22/08/16

To CCQE and Beyond

Results and prospects of latest CCQE-like analyses from the T2K near detectors

Stephen Dolan For the T2K Collaboration

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Overview

- The T2K experiment
- Motivation for measuring CCQE-like cross sections
- Flux and near detectors' details
- NuFact 15 recap
- CC0 π on water cross section using the PØD
- FGD1 analyses using proton information
 - CC0 π using proton kinematics
 - CC0 π using inferred kinematic imbalance
 - CC0 π using transverse kinematic imbalance
- Other analyses
 - CC0 π measurements from $\nu + \bar{\nu}$ joint fit
 - CC0 π at INGRID
 - Extraction of free nucleon cross section using δp_{TT}
- Summary and future work





The T2K Experiment



Far Detector

Super-Kamiokande





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Data Collection



- Continuous rise in beam power from ~225 kW (2014) to ~420 kW (2016)
- Using this to make world leading measurements of oscillation parameters (see talk by Benjamin Quilain – WG1, Tuesday, 14:00)

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Neutrino Scattering and OA

- Oscillation analysis (OA) requires E_{ν} spectrum (or similar)
- Can reconstruct using observed μ assuming stationary target and elastic scattering

$$E_{\nu,rec} = \frac{m_p^2 - m_n^2 - m_\mu^2 + 2m_n E_\mu}{2(m_n - E_\mu + p_\mu \cos(\theta_\mu))}$$

Bias due to Fermi Moton and CCnonQE components





Neutrino Scattering and OA

- Essential to understand νN scattering to control the bias
 - CCQE particularly important for T2K
- Probe using $CC0\pi$ cross sections
 - Less FSI model dependence
 - Simplest channel to probe nuclear effects.







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- Off-axis v_{μ} beam
 - Tightly-peaked at 600 MeV 2.5° off-axis towards SK
 - Low contamination from non- ν_{μ} components
 - Flux estimation aided by hadron production measurements from NA61/SHINE at CERN (see talk at WG1+2 session on Thursday 10:45)

Phys. Rev. D 87, 012001





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From Reconstruction to Truth



- Measure selected number of CC0π events in bins of a reconstructed quantity
- Need the **total** number of $CC0\pi$ events in bins of a **true** quantity

Two Methods



Matrix Unfold



- Use MC to build unsmearing matrix
- Apply unsmearing matrix + efficiency correction to data
- True bin \rightarrow Reco Template
- Vary MC template norm to fit data
- Apply efficiency correction

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Previously at NuFact

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INGRID On-Axis CCQE Result

- Carbon Target
- 2-bin measurement in neutrino energy
- Split into 1- and 2-track samples
- Used cuts on muon and proton kinematic variables to enhance purity
- Result depends on nuclear model used and the presence of 2p2h





ND280 Off-Axis CC0 π Result

- Uses FGD1 as a CH target alongside TPC for tracking
- Flux integrated doubledifferential **CC0** π cross section in final state muon kinematic variables $(p_{\mu}, \cos(\theta_{\mu}))$
- Split into two analyses with different selection and crosssection extraction strategies
 - Good agreement
- Results compared to 2p2h models



Detector: ND280 – FGD1 **Target:** Carbon **Signal:** $CC0\pi$ **Unfolding:** Matrix + Fit **Status:** Phys. Rev. D **93**, 112012

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ND280 Off-Axis CC0 π Result

- Results compared to Martini et al. model with(red)/without(black) 2p2h
- Data prefer a 2p2h contribution



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What next?



Interactions

Final State Interactions

Pauli

- Would like to disentangle the role of separate nuclear effects
- Current results provide an important piece of the puzzle

One cross-section measurement

Now need complementary measurements ...



Ongoing measurements New PØD $CC0\pi$ water cross section in muon kinematics analysis - Measure of A-scaling, invaluable for OA $CC0\pi$ measurement using muon + proton kinematics - Enhanced sensitivity to nuclear effects FGD1 ongoing analyses $CC0\pi$ measurement using composite variables using proton - Imbalance between the proton and muon can be information a precision probe of nuclear effects $CC0\pi$ using INGRID proton module Model-independent measurement at higher E_{ν} Other FGD1 / $CC0\pi$ neutrino/anti-neutrino joint fit INGRID - $\sigma_{np-nh}/\sigma_{CCOE}(E_{\nu})$ is substantially different for ν_{μ} and $\bar{\nu}_{\mu}$ analyses Measurement of free-nucleon cross section using δp_{TT}

