

Low-scale baryogenesis

Takehiko Asaka (Niigata Univ.)

**Xth Rencontres du Vietnam
Flavor Physics Conference**

**ICISE, QUY NHON, VN
(27 July-2 August 2014)**

Baryon asymmetry of the universe (BAU)

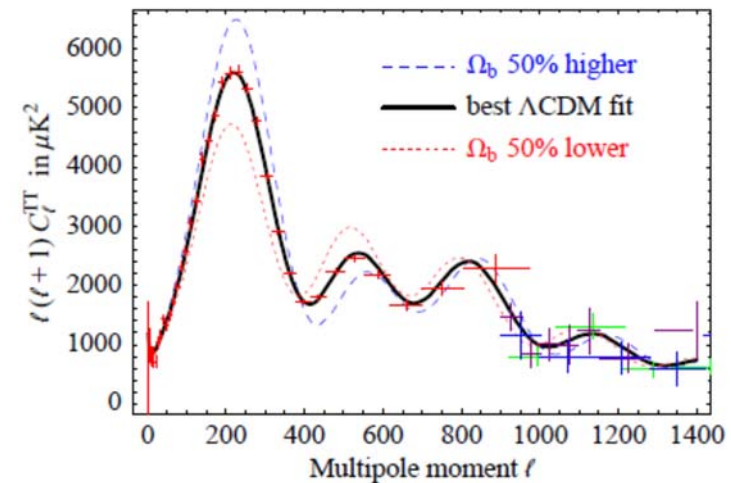
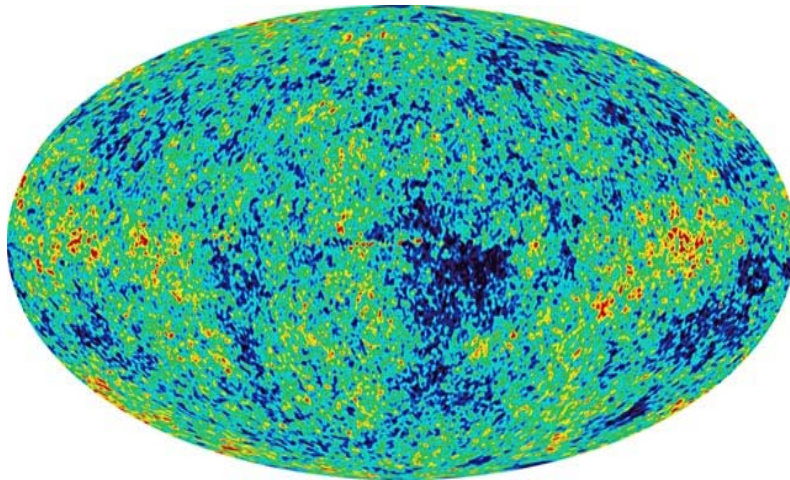
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Baryon Number $B = (\# \text{ of baryons}) - (\# \text{ of antibaryons})$

$$\frac{n_B}{s} = (8.579 \pm 0.109) \times 10^{-11}$$

Planck 2013 results
arXiv:1303.5076

n_B : Baryon number density
 s : Entropy density



[Strumia 06]

- **Neutrino mass scales**

- ▣ Atmospheric: $\Delta m_{\text{atm}}^2 \simeq 2.4 \times 10^{-3} \text{eV}^2$

- ▣ Solar : $\Delta m_{\text{solar}}^2 \simeq 7.5 \times 10^{-5} \text{eV}^2$

⇒ **Need for physics beyond the SM !**

- Important questions:

- ▣ ***“What is the origin of neutrino masses?”***

- ▣ ***“What is the implications of it?”***

- ▣ ***“How do we test it experimentally?”***

Extension by RH neutrinos ν_R

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$$\delta L = i\bar{\nu}_R \partial_\mu \gamma^\mu \nu_R - F \bar{L} \nu_R \Phi - \frac{M_M}{2} \bar{\nu}_R \nu_R^c + \text{h.c.}$$

Minkowski '77

Yanagida '79

Gell-Mann, Ramond, Slansky '79

Glashow '79

- Seesaw mechanism ($M_D = F\langle\Phi\rangle \ll M_M$)

$$-L = \frac{1}{2} (\bar{\nu}_L, \bar{\nu}_R^c) \begin{pmatrix} 0 & M_D \\ M_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} + \text{h.c.} = \frac{1}{2} (\bar{\nu}, \bar{N}^c) \begin{pmatrix} M_\nu & 0 \\ 0 & M_M \end{pmatrix} \begin{pmatrix} \nu^c \\ N \end{pmatrix} + \text{h.c.}$$

$$M_\nu = -M_D^T \frac{1}{M_M} M_D$$

$$U^T M_\nu U = \text{diag}(m_1, m_2, m_3)$$

□ Light active neutrinos ν_1, ν_2, ν_3

→ explain neutrino oscillations

□ Heavy neutral leptons N_1, N_2, N_3 ($N \simeq \nu_R$)

