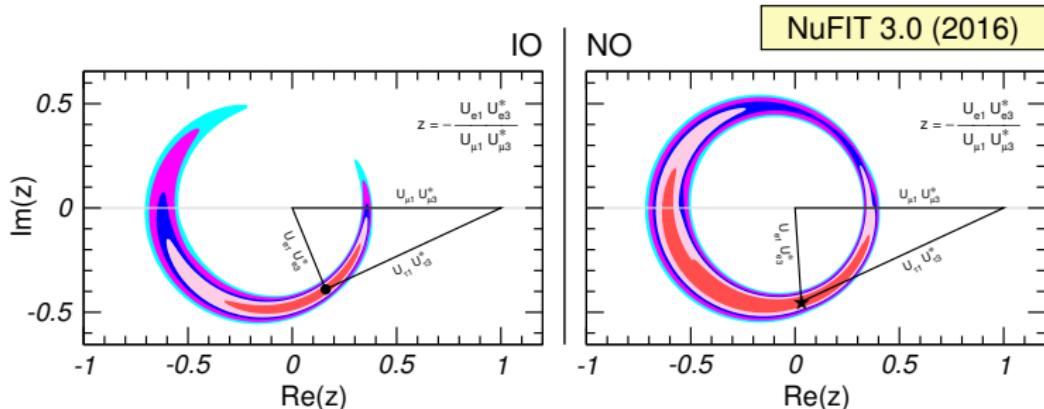


$$J_{CP}^{\max} = 0.0329 \pm 0.0007 \quad (+0.0021) \quad (-0.0024)$$

$$J_{CP}^{bf} = -0.033$$

About  $10^3$  larger than  $J_{CP}^{\text{quarks}} = (3.04^{+0.21}_{-0.20}) \times 10^{-5}$

# Unitarity Triangle

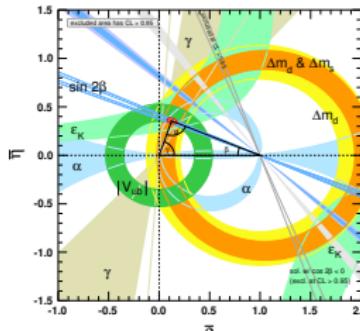


$$z = - \frac{U_{e1} U_{e3}^* U_{\mu 1}^* U_{\mu 3}}{|U_{\mu 1}|^2 |U_{\mu 3}|^2} \implies \text{Im}(z) = \frac{J_{CP}}{|U_{\mu 1}|^2 |U_{\mu 3}|^2}$$

Regions defined with respect to the global minimum (NO)

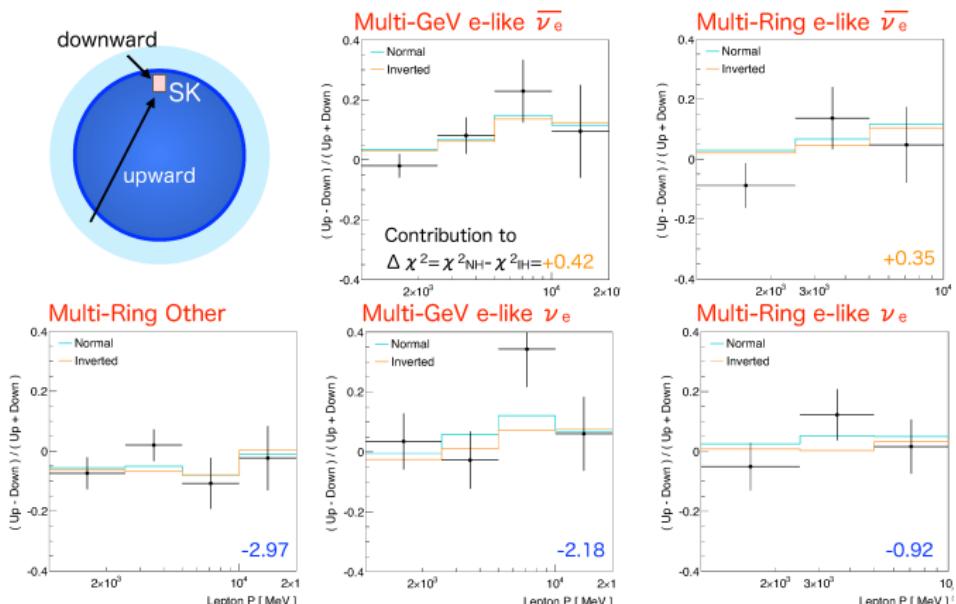
Quark sector:

[PDG 2016]



# Mass Ordering

- NuFIT 3.0 (2016):  $\Delta\chi^2_{\text{IO-NO}} \simeq 1$  Without SK atmospheric data
- Bari 2017:
  - $\Delta\chi^2_{\text{IO-NO}} = 3.6$  ( $\sim 2\sigma$ ) With analyzable subset of SK atmospheric data
  - $\Delta\chi^2_{\text{IO-NO}} = 1.1$  Without SK atmospheric data
- SK atmospheric preference for NO due to excess of e-like events

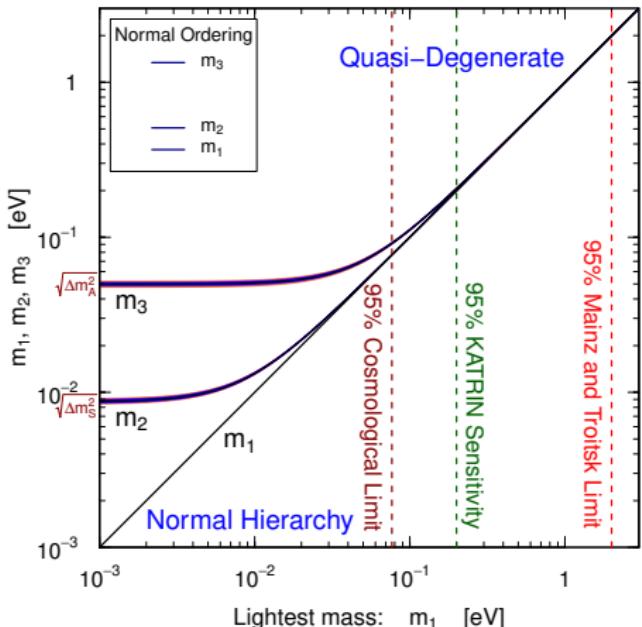


Super-Kamiokande

$$\Delta\chi^2_{\text{IO-NO}} = 5.2$$

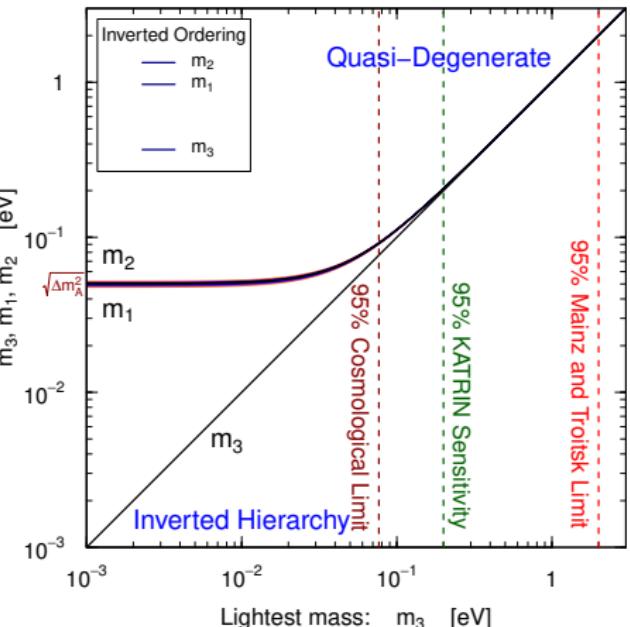
[Koshio @ NOW2016]