N.B: T2K employs a blind cross-section analysis strategy

- Ongoing or recently completed analyses not applied to real data (remain "blind")

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$CC0\pi$ water cross section

Contact: Tianlu Yuan tianlu.vuan@colorado.edu

- Isolate CC0 π events starting in the PØD
- Separate data taking periods into when PØD water target is full/empty \rightarrow subtract to get water cross section

0.6

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Target: Water

Event Selection

- Uses PØD as a target, requires TPC for tracking
- Aim to find single μ only
- Two control samples
 - CC1 π : look for 2 PØD tracks and Michel e^-
 - **CCOther**: look for >2 PØD tracks



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Detector: ND280 - PØD

$CC0\pi$ water cross section

Contact: Tianlu Yuan **tianlu.yuan@colorado.edu**

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- Construct flux integrated doubledifferential cross section
- Results compared to GENIE and NEUT predictions
- Can also compare to FGD1 CC0π on Carbon result

 Similar studies underway using FGD2 water layers to extract Oxygen:Carbon cross section ratio



Detector: ND280 - PØD



$CC0\pi$ water cross section

Contact: **Tianlu Yuan** tianlu.yuan@colorado.edu

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TZ

- Compare results to RPA/RPA+2p2h on Carbon
- Data prefer 2p2h contribution
- Difficult to untangle role of A-Scaling





FGD1 - CC0 π + Np

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FGD1 CC0*π* Analyses



- Require one μ -like track or one μ -like and p-like track(s) starting in FGD1
- Use a Michel electron tag and ECal EM shower veto to reject 1π backgrounds
- Use of many samples gives wide kinematic acceptance

Sidebands

 Require extra π-like track(s)





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$CC0\pi$ using $\mu + p$ kinematics



• Measuring *p* kinematics allows us to move beyond these assumptions

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0.6 0.8

 $cos(\theta_{p}^{true})$

1.0

-1.0 -0.8 -0.6 -0.4 -0.2 0.0 0.2 0.4

$CC0\pi$ using $\mu + p$ kinematics

- Aim to measure $CC0\pi$ cross section in bins of $\cos(\theta_{\mu}), \cos(\theta_{p}), p_{p}$ in samples with proton
- Use $\cos(\theta_{\mu})$, p_{μ} when no proton reconstructed
- Also measure proton multiplicity
- Construct flux-integrated double/triple-differential cross-section
- Fake data: GENIE*
- Nominal MC: NEUT

Detector: ND280 - FGD1

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Contact:

Pierre Bartet

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 $CC0\pi$ and inferred kinematic imbalance

- Can use proton and muon kinematics to form variables specifically engineered to probe nuclear effects
- Under stationary target and elastic scattering assumptions can infer proton kinematics from measured μ
- Non-zero imbalance between inference and measured proton indicates presence of nuclear effects or CC-non-QE interaction
- Measure:



$CC0\pi$ and inferred kinematic imbalance

Contact: Jiae Kim

flux

FSI

······ MC Truth

xsec detector

data stat MC stat

Unfolded

Fakedata truth

 $0.8 < \cos \theta_{\mu} < 1.0, p_{..} > 750 \text{ MeV}$

T2K Work

In Progress

jiae@phas.ubc.ca



- Measure inferred kinematics in bins of p_{μ} , $\cos(\theta_{\mu})$
- Fake data: GENIE*
- Nominal MC: NEUT



 $(10^{-38}cm^2GeV^{-1}Nucleon^{-1})$

0.06

0.05

0.04

0.03

0.02

Single Transverse Variables



Detector: ND280 – FGD1Target: CarbonSignal: CC0π+NpUnfolding: FitStatus: BlindStephen DolanNuFact 2016, Quy Nhon, Vietnam27T2K







$CC0\pi$ and transverse imbalance Contact: Stephen Dolan Stephen Dolan Stephen Dolan

- 3 single transverse variables (STV) characterise imbalance in plane transverse to incoming v^*
- For CCQE case any deviation from $\delta p_T = 0$, $\delta \phi_T = 0$ is indicative of nuclear effects

• Minimal dependence on E_{ν} for δp_T and $\delta \alpha_T$



 δp_T

δατ

$CC0\pi$ in STV - Fermi Motion and FSI

Moving from CCQE→CC0Pi+Np, STV still a probe of nuclear effects



Detector: ND280 – FGD1	Target: Carbon	Signal: CC0π+Np	Unfolding: Fit	Status: Blind
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$CC0\pi$ in STV - 2p2h and M_A