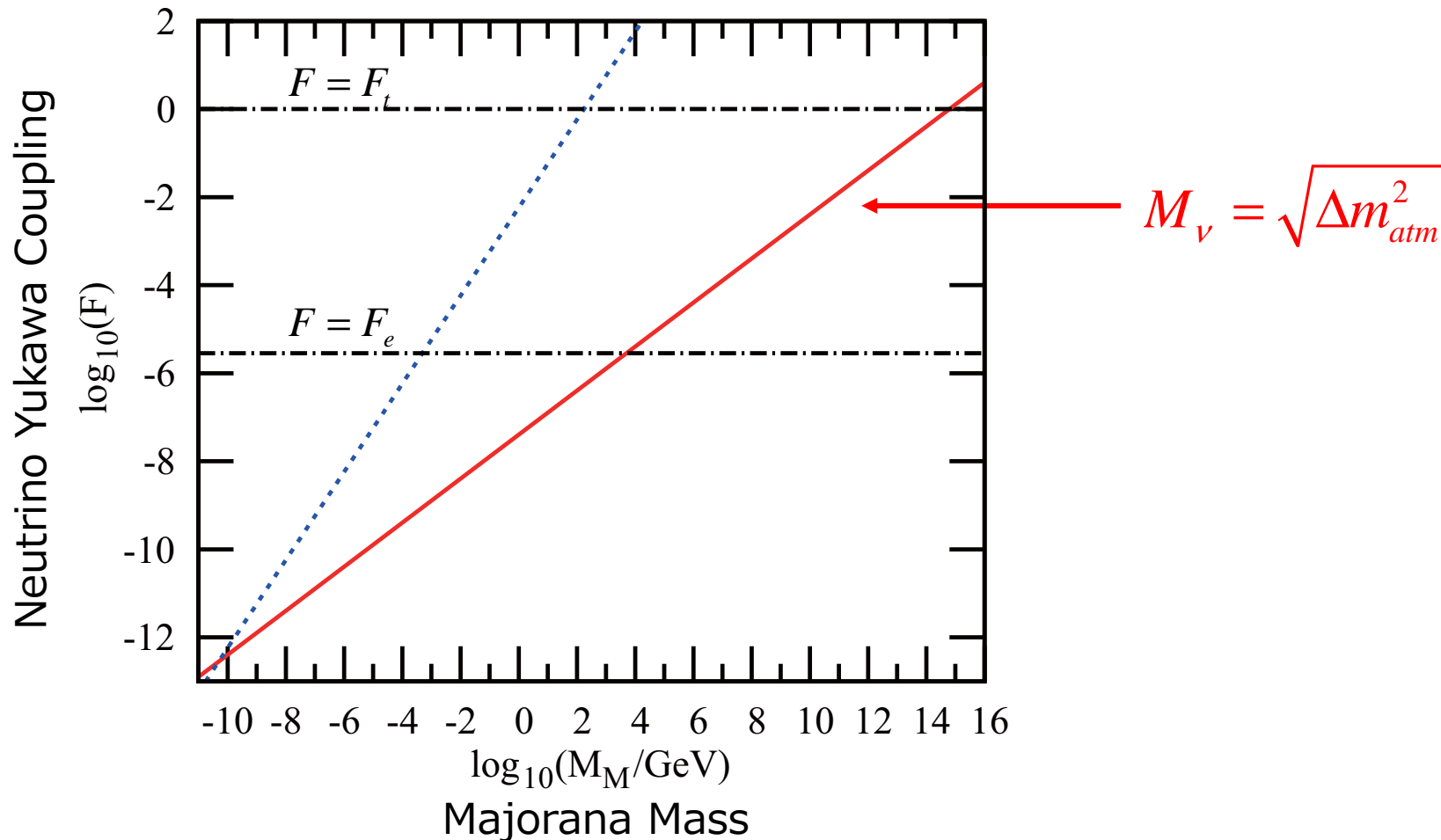
- Mass M_M
- Mixing $\Theta = M_D/M_M$

mixing in CC current $\nu_L = U \nu + \Theta N^c$

Scale of Majorana mass

- The simplest case: one pair of ν_L and ν_R

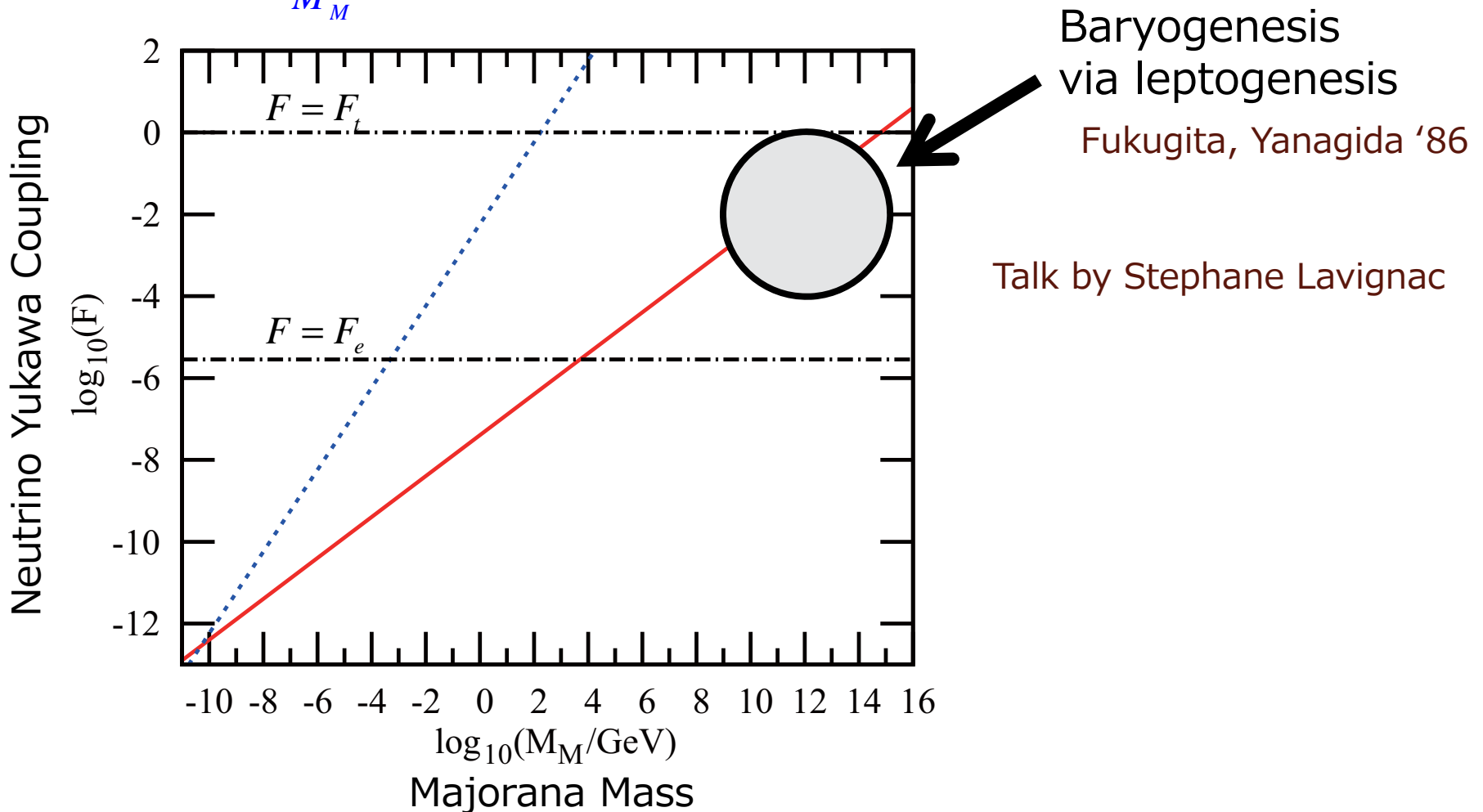
$$M_\nu = -M_D^T \frac{1}{M_M} M_D \Rightarrow F^2 = M_M M_\nu / \langle \Phi \rangle^2$$



