Investigating Cluster Astrophysics and Cosmology with Cross-Correlation of Thermal Sunyaev-Zel'dovich Effect and Weak Lensing

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- 1. Weak Lensing, thermal Sunyaev-Zel'dovich Effect and cross-correlations
- 2. Construction of Model and Covariance
- 3. Results: Implications to Cluster Astrophysics and Cosmology
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# Weak Lensing: Cosmic Shear

'Cosmic Shear': Gravitational lensing by large-scale structure. **unbiased tracer** of dark matter.

galaxies



position kernel density



Convergence map from HSC (Oguri+, 2017)



Sunyaev and Zel'dovich (1972, 1980)

# tSZ by Planck

 The all-sky thermal Sunyaev-Zel'dovich effect has already been measured by Planck. The power spectrum of Compton-y can place a tight constraint on σ<sub>8</sub>. (c.f. Komatsu and Seljak, 2002).





### Cross-Correlation of tSZ and WL

#### Astrophysics and Cosmology with WL and tSZ

Power spectra of WL and tSZ give us information of the large-scale structure in the Universe and are very useful to constrain cosmological models. However, the combination of them (cross-correlation) provides us with additional and independent information.

Furthermore, tSZ directly reflects gas distribution. We can obtain implications on **cluster astrophysics** as well.

#### From the observational aspect

Cross-correlation does not suffer from noise auto-correlation in the assumption that noises of different observables are uncorrelated.

$$\langle A_{\rm obs} B_{\rm obs} \rangle = \langle AB \rangle, \ A_{\rm obs} = A + N_A, \ B_{\rm obs} = B + N_B$$

### Measurements of Cross-Correlation

Several groups have already reported the detection of the cross-correlation.

RCSLenS x Planck CFHTLenS x Planck (Van Waerbeke+, 2014) (Hojjati+, 2016) 0.T0\_ 5·10<sup>-9</sup> Planck A  $imes 10^{-9}$ 4·10<sup>-9</sup> 3.53·10<sup>-9</sup> Planck 130 significance  $\xi_{yk}(r)$ 3.02·10<sup>-9</sup>  $1.10^{-9}$ 2.5 $0.10^{0}$ **2PCFs** -1·10<sup>-9</sup> 2.0 $\xi^{y-\kappa}(\vartheta)$ -2·10<sup>-9</sup> PlanckJ 1.54·10<sup>-9</sup> 1.03·10<sup>-9</sup> 0.5 $\xi_{yk}(r)$ 2·10<sup>-9</sup>  $1.10^{-9}$ 0.0 $0.10^{0}$ -0.5204060 80 100 120 140 160 180 -1·10<sup>-9</sup> 30 40 50 60 0 10 20 separation [arcmin] separation [arcmin]

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# Models

First, let us consider the model of the signal.

Theoretical prediction of auto- and cross-power spectra

Spectra can be decomposed into two terms based on halo model.

$$\begin{split} C_{\ell}^{y\kappa} &= C_{\ell}^{y\kappa(1\mathrm{h})} + C_{\ell}^{y\kappa(2\mathrm{h})}, & \text{well calibrated by N-body sim.} \\ C_{\ell}^{y\kappa(1\mathrm{h})} &= \int_{0}^{z_{\mathrm{dec}}} dz \, \frac{d^2 V}{dz d\Omega} \int_{M_{\mathrm{min}}}^{M_{\mathrm{max}}} dM \, \frac{dn}{dM} y_{\ell}(M, z) \kappa_{\ell}(M, z), \end{split}$$
 $C_{\ell}^{y\kappa(2h)} = \int_{0}^{z_{dec}} dz \, \frac{d^2V}{dzd\Omega} P_{\rm m}(k = \ell/d_A, z) \quad \begin{array}{l} \mbox{Fourier transform of} \\ \mbox{Compton-y of halo} \end{array}$ **Compton-y of halo**  $\times \int_{M_{\min}}^{m_{\max}} dM \, \frac{dn}{dM} b(M,z) \kappa_{\ell}(M,z) \quad \begin{array}{l} \clubsuit \text{We need pressure} \\ \text{profile of halo.} \end{array}$  $\times \int_{M}^{M} dM \frac{dn}{dM} b(M,z) y_{\ell}(M,z).$ 

## Analytic Model of Gas Profiles

 We employ the analytic gas density/pressure profile model of individual halo, which is proposed by Shaw+ (2010) and improved by Flender+ (2017). The model contain six free parameters, each of which describes a physical process (e.g., SNe/AGN feedback, non-thermal pressure).