Target: Carbon

M. Martini, M. Ericson, G. Chanfray, and J. Marteau, Phys. Rev. C 80, 065501 (2009)

Detector: ND280 - FGD1

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J. Nieves, I. R. Simo, and M. J. V. Vacas, Phys. Rev. C 83, 045501 (2011)



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Signal: $CC0\pi + Np$

Unfolding: Fit

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Status: Blind

$CC0\pi$ in STV

Contact: Stephen Dolan s.dolan@physics.ox.ac.uk

 $p_{\mu} > 250 \; MeV/c$

 $\cos(\theta_u) > -0.6$

 $\cos(\theta_p) > 0.4$

 $450 \ MeV/c < p_{\mu} < 1 \ GeV/c$

Restrict cross section to ND280 acceptance —

- Use a regularised template fit to unfold
 - Useful to deal with large STV smearing
 - Regularisation insists cross-sections should be smooth
- Fake data: GENIE*
- Nominal MC: NEUT



FGD - CC0 π + Np

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Summary

- T2K is measuring cross-sections of exclusive final-state topologies - This talk: $CC0\pi$
 - Talk by Erez Reinherz-Aronis: CCInc, CC1 π +, NCE (WG2, Friday, 10:45)
- Lots going on in T2K cross-section analyses many results coming soon!
- First CC0 π water cross section has been measured
- Many new techniques in use to complement each other and existing results
 - Analyses specifically engineered to probe nuclear effects



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The Future



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Thank you for listening



Cảm ơn

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Ongoing measurements

•	CCOπ water cu - Measure of A	New PØD					
	Detector: ND280 – PØD	Target: Water	Signal: CC0 π	Unfolding: Matrix	Status: Unblind		
•	$CC0\pi$ measure						
	- Enhanced ser						
	Detector: ND280 – FGD1	Target: Carbon	Signal: CC0 <i>π</i> (+N	p) Unfolding: Fit	Status: Blind		
•	CC0π measure - Imbalance as	Ongoing FGD1 analyses					
[Detector: ND280 – FGD1	Target: Carbon	Signal: CC0 <i>π</i> +Np	Unfolding : Matrix	Status: Blind	information	
•	- Imbalance as						
	Detector: ND280 – FGD1	Target: Carbon	Signal: CC0 <i>π</i> +Np	Unfolding: Fit	Status: Blind		
•	CC0π using IN - Model indep						
	Detector: ND280 - FGD1	Target: Carbon	Signal: CC0 π	Unfolding: Fit	Status: Blind		
•	$CC0\pi$ neutrino	Other FGD1 /					
	- $\sigma_{np-nh}/\sigma_{CCQE}$	- INGRID					
[Detector: INGRID	Target: Carbon	Signal: CC0 π	Unfolding : Matrix	Status: Blind	analyses	
•	Measurement	leasurement of free-nucleon cross-section using δp_{TT}					
[Detector: ND280 – FGD1	Target: Hydrogen	Signal: CC1 <i>π</i> 1p	Unfolding: Fit	Status: Blind		

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BACKUPS

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Likelihood Fitting



- Each true bin has some template in reconstructed bins.
- Varying the number of events in a true bin applies a normalisation factor to the corresponding recon template.
- Vary templates until reconstructed distribution best fits the data

• I.e minimise:
$$\chi^2_{stat} = \sum_{j}^{recobins} 2(N_j^{MC} - N_j^{obs} + N_j^{obs} ln \frac{N_j^{obs}}{N_j^{MC}})$$

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Fitting summary

• The best fit parameters are those that minimise the following likelihood:

$$\chi^2 = \chi^2_{stat(fit\,goodness)} + \chi^2_{syst(penalty)} + \chi^2_{reg}.$$

$$\chi_{stat}^2 = \sum_{j}^{recobins} 2(N_j^{MC} - N_j^{obs} + N_j^{obs} ln \frac{N_j^{obs}}{N_j^{MC}})$$

$$\chi^2_{syst} = (\vec{a}^{syst} - \vec{a}^{syst}_{prior})(V^{syst}_{cov})^{-1}(\vec{a}^{syst} - \vec{a}^{syst}_{prior})$$

$$\chi^2_{reg} = p_{reg} \sum_i (c_i - c_{i-1})^2$$



Systematics in the fitter

• Add term to the fit

$$\chi^2_{syst} = (\vec{a}^{syst} - \vec{a}^{syst}_{prior})(V^{syst}_{cov})^{-1}(\vec{a}^{syst} - \vec{a}^{syst}_{prior})$$