Scale of Majorana mass

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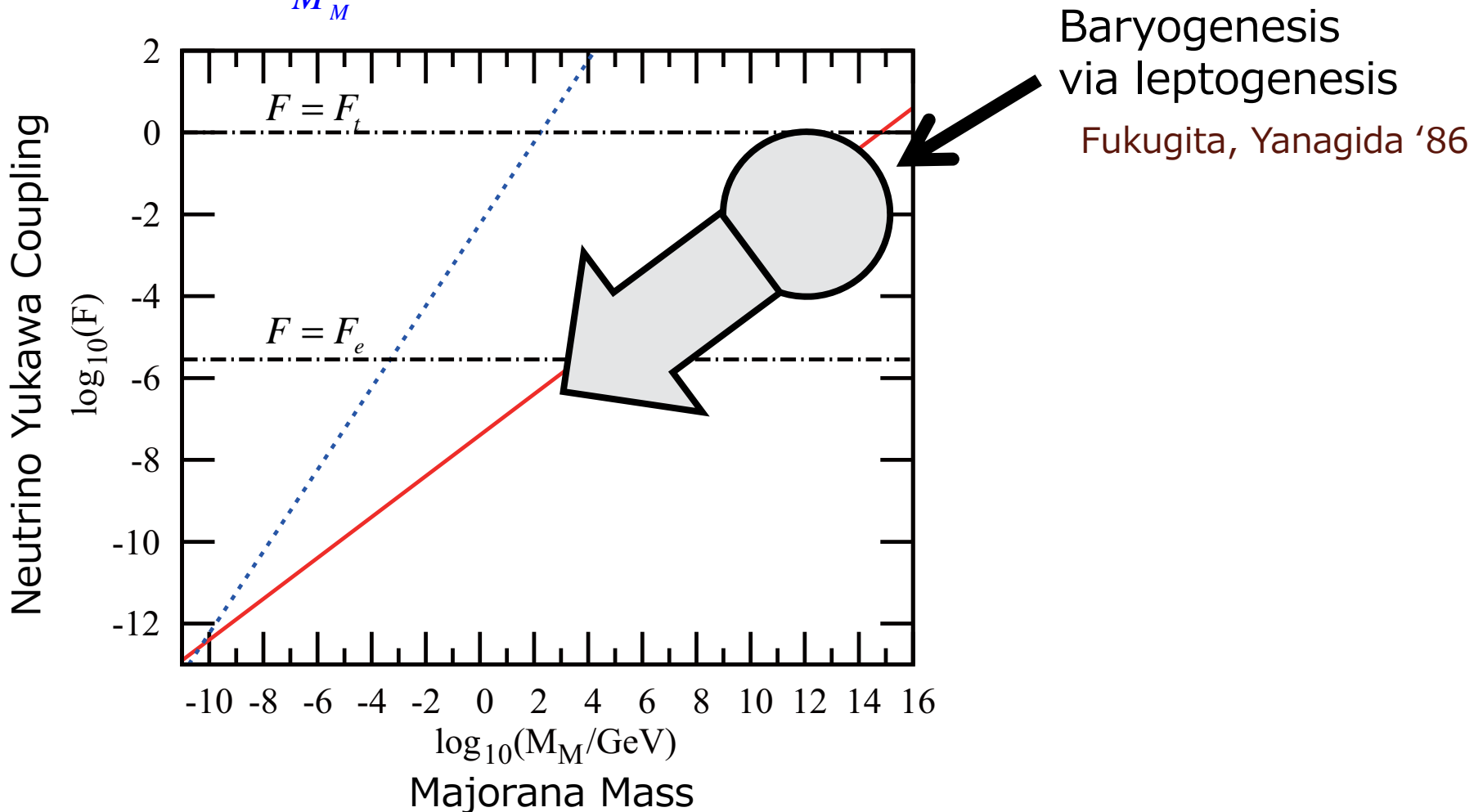
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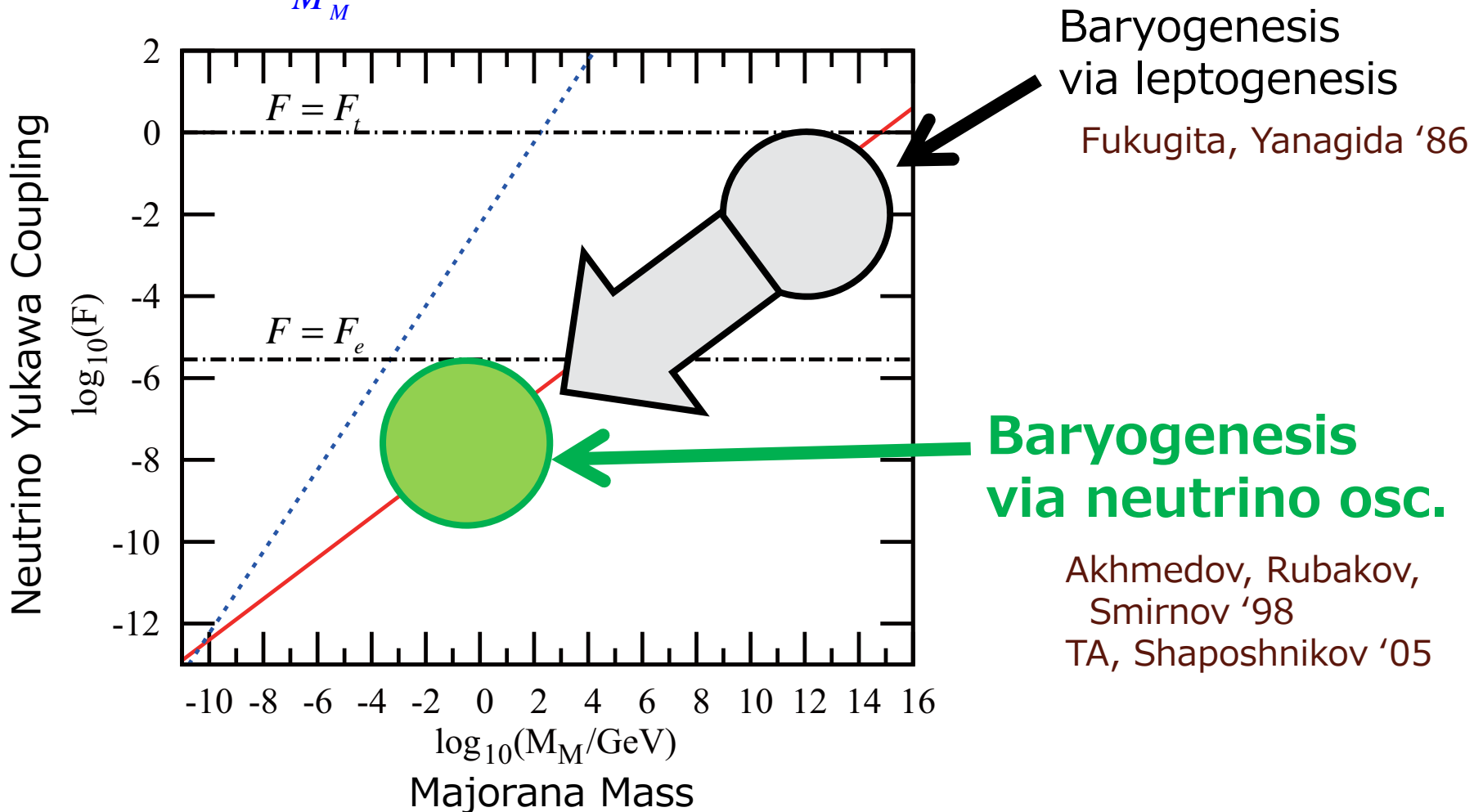
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Scale of Majorana mass

- The simplest case: one pair of ν_L and ν_R

$$M_\nu = -M_D^T \frac{1}{M_M} M_D \Rightarrow F^2 = M_M M_\nu / \langle \Phi \rangle^2$$



- $M_M \lesssim M_W$:no new mass scale between M_W and M_{pl}

$$F \simeq 4 \times 10^{-7} \left(\frac{M_M}{100\text{GeV}} \right)^{\frac{1}{2}} \left(\frac{M_\nu}{0.05\text{eV}} \right)^{\frac{1}{2}}$$

- Lightest one N_1 among 3 heavy neutral leptons with keV mass can be **dark matter !!**
Dodelson, Widrow '94, ...
- Oscillation of heavy neutral leptons N_2 and N_3 can generate **baryon asymmetry of the universe !**
Akhmedov, Rubakov, Smirnov '98, TA, Shaposhnikov '05
- **Dark matter N_1 can be tested by cosmic X-ray obs.**
Talk by Oleg Ruchayskiy
- Physics of heavy neutral leptons N_2 and N_3 can be **tested by experiments !!!**

§

Baryogenesis via Neutrino Oscillation

Conditions for Baryogenesis (Sakharov '67)

- **B and L violations**
 - ▣ B violation due to EW sphaleron
 - ▣ L violation due to Majorana masses
 - negligible for baryogenesis temperature since $M_M \lesssim M_W$
- **C and CP violations**
 - ▣ 1 CP phase in quark sector
 - ▣ 6 CP phases in lepton sector
 - 3 CP phases associated with $N_{2,3}$ are relevant for baryogenesis
- **Out of equilibrium**
 - ▣ No 1st order EW phase transition as in the SM
 - ▣ $N_{2,3}$ can be out of equilibrium for $T \gtrsim M_W$,
if Yukawa couplings are small enough

Baryogenesis via neutrino osc.

- Oscillation of HNLs can be a source of BAU

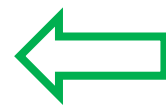
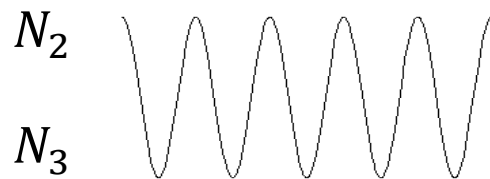
Akhmedov, Rubakov, Smirnov ('98) / TA, Shaposhnikov ('05)

Shaposhnikov ('08), Canetti, Shaposhnikov ('10)

TA, Ishida ('10), Canetti, Drewes, Shaposhnikov ('12), TA, Eijima, Ishida ('12)

Canetti, Drewes, Shaposhnikov ('12), Canetti, Drewes, Frossard, Shaposhnikov ('12)

- ▣ Oscillation starts at $T_{osc} \sim (M_0 M_N \Delta M)^{1/3}$

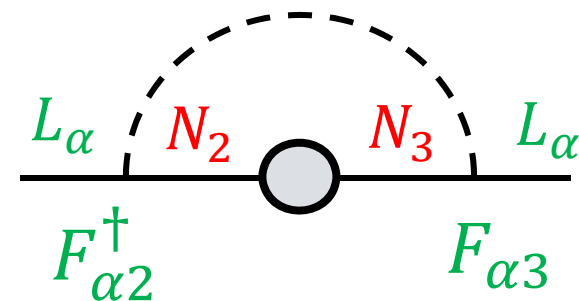
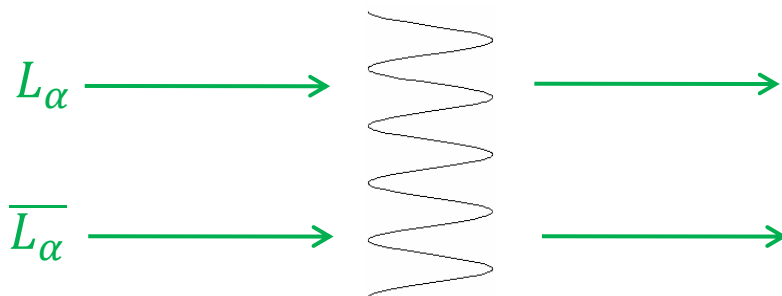