# Neutrino Masses



$$m_2^2 = m_1^2 + \Delta m_{21}^2 = m_1^2 + \Delta m_S^2$$

$$m_3^2 = m_1^2 + \Delta m_{31}^2 = m_1^2 + \Delta m_A^2$$



$$m_1^2 = m_3^2 - \Delta m_{31}^2 = m_3^2 + \Delta m_A^2$$

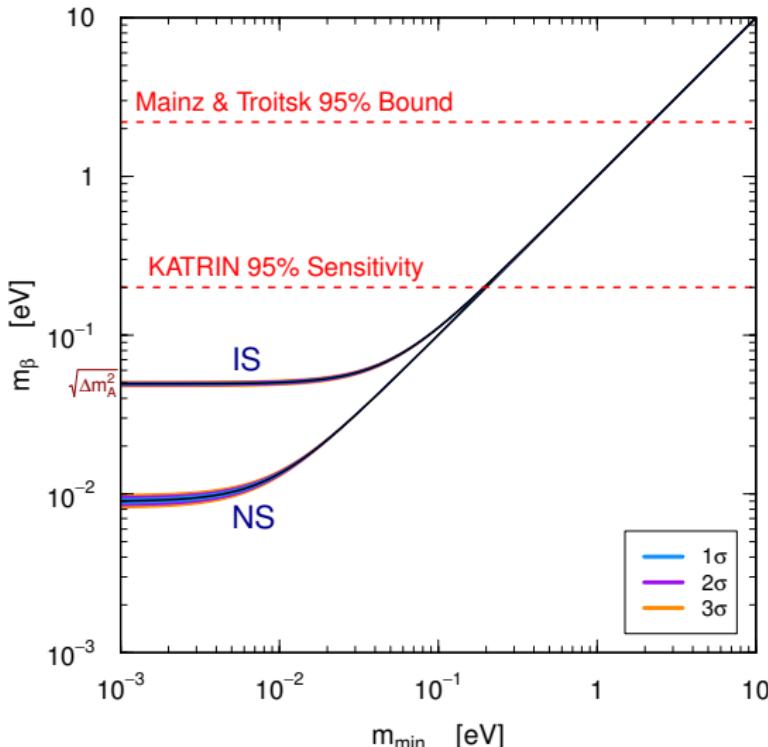
$$m_2^2 = m_1^2 + \Delta m_{21}^2 \simeq m_3^2 + \Delta m_A^2$$

Quasi-Degenerate for  $m_1 \simeq m_2 \simeq m_3 \simeq m_\nu \gtrsim \sqrt{\Delta m_A^2} \simeq 5 \times 10^{-2}$  eV

95% Cosmological Limit: Planck TT + lowP + BAO [arXiv:1502.01589]

# Beta Decay

$$m_\beta^2 = |U_{e1}|^2 m_1^2 + |U_{e2}|^2 m_2^2 + |U_{e3}|^2 m_3^2$$



► Quasi-Degenerate:

$$m_\beta^2 \simeq m_\nu^2 \sum_k |U_{ek}|^2 = m_\nu^2$$

► Inverted Hierarchy:

$$m_\beta^2 \simeq (1 - s_{13}^2) \Delta m_A^2 \simeq \Delta m_A^2$$

► Normal Hierarchy:

$$\begin{aligned} m_\beta^2 &\simeq s_{12}^2 c_{13}^2 \Delta m_S^2 + s_{13}^2 \Delta m_A^2 \\ &\simeq 2 \times 10^{-5} + 6 \times 10^{-5} \text{ eV}^2 \end{aligned}$$