- The fit is able to constrain systematic parameters (mostly through control regions) but picks up a penalty if it moves far from the prior.
- **Detector Systematics** (e.g. TPC momentum resolution):
 - Make many toy experiments, each varying detector properties we are unsure of.
 - Produce covariance matrix which tells the fit the overall uncertainty in the number of events in each bin and how this correlates between bins.
- **Model Systematics** (e.g. *M_{A,RES}*, pion FSI):
 - Use covariance matrix produced from external data fits which tells us the uncertainty on model parameters.
 - Make splines that tell the fitter how to reweight the MC if we alter model parameters that describe the background In the fit.
 - Flux Systematics:
 - Use covariance matrix produced by beam group which tells us the uncertainty on the flux in bins of neutrino energy.
 - Store the neutrino energy of each event in the fit.



Regularising the fitter

- Reaching a fit result is an 'ill posed problem' -> there are often degeneracies in the fit solutions.
- E.g. can often lower a particular template scaling parameter so long as we raise the adjacent parameters
 - Strong anti correlations between bins!
- Can resolve with regularisation:
 - Add another term
 - Regularisation loosely ties bins together
- But how to choose the best p_{reg} ?
- Want to have the maximum smoothing impact with the minimal effect on the χ^2 of the fit.
 - This problem is well studied, can choose p_{reg} using the "L-Curve".
- Toy example fitting a Gaussian from a flat prior in backups.

$$\chi^2_{reg} = p_{reg} \sum_i (c_i - c_{i-1})^2$$



Regularising the fitter







SIAM Rev., 34(4), 561580 (1992).

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Regularisation

More detail:

https://www.mat.tuhh.de/lehre/material/RegLS.pdf

• Penalty term in my fit from regularisation:

$$\chi^2_{reg} = p_{reg} \sum_i (c_i - c_{i-1})^2$$

This looks like a rather non χ^2 -like term slapped onto the fit...

But could also write it as:
$$\chi^2_{reg} = p_{reg} (p - p_{prior}) (V_{cov})^{-1} (p - p_{prior})$$

$$(V_{cov})^{-1} = \begin{bmatrix} 1 & -1 & 0 & \cdots & 0 \\ -1 & 2 & -1 & \dots & 0 \\ 0 & -1 & 2 & \ddots & 0 \\ \vdots & \vdots & \ddots & \ddots & -1 \\ 0 & 0 & 0 & -1 & 1 \end{bmatrix}$$
 Or to make the matrix non-singular:
$$(V_{cov})^{-1} = \begin{bmatrix} 1 & -1 & 0 & \cdots & 0 \\ -1 & 2 & -1 & \dots & 0 \\ 0 & -1 & 2 & \ddots & 0 \\ \vdots & \vdots & \ddots & \ddots & -1 \\ 0 & 0 & 0 & -1 & 2 \end{bmatrix}$$

- Applying a penalty term in this way makes regularisation enter the fit identically to model parameters.
- In fact in some sense the application of regularisation is a model that says cross sections should be smooth relative to their prior.
- The uncertainty in the smoothness model is then $1/\sqrt{p_{reg}}$.





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T2K



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FGD2 Oxygen Cross Section

Contact: Margherita Buizza Avanzini **buizza@llr.in2p3.fr**

• Measure the cross section on **oxygen** and **carbon** simultaneously fitting events starting in FG2 water and carbon layers



 Extract flux integrated Oxygen cross section and ratio oxygen/carbon flux integrated cross sections

$$\frac{d\sigma}{dp_{\mu}} = \frac{o_i N_{MC}^{sig,O}}{\Phi \cdot N_{neutrons}^{FV,O} \cdot \Delta p_{\mu}}$$





INGRID CCOPi Analysis



- Select 1 μ -like track beginning in the proton module
- Measure θ_{μ}^{rec} and d_{μ}^{rec} distance penetrated through the iron (no B field)
- Unfold into θ_{μ}^{true} and p_{μ}^{true}
- Build double differential cross-section
- With 4 momentum bins and 5 angular bins uncertainty is 10%-20%
- Blind analysis



CCOPi $v + \bar{v}$ joint fit

- 2p2h contribution is different for v and \bar{v} *
- Comparison of $v + \bar{v}$ can help identify 2p2h
- Aim to extract $\bar{\nu}$ double differential crosssection in p_{μ} , $\cos(\theta_{\mu})$ alongside $\nu + \bar{\nu}$ sum, difference and asymmetry
- Uses extra high angle and backward $\mu^{+/-}$ samples
- Cross-section extraction via a likelihood template fit
- Blind analysis