$$V_N = \frac{T^2}{8k} F^\dagger F$$

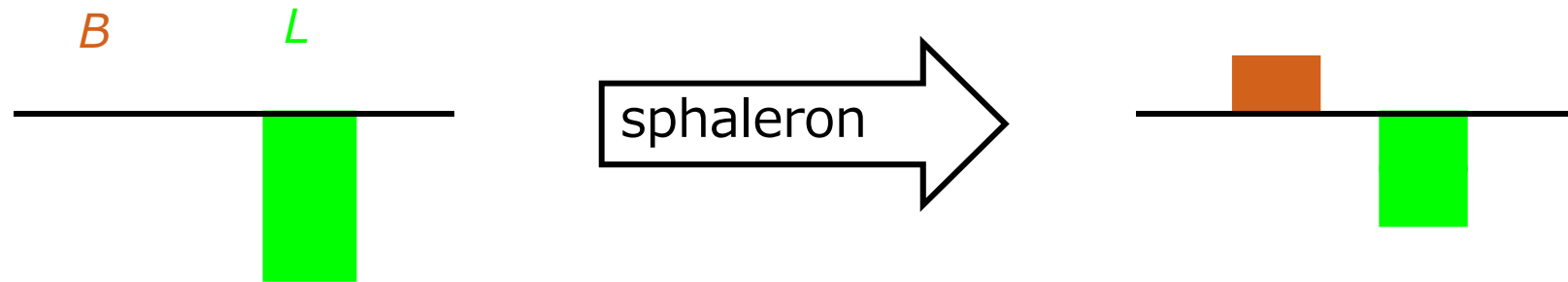
Medium effects



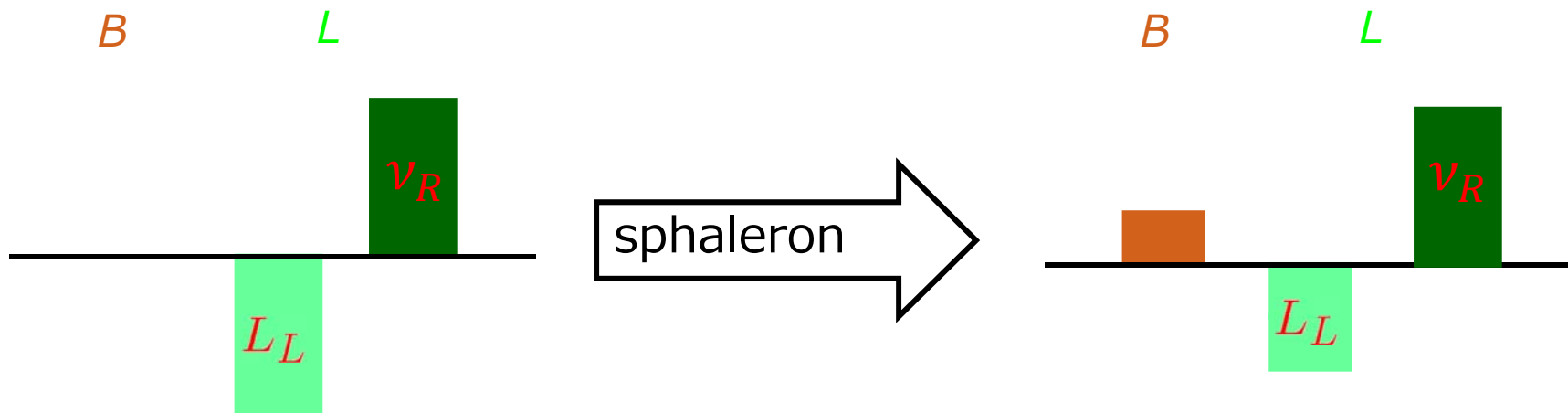
- ▣ Asymmetries are generated since evolution rates of L_α and \overline{L}_α are different due to CPV



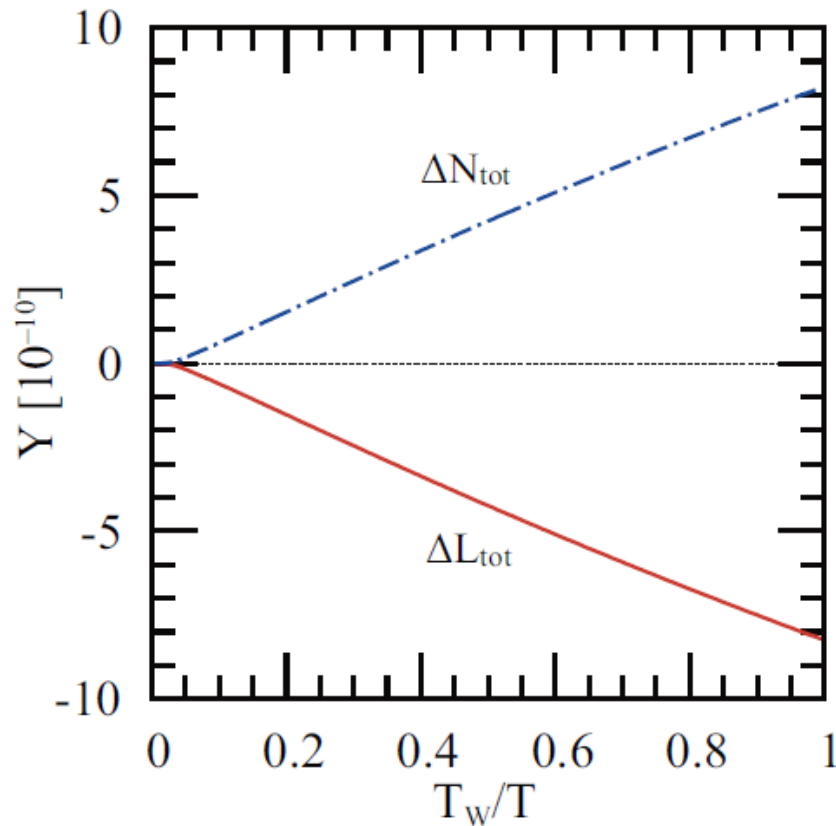
Baryogenesis via leptogenesis



Baryogenesis via neutrino osc.



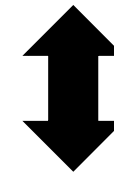
Evolution of asymmetries



Shaleron converts ΔL_{tot} partially into baryon asymmetry

[Kuzmin, Rubakov, Shaposhnikov]

$$\frac{n_B}{s} = -2.5 \times 10^{-4} \Delta L_{tot}(T_W)$$



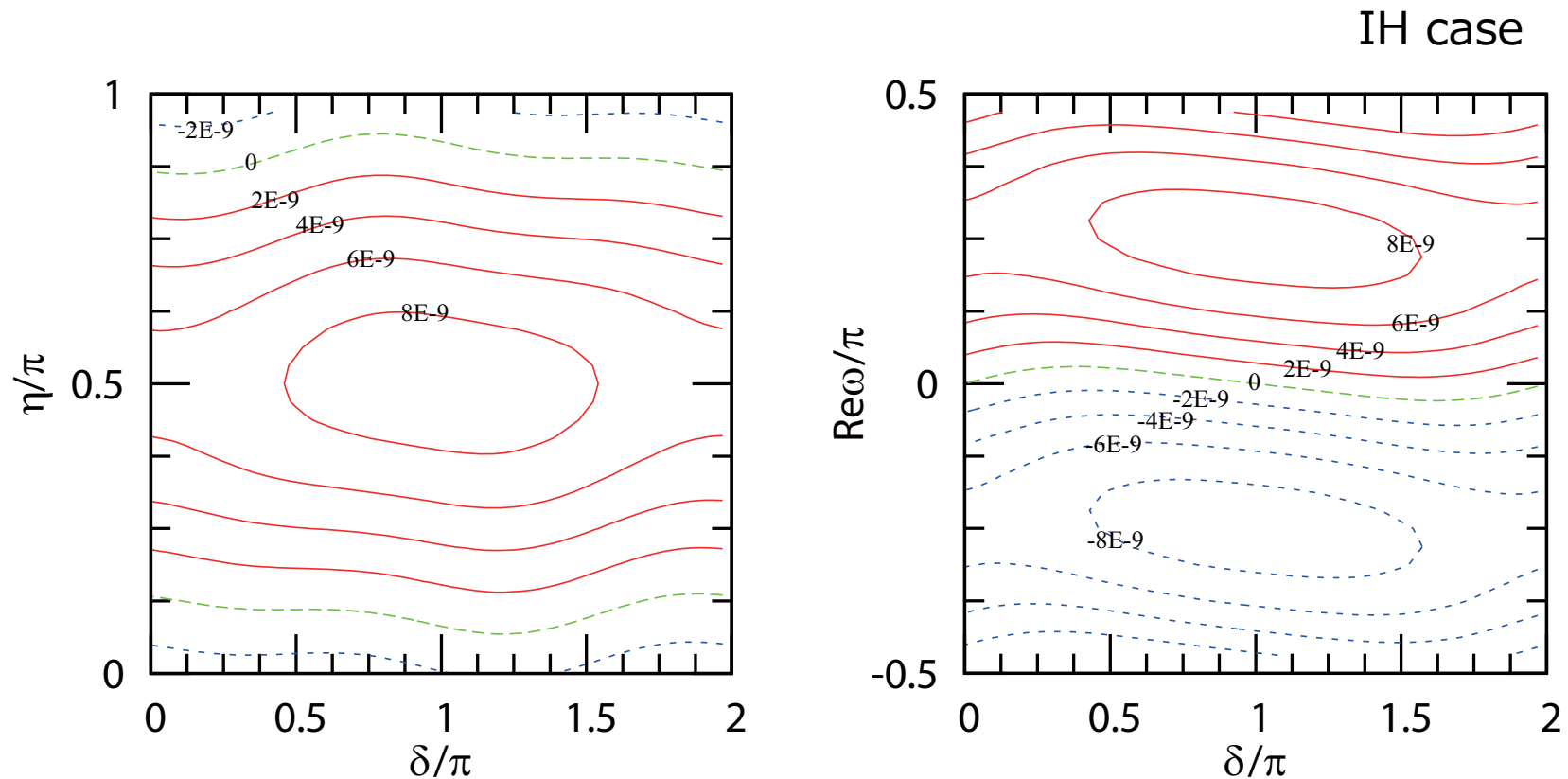
$$\frac{n_B}{s} = (8.55 \sim 9.00) \times 10^{-11}$$