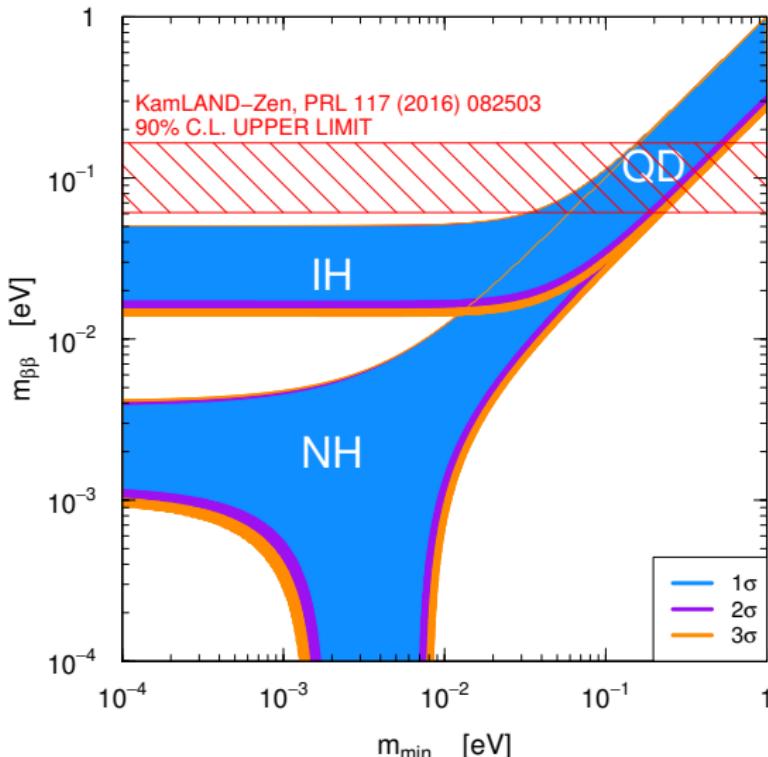
► If  $m_\beta \lesssim 4 \times 10^{-2}$  eV



Normal Spectrum

# Neutrinoless Double-Beta Decay

$$m_{\beta\beta} = |U_{e1}|^2 m_1 + |U_{e2}|^2 e^{i\alpha_2} m_2 + |U_{e3}|^2 e^{i\alpha_3} m_3$$



► Quasi-Degenerate:

$$|m_{\beta\beta}| \simeq m_\nu \sqrt{1 - s_{2\vartheta_{12}}^2 s_{\alpha_2}^2}$$

► Inverted Hierarchy:

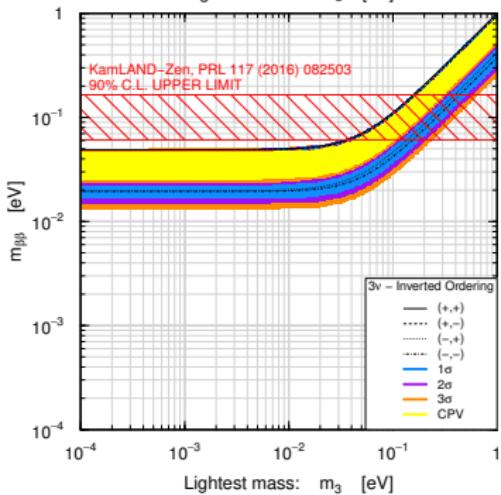
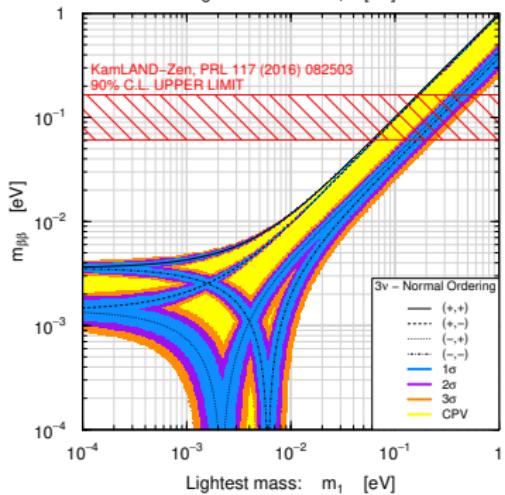
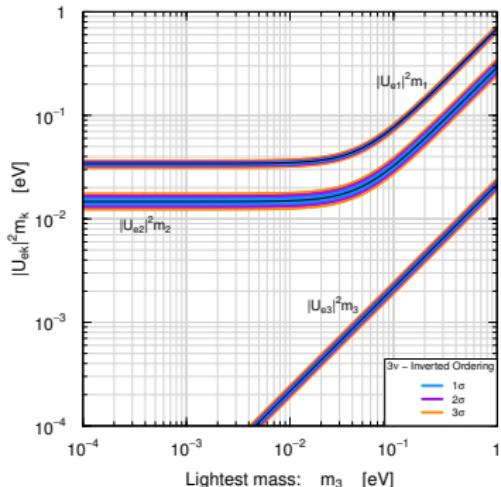
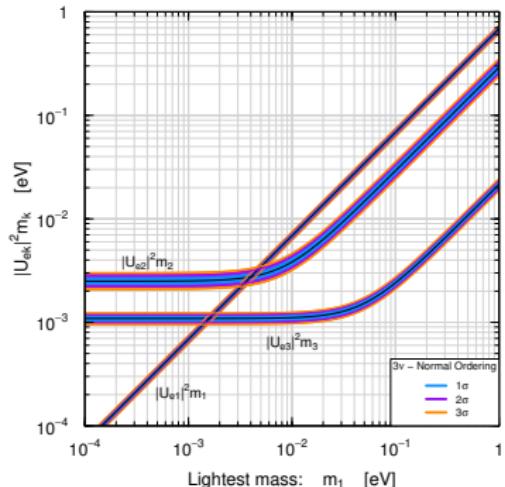
$$|m_{\beta\beta}| \simeq \sqrt{\Delta m_A^2 (1 - s_{2\vartheta_{12}}^2 s_{\alpha_2}^2)}$$