* M Martini: PHYSICAL REVIEW C 80, 065501 (2009), PHYSICAL REVIEW C 81, 045502 (2010)

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3.24x10^20 POT



CCOPi $v + \bar{v}$ joint fit measurements

$$c_i^{\bar{v}} N_i^{\bar{v}MCCC0\pi}, c_i^{v} N_i^{vMCCC0\pi}$$

= (what we measure) number of $CC0\pi$ events in 'true' muon p, $\cos\theta$ bins

Extract 3 measurements:

• CC0 $\pi \bar{v}$ flux integrated cross-section

$$\frac{d\,\sigma^{\nabla}}{dp_{\mu}d\cos\theta_{\mu}} = \frac{c_{i}^{\nabla}N_{i}^{\nabla MCCC0\pi}}{\Phi^{\nabla} \cdot N_{protons}^{FV} \cdot \Delta p_{\mu}\Delta\cos\theta_{\mu}}$$

 sum, difference and v-v xsec → allow to disentangle different terms of xsec (compare with 2p2h models)

$$\frac{d\left(\sigma^{\mathsf{v}}\pm\sigma^{\overline{\mathsf{v}}}\right)}{dp_{\mu}d\cos\theta_{\mu}} = \frac{1}{\Delta p_{\mu}\Delta\cos\theta_{\mu}} \left[\frac{c_{i}^{\overline{\mathsf{v}}}N_{i}^{\overline{\mathsf{v}}MCCC0\pi}}{\Phi^{\overline{\mathsf{v}}}\cdot N_{protons}^{FV}} \pm \frac{c_{i}^{\mathsf{v}}N_{i}^{\mathsf{v}MCCC0\pi}}{\Phi^{\mathsf{v}}\cdot N_{neutrons}^{FV}} \right]$$

- asymmetry of v- \bar{v} xsec \rightarrow direct effect on δ_{CP} measurement

$$\frac{d\left(\sigma^{\mathsf{v}}-\sigma^{\bar{\mathsf{v}}}\right)}{d\left(\sigma^{\mathsf{v}}+\sigma^{\bar{\mathsf{v}}}\right)} = \frac{c_{i}^{\bar{\mathsf{v}}} N_{i}^{\bar{\mathsf{v}} MC CC0 \pi} / (\Phi^{\bar{\mathsf{v}}} \cdot N_{protons}^{FV}) - c_{i}^{\mathsf{v}} N_{i}^{\bar{\mathsf{v}} MC CC0 \pi} / (\Phi^{\bar{\mathsf{v}}} \cdot N_{neutrons}^{FV})}{c_{i}^{\bar{\mathsf{v}}} N_{i}^{\bar{\mathsf{v}} MC CC0 \pi} / (\Phi^{\bar{\mathsf{v}}} \cdot N_{protons}^{FV}) + c_{i}^{\mathsf{v}} N_{i}^{\bar{\mathsf{v}} MC CC0 \pi} / (\Phi^{\bar{\mathsf{v}}} \cdot N_{neutrons}^{FV})}$$



δp_{TT}



FIG. 1. Schematic illustration of the double-transverse kinematics. The incoming and outgoing particle momenta are represented by \vec{p}_{ν} and \vec{p}_{l} , $\vec{p}_{\rm p}$ and \vec{p}_{π} , respectively. The double-transverse momentum imbalance, $\delta p_{\rm TT}$ is given by $p_{\rm TT}^{\rm p} + p_{\rm TT}^{\pi}$ with respect to the axis $\vec{z}_{\rm TT}$ defined by $\vec{p}_{\nu} \times \vec{p}_{l}$.



Phys. Rev. D 92, 051302(R)

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$CC1p1\pi$ on hydrogen

- Cross-section measurement on H allows measure of $\sigma(v + p)$ without nuclear effects
- Can use δp_{TT} to isolate H content of composite target *



• Aim to measure Δ^{++} production on H

* Phys. Rev. D 92, 051302(R)

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Reconstructing the Neutrino Direction



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Pauli Blocking



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T2/K





0.5



1.5

2.5

 E_{ν} (GeV)

Interaction Modes in selected 1 ring μ -like events at SuperK(NEUT):

0^L0

0.5

1.5

2

 $\begin{array}{c} 2.5 \\ E_{\rm v} \, ({\rm GeV}) \end{array}$



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ND280 Off-Axis CC0 π Result



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