[Planck 2013]

Yield of BAU depends on $F_{\alpha I}$ (even PMNS matrix !!) and masses, especially, CP violating parameters and mass difference

$$T_{osc} \sim (M_0 M_N \Delta M)^{1/3}$$

BAU and Phases δ and η

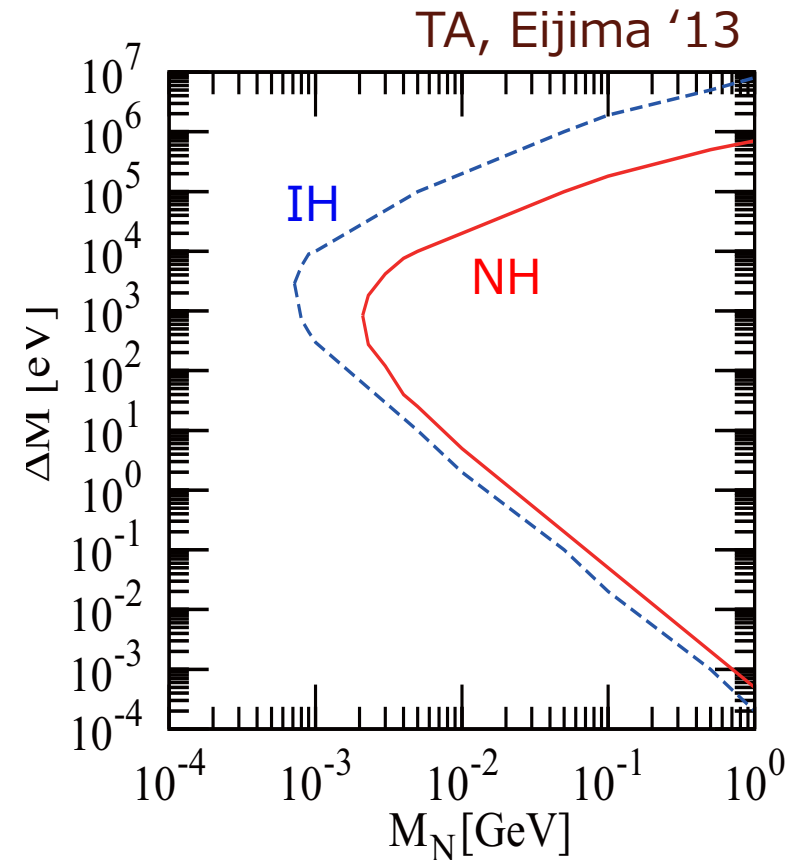
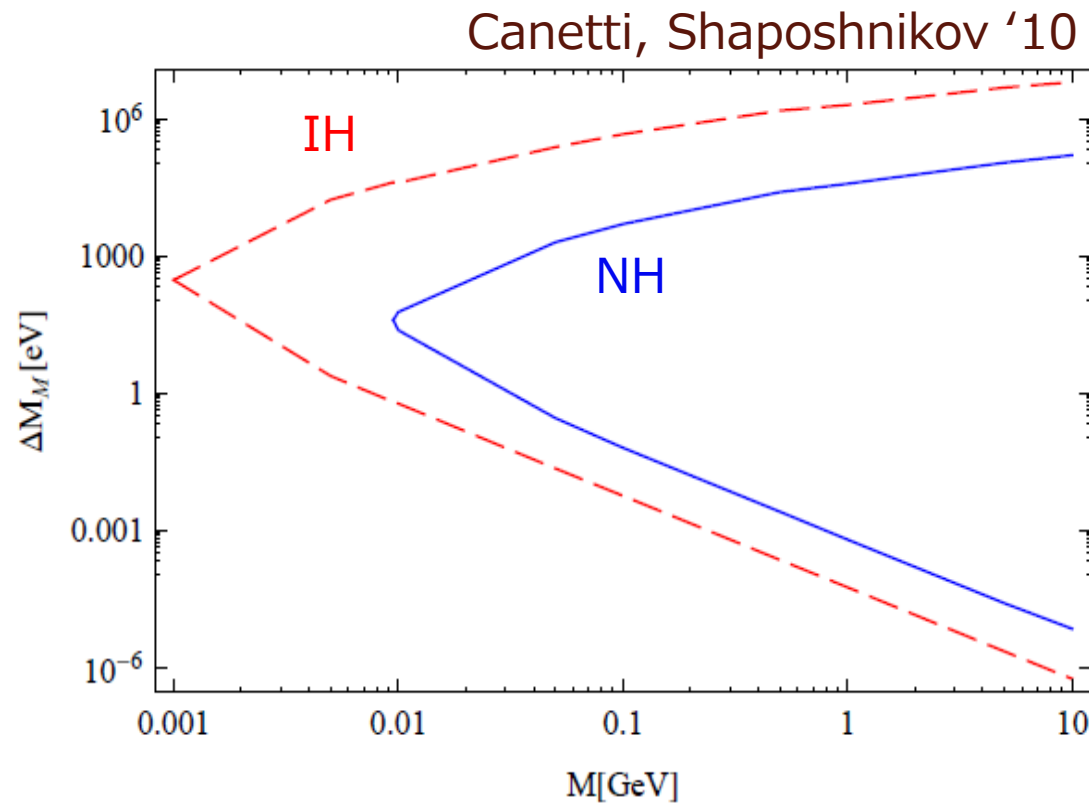


Yield of BAU depends on $F_{\alpha I}$ (even PMNS matrix !!) and masses, especially, CP violating parameters and mass difference

$$T_{\text{osc}} \sim (M_0 M_N \Delta M)^{1/3}$$

Baryogenesis via neutrino osc.