#### Basic ideas

DM density follows NFW profile.  $\rho_{\rm DM} = \frac{\rho_s}{(r/r_s)(1+r/r_s)^2}$ Gas profile is determined from Euler eq. with polytope relation.  $\frac{dP_{\rm tot}}{dr} = -\rho_g(r)\frac{d\Phi(r)}{dr}, \ P_{\rm tot} \propto \rho_g^{\Gamma}$  hydro. sim. suggests  $\Gamma$ ~1.2 To determine the normalization, stellar and AGN feedback  $E_{g,f} = E_{g,i} + \varepsilon_{\rm DM} |E_{\rm DM}| + \varepsilon_{\rm f} M_* c^2 + \Delta E_p$ Energy of gas dynamical friction between gas and DM work by gas expansion

# Gas Profile Model

#### Non-thermal pressure

Turbulent motion also can support the self-gravity of the halo.

This effect is parametrized as,

$$\frac{P_{\rm nth}}{P_{\rm tot}}(r) = \alpha (1+z)^{\beta} \left(\frac{r}{R_{500}}\right)^{0.8}$$

Note: only thermal pressure contributes as tSZ.

 Free parameters are calibrated by gas density and gas fraction of X-ray clusters. We fix parameters other than alpha.



# Covariance Estimation

#### Let us move on covariance matrix estimation.

#### Simulations

In order to estimate covariance matrix, we employ **N-body simulations.** Box size: (1 Gpc/h)<sup>3</sup> # of particles: 2048<sup>3</sup> Cosmology: Planck 2015 **The snapshots cover up to z=4.13**.



#### Mock observations

WL: Multiple plane method (White & Hu, 2000)

tSZ: First, we find halos with Rockstar (Behroozi+, 2013). For each halo, we solve the analytic model, and assign thermal pressure to the member particles.

## Covariance Estimation



#### Mock observations

- WL: Multiple plane method (White & Hu, 2000)
- tSZ: First, we find halos with Rockstar (Behroozi+, 2013). For each halo, we solve the analytic model, and assign thermal pressure to the member particles.

### Covariance Estimation



- WL: Multiple plane method (White & Hu, 2000)
- tSZ: First, we find halos with Rockstar (Bohroozi+, 2013).

For each halo, we solve the analytic model, and assign the covariance matrix thermal pressure to the member particles.

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Compton-y Spectra



## **Cross-Correlation Function**



### Constraints on Non-thermal Pressure

#### Non-thermal pressure

Non-thermal contribution is hard to measure by conventional X-ray observations of clusters.

However, the power spectrum and the cross-correlation

are sensitive to it and can be used to estimate its contribution.

Note: Non-thermal pressure and  $\sigma_8$  are strongly degenerate.



### **Constraints from Cross-Correlation**

• We can constrain the amplitude of non-thermal pressure and  $\sigma_8$  with power spectrum and cross-correlation.



### Removing Small Scales Correlations

- 1. Halo model assumption breaks down.
- 2. Incomplete separation from foreground contamination



# Source of Tension

• The analytic model is calibrated with **galaxy clusters** and **low-z galaxy groups**. However, **high-z groups** contribute to a substantial fraction of signal. At this range, analytic model inevitably contains uncertainty.

• Contribution from high-z groups (z > 0.2 and  $M_{500c} < 4 \times 10^{14} M_{sun}/h$ )



# Possible Solution of Tension

 Let us consider an extreme case of enhanced star formation for group size halos.

That corresponds to severe depletion of hot gas for such halos.



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# Summary

- •Weak lensing and thermal Sunyaev-Zel'dovich effect are promising probes into the large-scale structure.
- •Theoretical modeling and simulations can be used to estimate the signal and covariance matrix.
- •The cross-correlation of them can provide us with additional information of cosmology and cluster astrophysics.
- •We are currently working on the measurement of the cross-correlation with HSC and Planck data.