► Normal Hierarchy:

$$\begin{aligned} |m_{\beta\beta}| &\simeq |s_{12}^2 \sqrt{\Delta m_S^2} + e^{i\alpha} s_{13}^2 \sqrt{\Delta m_A^2}| \\ &\simeq |2.7 + 1.2 e^{i\alpha}| \times 10^{-3} \text{ eV} \end{aligned}$$

► If  $|m_{\beta\beta}| \lesssim 10^{-2} \text{ eV}$

Normal Spectrum



# Conclusions

- ▶ Robust  $3\nu$ -Mixing Paradigm
- ▶ Precise determination of the mixing parameters
  - ▶  $\Delta m_{21}^2 \simeq 7.4 \pm 0.2 \times 10^{-5} \text{ eV}^2$  ( $\sim 3\%$ )
  - ▶  $\sin^2 \vartheta_{12} \simeq 0.30 \pm 0.02$  ( $\sim 6\%$ )
  - ▶  $\frac{|\Delta m_{31}^2|}{10^{-3} \text{ eV}^2} \simeq \begin{cases} 2.54 \pm 0.04 & (\text{NO}) \quad (\sim 2\%) \\ 2.45 \pm 0.04 & (\text{IO}) \quad (\sim 2\%) \end{cases}$
  - ▶  $\sin^2 \vartheta_{13} \simeq 0.0216 \pm 0.0007$  ( $\sim 3\%$ )
- ▶ Open Problems:
  - ▶  $\vartheta_{23} \leq 45^\circ$  ? [T2K, NO $\nu$ A, ...]
  - ▶ CP violation ?  $\delta_{13} \approx 3\pi/2$  ? [T2K, NO $\nu$ A, DUNE, HyperK]
  - ▶ Mass Ordering ? [JUNO, RENO-50, PINGU, ORCA, INO]
  - ▶ Absolute Mass Scale ? [ $\beta$  Decay, Neutrinoless Double- $\beta$  Decay, Cosmology]
  - ▶ Dirac or Majorana ? [Neutrinoless Double- $\beta$  Decay]
  - ▶ Physics Beyond Three-Neutrino Mixing ?
  - ▶ Theory: Why lepton mixing  $\neq$  quark mixing ? Is there any connection ?
  - ▶ Why  $0 < \sin^2 \vartheta_{13} \ll \sin^2 \vartheta_{12} < \sin^2 \vartheta_{23} \simeq 0.5$  ?