Region accounting for $\frac{n_B}{s} = (8.55-9.00) \times 10^{-11}$



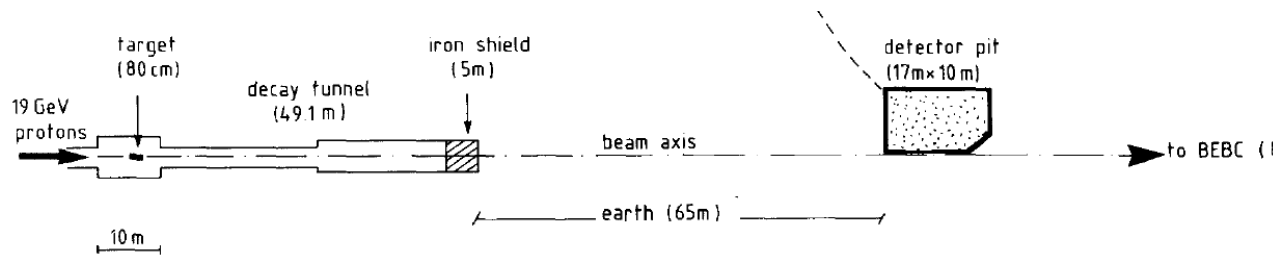
$$M_N > 2.1 \text{ MeV (NH)}$$

$$M_N > 0.7 \text{ MeV (IH)}$$

Direct search experiment

■ PS191 [Bernardi et al '86, '88]

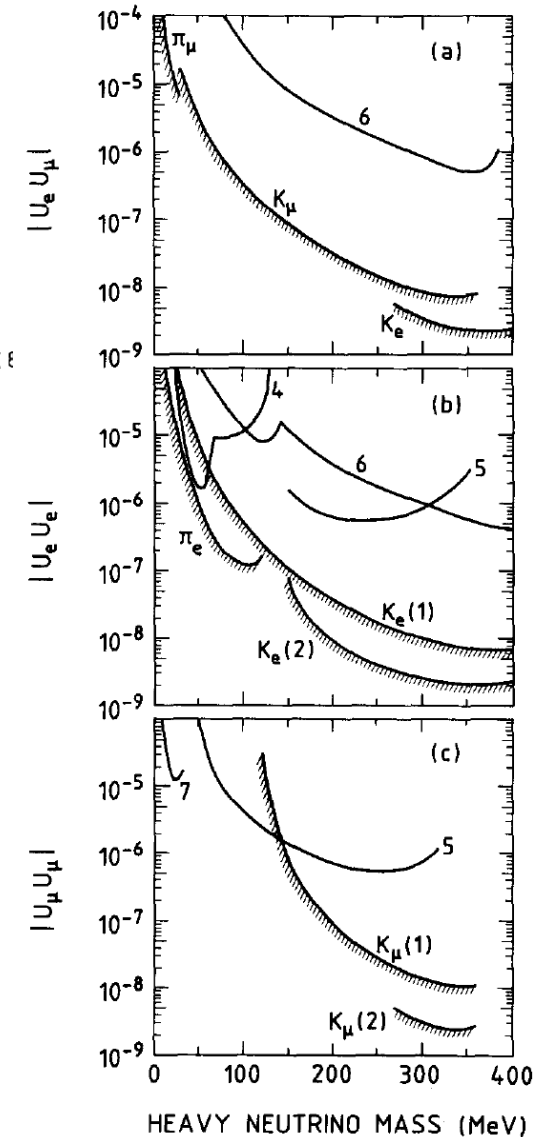
Beam dump experiment
performed at CERN in 1984



- Production $\pi^+, K^+ \rightarrow e^+ N$
- Detection $N \rightarrow \ell^+ \ell^- \nu, \ell^- \pi^+$

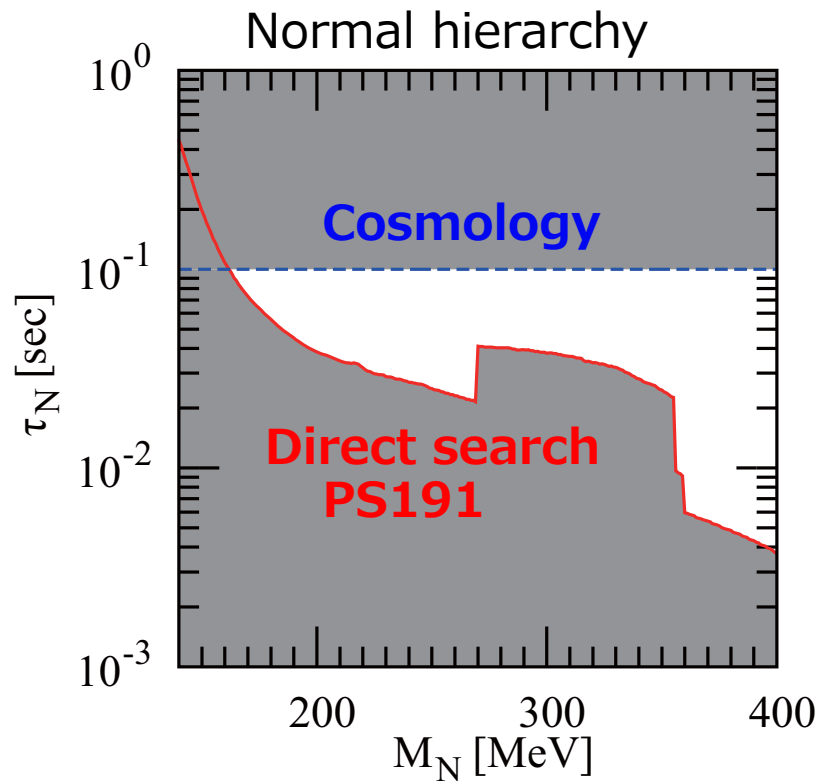
■ Upper bounds mixing elements Θ

⇒ Lower bound on lifetime of $N_{2,3}$

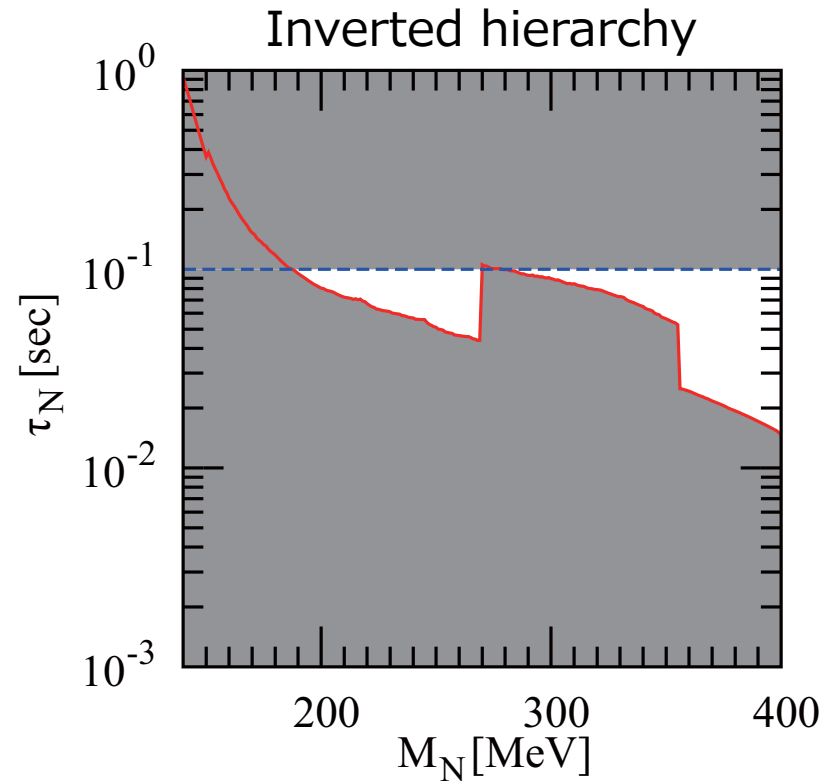


Constraints on light RH neutrinos

TA, Eijima '13



$M_N > 163 \text{ MeV}$



$M_N = 188 - 269 \text{ MeV}$
 $M_N > 285 \text{ MeV}$



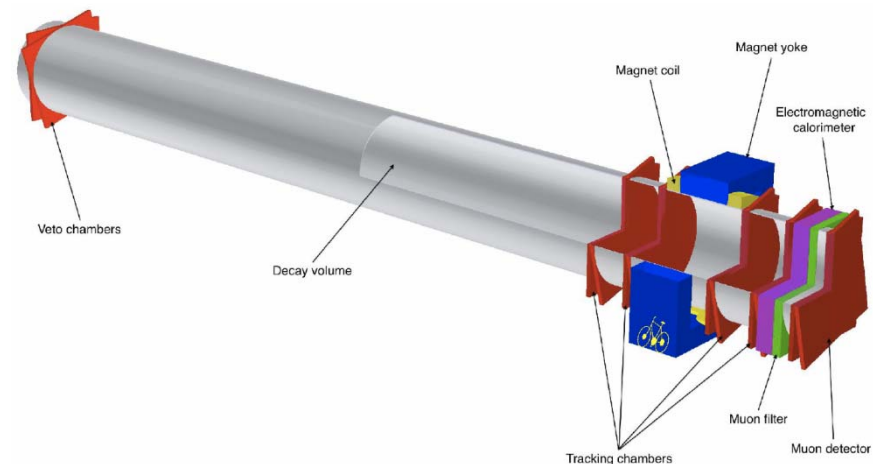
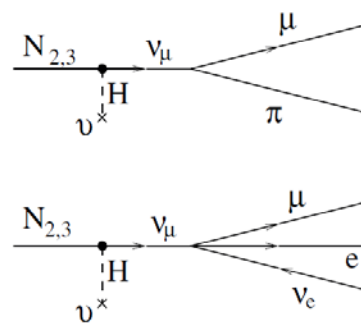
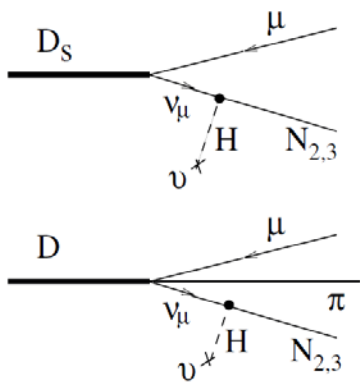
Search for heavy neutral leptons (HNLs)

- A new fixed-target experiment at the CERN SPS accelerator is proposed that will use decays of charm mesons to search for Heavy Neutral Leptons



SHIP - Search for Hidden Particles

CERN, Universität Zürich, EPFL Lausanne, INFN Cagliari,
Università Federico II and INFN Napoli, Imperial College London



Takehiko Asaka (Niigata Univ.)

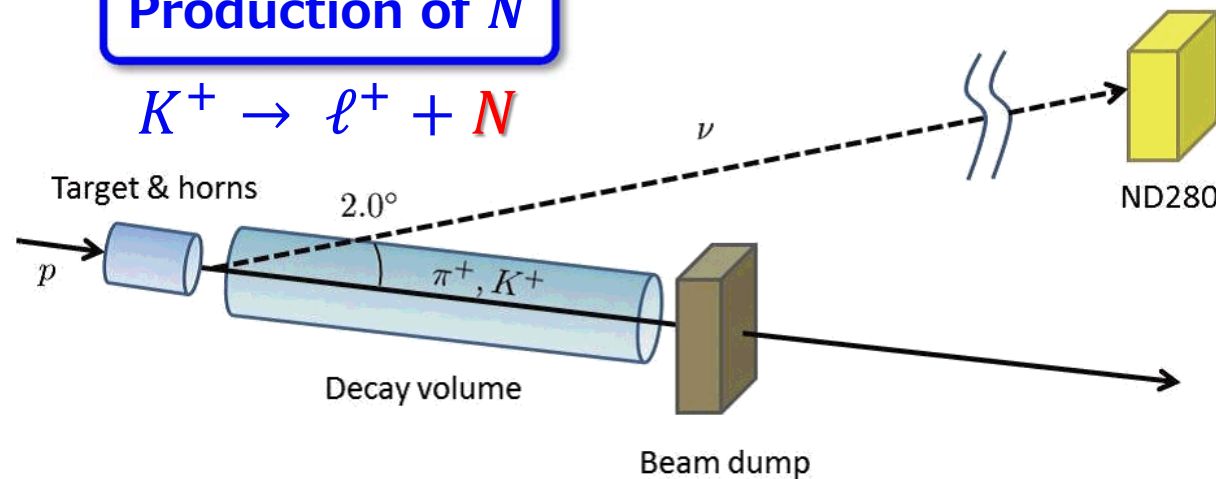
2014/08/01

Search for HNLs at T2K

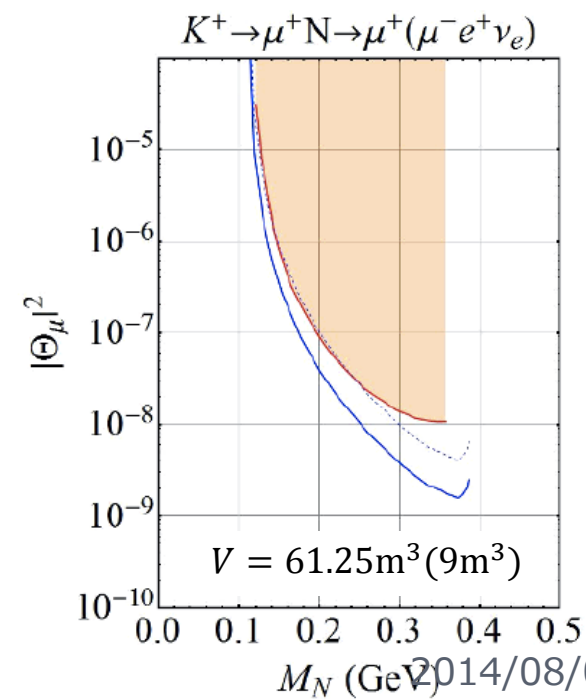
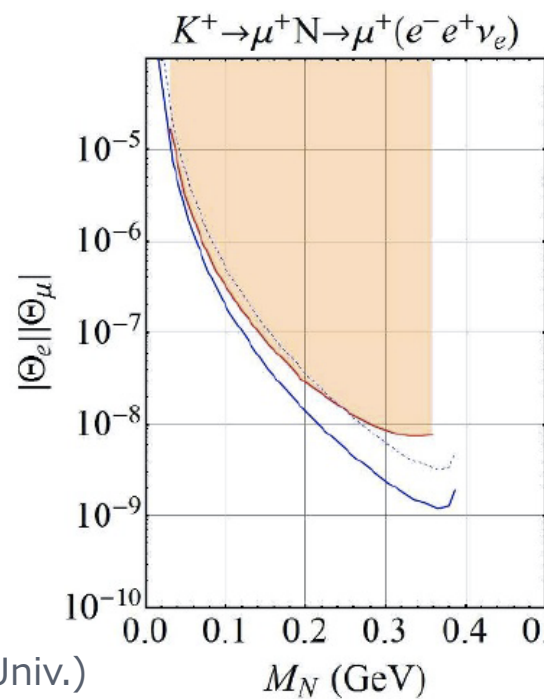
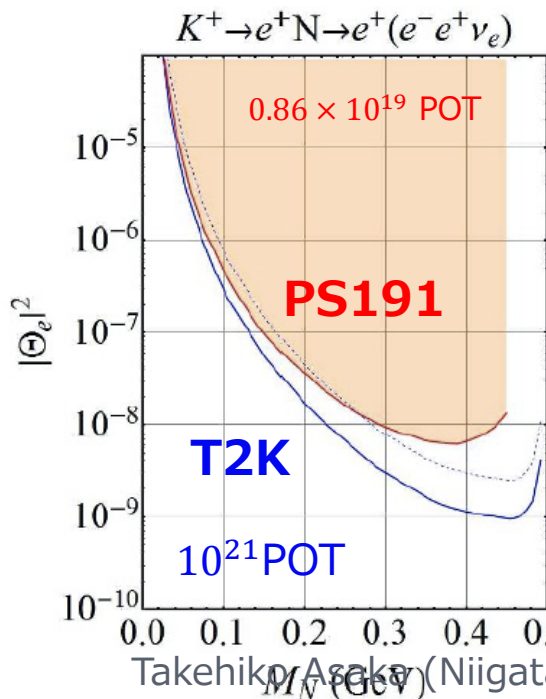
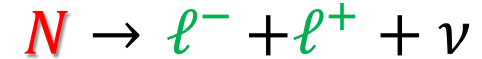
TA, Eijima, Watanabe
[JHEP1303 (2013) 125]

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Production of N

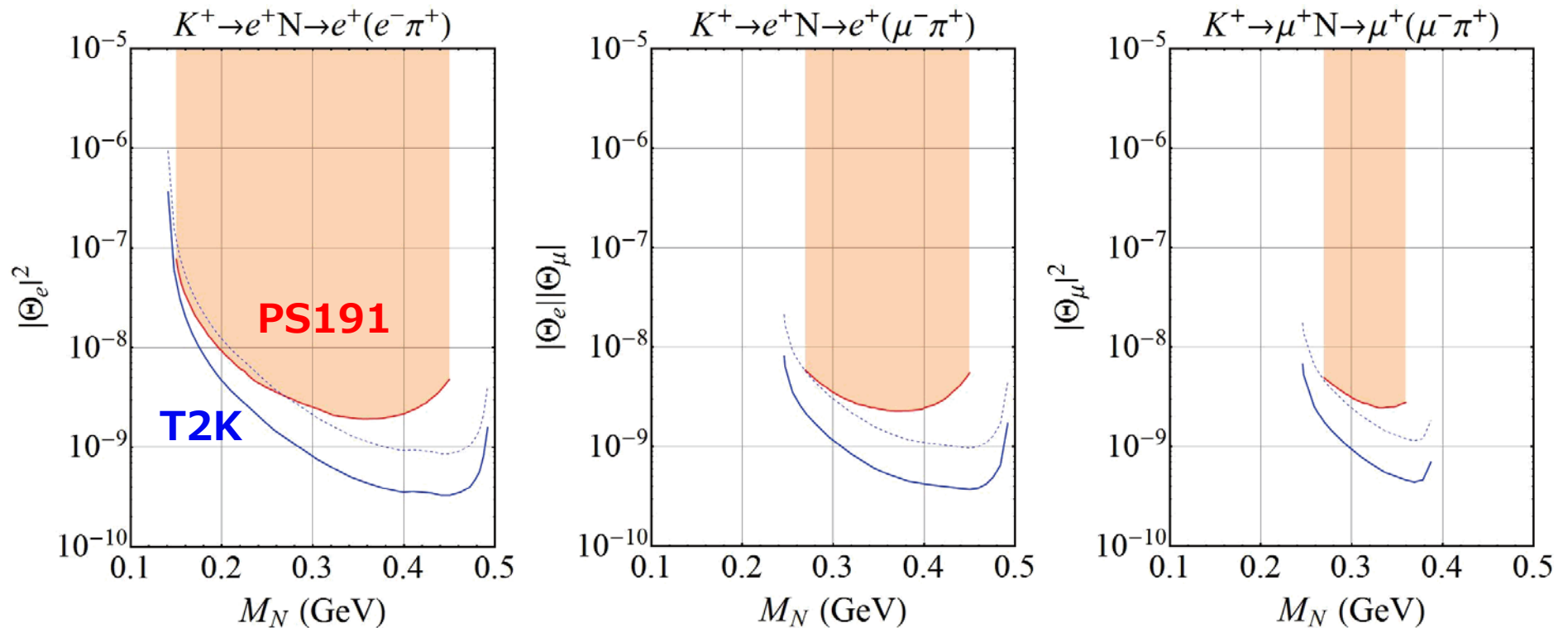


Detection of N



Sensitivity: PS191 vs T2K

TA, Ejima, Watanabe '13



T2K at 10^{21} POT has a better sensitivity than PS191 (0.86×10^{19} POT) !

- We have considered the ν MSM with three right-handed neutrinos which are lighter than weak scale.

- Neutrino masses by seesaw mechanism
- Dark matter (lightest HNL N_1 with \sim keV mass)
- Baryogenesis via neutrino oscillations of $N_{2,3}$
- Direct search of $N_{2,3}$ is possible

- We have found the possible region for neutrino oscillations and BAU, allowed from search and cosmological constraints.

- $M_N > 163$ MeV (NH)
 $M_N = 188 - 269$ MeV and $M_N > 285$ MeV (IH)

- ***Search for these heavy neutral leptons are crucial to solve the origin of neutrino masses as well as the mysteries of our universe, DM and BAU !!!***



Backup

BBN constraint on lifetime

- Long-lived $N_{2,3}$ may spoil the success of BBN

- ▣ Speed up the expansion of the universe

- $\rho_{\text{tot}} = \rho_{\text{MSM}} + \rho_{N_{2,3}} \Rightarrow H^2 = \frac{\rho_{\text{tot}}}{3 M_P^2}$

- p-n conv. decouples earlier \Rightarrow overproduction of ${}^4\text{He}$

$$n + \nu \leftrightarrow p + e^-, \dots$$

- ▣ Distortion of spectrum of active neutrinos

- $N_{2,3} \rightarrow \nu \bar{\nu} \nu, e^+ e^- \nu, \dots$

- Additional neutrinos may not be thermalized

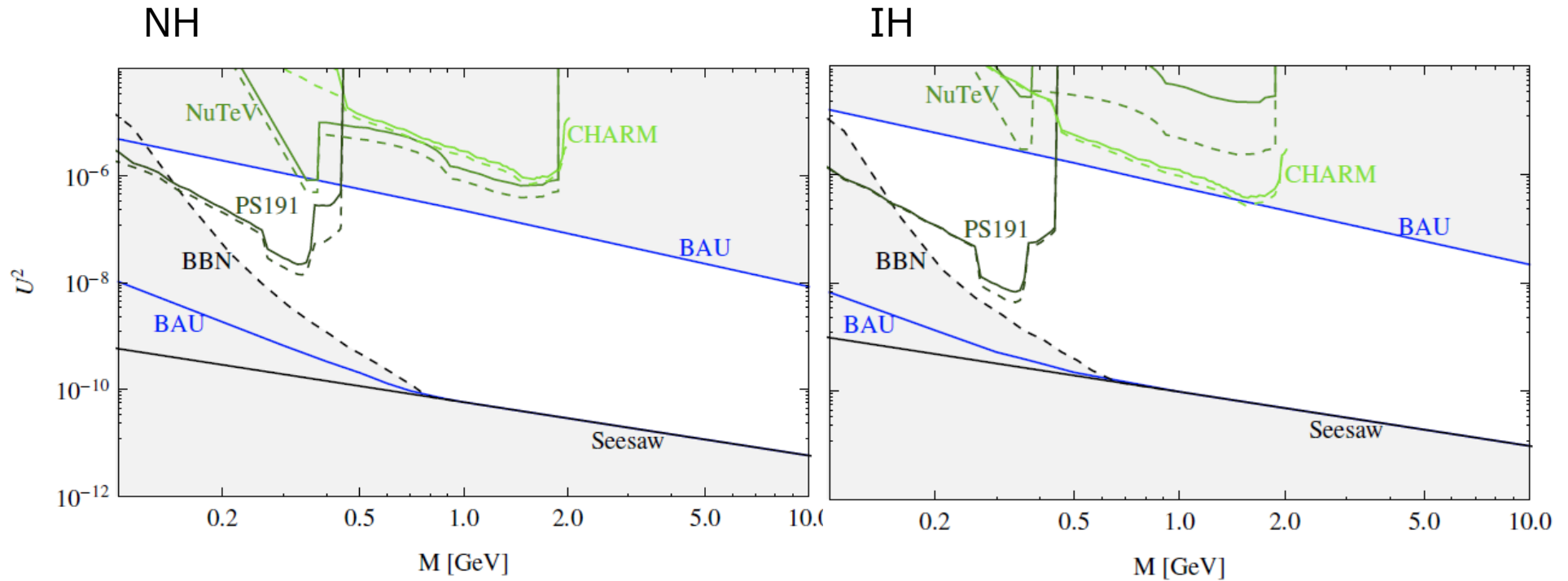
\Rightarrow Upper bound on lifetime of $N_{2,3}$

- Dolgov, Hansen, Rafflet, Semikoz ('00)

- ▣ One family case: $\tau_N < 0.1 \text{ sec for } M_N > m_\pi$

Constraints on HNLs N_2 and N_3

Canetti, Drewes, Frossard, Shaposhnikov '13



See Gorbunov and Shaposhnikov '07 [arXiv:0705.1729]
Atre, Han, Pascoli, Zhang '09 [arXiv:0901.3589]

Neutrino Yukawa couplings for $N_{2,3}$

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$$F = U_{\text{PMNS}} D_\nu^{1/2} \Omega D_N^{1/2} / \langle \Phi \rangle \quad (\text{in NH})$$

[Casas, Ibarra '01]

Parameters of active neutrinos

$D_\nu^{1/2} = \text{diag}(\sqrt{m_1} = 0, \sqrt{m_2}, \sqrt{m_3})$: active ν masses

$$U_{\text{PMNS}} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -c_{23}s_{12} - s_{23}c_{12}s_{13}e^{i\delta} & c_{23}c_{12} - s_{23}s_{12}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{23}s_{12} - c_{23}c_{12}s_{13}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{12}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} 1 \\ e^{i\eta} \\ 1 \end{pmatrix}$$

Dirac phase δ

Majorana phase η

Parameters of sterile neutrinos

$D_N^{1/2} = \text{diag}(\sqrt{M_2}, \sqrt{M_3})$: sterile ν masses

$$\Omega = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix} \quad \begin{array}{l} \omega: \text{complex number} \\ \xi = \pm 1 \end{array}$$

$\text{Im}\omega$